

Quantum field theory
Exercises 4. Solutions
2005-11-21

• **Exercise 4.1.**

- In the rest frame $E = m$, $\mathbf{p} = 0$. Performing a boost with rapidity η (along x axis for example) we have $E = m \cosh \eta$, $p_x = m \sinh \eta$. Thus, $(E + p)/(E - p) = e^{2\eta}$.
- Performing a second boost with rapidity η' along x we get $E \rightarrow E \cosh \eta' + p \sinh \eta'$ and $p \rightarrow E \sinh \eta' + p \cosh \eta'$, so

$$\begin{aligned} e^{2\eta} &\rightarrow \frac{(E \cosh \eta' + p \sinh \eta') + (E \sinh \eta' + p \cosh \eta')}{(E \cosh \eta' + p \sinh \eta') - (E \sinh \eta' + p \cosh \eta')} \\ &= e^{2\eta'} \frac{E + p}{E - p} = e^{2(\eta + \eta')} . \end{aligned}$$

• **Exercise 4.2.**

Let $v^0 = \xi_R^\dagger \psi_R$ and $v^i = \xi_R^\dagger \sigma^i \psi_R$. Using the rules for spinor transformations under infinitesimal boosts and rotations

$$\psi_R \rightarrow e^{(-i\theta_i + \eta_i) \frac{\sigma^i}{2}} \psi_R = (1 + (-i\theta_i + \eta_i) \frac{\sigma^i}{2} + \dots) \psi_R ,$$

we get for v^0

$$\begin{aligned} v^0 &\rightarrow \xi_R^\dagger (1 + (+i\theta_i + \eta_i) \frac{\sigma^i}{2}) (1 + (-i\theta_i + \eta_i) \frac{\sigma^i}{2}) \eta_R = \xi_R^\dagger (1 + \eta_i \sigma^i + \dots) \eta_R \\ &= v^0 + \eta_i v^i + \dots \end{aligned}$$

and for v^i

$$\begin{aligned} v^i &\rightarrow \xi_R^\dagger (1 + (i\theta_j + \eta_j) \frac{\sigma^j}{2}) \sigma^i (1 + (-i\theta_k + \eta_k) \frac{\sigma^k}{2}) \eta_R \\ &= \xi_R^\dagger (\sigma^i + \frac{i}{2} \theta_j [\sigma^j, \sigma^i] + \frac{1}{2} \eta_j \{\sigma^j, \sigma^i\} + \dots) \eta_R = \xi_R^\dagger (\sigma^i - \theta_j \varepsilon^{jik} \sigma^k + \eta_j \delta^{ji} + \dots) \eta_R \\ &= v^i - \theta_j \varepsilon^{jik} v^k + \eta_j \delta^{ji} v^0 + \dots . \end{aligned}$$

This is exactly transformation properties of a (contravariant) four-vector under infinitesimal boost η_i (boost of rapidity $|\eta|$ in the direction of $\eta^i/|\eta|$). We also checked that it has correct properties under usual rotations θ^i .

Calculations for $\xi_L^\dagger \bar{\sigma}^\mu \eta_L$ are completely analogous.

You can also check everything explicitly for finite boosts, using the formula

$$e^{\eta_i \sigma^i} = \cosh \eta + \frac{\eta_i}{\eta} \sigma^i \sinh \eta ,$$

where $\eta = \sqrt{\eta^i \eta^i}$ is rapidity of the boost and η_i/η — its direction.

• **Exercise 4.3.**

Transformations of the double-contravariant tensor are

$$F^{\mu\nu} \rightarrow \Lambda^{\mu}_{\rho} \Lambda^{\nu}_{\sigma} F^{\rho\sigma} ,$$

where $\Lambda^{\mu}_{\rho} = \exp\{-(i/2)\omega_{\alpha\beta}(J^{\alpha\beta})^{\mu}_{\rho}\}$, and

$$(J^{\alpha\beta})^{\mu}_{\rho} = i(\eta^{\alpha\mu}\delta^{\beta}_{\rho} - \eta^{\beta\mu}\delta^{\alpha}_{\rho}) .$$

Then, to the first order

$$\delta F^{\mu\nu} = \omega^{\mu}_{\rho} F^{\rho\nu} - \omega^{\nu}_{\rho} F^{\rho\mu} .$$

Using the explicit form $\omega_{i0} = -\omega^{i0} = -\eta^i$ and $\omega_{ij} = \omega^{ij} = \epsilon^{ijk}\theta^k$ one gets for **B** and **E**

$$\begin{aligned} \delta \mathbf{E} &= -[\boldsymbol{\eta} \times \mathbf{B}] + [\boldsymbol{\theta} \times \mathbf{E}] , \\ \delta \mathbf{B} &= +[\boldsymbol{\eta} \times \mathbf{E}] + [\boldsymbol{\theta} \times \mathbf{B}] . \end{aligned}$$