

# Quantum Field Theory

## Set 9

### Exercise 1: $SU(2)$ Noether's current

Consider a pair of complex scalar fields  $\Phi = \begin{pmatrix} \phi_1 \\ \phi_2 \end{pmatrix}$  transforming according to the representation  $j = 1/2$  of a global  $SU(2)$  symmetry (do not confuse this symmetry with the Lorentz transformations: they are completely uncorrelated):

$$\begin{aligned} x^\mu &\longrightarrow x'^\mu = x^\mu, \\ \Phi(x) &\longrightarrow \Phi'(x') = \mathcal{U}\Phi(x), \\ \phi_a(x) &\longrightarrow \phi'_a(x') = \mathcal{U}_a{}^b \phi_b(x), \quad a, b = 1, 2. \end{aligned}$$

where  $\mathcal{U}$  is a  $SU(2)$  matrix. Given the Lagrangian density

$$\mathcal{L} = \partial^\mu \Phi^\dagger \partial_\mu \Phi - m^2 \Phi^\dagger \Phi - \frac{\lambda}{2} (\Phi^\dagger \Phi)^2,$$

- Show that the Lagrangian density is invariant under  $SU(2)$  transformations.
- Compute the Noether current for this symmetry.
- Add to the above Lagrangian the interaction with a real *triplet*  $\vec{A} = (A_1, A_2, A_3)$  of  $SU(2)$

$$\tilde{\mathcal{L}} = \mathcal{L} + \frac{1}{2} \partial^\mu A^T \partial_\mu A + A_i \Phi^\dagger \sigma^i \Phi. \quad (1)$$

Is the new Lagrangian density invariant under  $SU(2)$  transformations involving the field  $\Phi$  and  $\vec{A}$ ?

### Exercise 2: Invariant measure

Consider the measure on three dimensional momentum space  $\frac{d^3 k}{(2\pi)^3 2k^0}$ , where  $k_0 = \sqrt{|\vec{k}|^2 + m^2}$ .

- Performing a boost in a particular direction, say  $\hat{k}_3$ , show the invariance of the measure under Lorentz transformations.  
(Hint: use the form of the boost transformations  $k'_0 = \cosh(\eta) k_0 + \sinh(\eta) k_3$ ,  $k'_3 = \cosh(\eta) k_3 + \sinh(\eta) k_0$ ).
- Show that the product  $\delta^3(\vec{k}) d^3 k$  is invariant under Lorentz transformation and deduce that the distribution  $(2\pi)^3 2k^0 \delta^3(\vec{k})$  is invariant as well.

### Exercise 3: Noether's current: a different approach

Given a Lagrangian density  $\mathcal{L}(\phi_a, \partial_\mu \phi_b)$  consider the infinitesimal symmetry transformation defined by

$$\begin{aligned} x^\mu &\longrightarrow x'^\mu = f(x) \simeq x^\mu - \epsilon_i^\mu(x) \beta^i(x), \\ \phi_a(x) &\longrightarrow \phi'_a(x') = D[\phi](x) \simeq \phi_a(x') + \beta^i(x') \Delta_{ai}(x'), \end{aligned}$$

where the infinitesimal parameters of the transformation depend explicitly on the spacetime coordinates. The request that the Lagrangian is invariant under the transformation where  $\beta$  is a *constant* parameter implies the equation (see the proof of Noether's theorem)

$$\frac{\partial \mathcal{L}}{\partial \phi_a} \Delta_{ai}(x) + \frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi_a)} \partial_\mu \Delta_{ai}(x) - \partial_\mu (\mathcal{L} \epsilon_i^\mu) = 0$$

- Using the above equation show that the change in the action is given by

$$\int_{\Omega'} \mathcal{L}[\phi'_a](x') d^4 x' - \int_{\Omega} \mathcal{L}[\phi_a](x) d^4 x = \int_{\Omega} J_i^\mu \partial_\mu \beta^i(x) d^4 x,$$

where  $J_i^\mu$  is the usual Noether's current defined for the *global* symmetry (the same symmetry but with constant parameters).

- Apply explicitly this procedure to find the Noether's current associated to the  $U(1)$  symmetry in a complex scalar field theory.

#### Exercise 4: Decomposition in Fourier modes

Given the decomposition of a real scalar field in a finite cubic volume  $V = L^3$

$$\phi(t, \vec{x}) = \sum_{\vec{n} \in \mathbb{Z}^3} \frac{1}{\sqrt{V}} \phi_n(t) e^{i \frac{2\pi}{L} \vec{n} \cdot \vec{x}}$$

- Expand  $\int d^3 x (\vec{\nabla} \phi)^2(t, \vec{x})$  in Fourier modes  $\phi_n(t)$ .

#### Exercise 5: Ladder operators

Given the algebra of the ladder operators (non relativistic normalization)

$$[a(\vec{q}), a^\dagger(\vec{p})] = (2\pi)^3 \delta^3(\vec{q} - \vec{p}),$$

and the expression of the Hamiltonian

$$H = \int \frac{d^3 k}{(2\pi)^3} \omega(\vec{k}) a^\dagger(\vec{k}) a(\vec{k}),$$

compute the commutation relations

$$[H, a(\vec{p})] = ? \quad [H, a^\dagger(\vec{p})] = ?$$