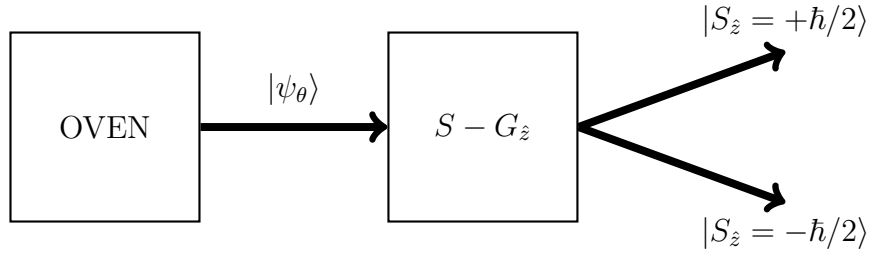


1. The initial state of an electron leaving the oven is

$$|\psi_\theta\rangle = \cos(\theta)|S_z = +\hbar/2\rangle + \sin(\theta)|S_z = -\hbar/2\rangle ,$$

with θ distributed according the probability density function

$$\rho(\theta) = \begin{cases} \frac{1}{2\pi} & \text{if } \theta \in [0, 2\pi] \\ 0 & \text{otherwise.} \end{cases}$$



Hence, the probability that the first Stern and Gerlach experiment measures $S_z = +\hbar/2$, for a fixed θ , is given by

$$\mathcal{P}(S_z = +\hbar/2; \theta) = |\langle S_z = +\hbar/2 | \psi_\theta \rangle|^2 = \cos^2(\theta) ,$$

while the probability that the first Stern and Gerlach experiment measures $S_z = -\hbar/2$, for a fixed θ , is given by

$$\mathcal{P}(S_z = -\hbar/2; \theta) = |\langle S_z = -\hbar/2 | \psi_\theta \rangle|^2 = \sin^2(\theta) .$$

Averaging over $\rho(\theta)$ yields how the electrons distribute among the two beams:

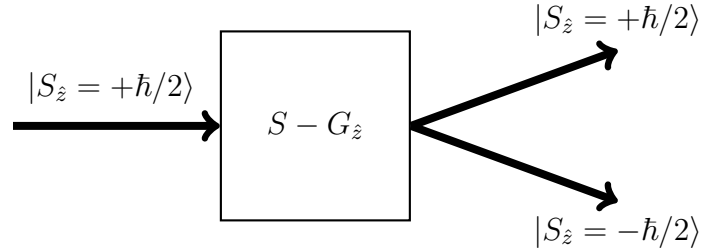
$$\begin{aligned} \mathcal{P}(S_z = +\hbar/2) &= \int_{-\infty}^{\infty} \mathcal{P}(S_z = +\hbar/2; \theta) \rho(\theta) d\theta \\ &= \frac{1}{2\pi} \int_0^{2\pi} \cos^2(\theta) d\theta \\ &= \frac{1}{2\pi} \int_0^{2\pi} \frac{1 + \cos(2\theta)}{2} d\theta \\ &= \frac{1}{2\pi} \left[\frac{\theta}{2} + \frac{\sin(2\theta)}{4} \right]_0^{2\pi} \\ &= \frac{1}{2} ; \end{aligned}$$

$$\begin{aligned}
\mathcal{P}(S_{\hat{z}} = -\hbar/2) &= \int_{-\infty}^{\infty} \mathcal{P}(S_{\hat{z}} = -\hbar/2; \theta) \rho(\theta) d\theta \\
&= \frac{1}{2\pi} \int_0^{2\pi} \sin^2(\theta) d\theta \\
&= \frac{1}{2\pi} \int_0^{2\pi} \frac{1 - \cos(2\theta)}{2} d\theta \\
&= \frac{1}{2\pi} \left[\frac{\theta}{2} - \frac{\sin(2\theta)}{4} \right]_0^{2\pi} \\
&= \frac{1}{2} .
\end{aligned}$$

2. Now, the initial state is $|S_{\hat{z}} = +\hbar/2\rangle$, therefore the probabilities that the second Stern and Gerlach experiment measures $S_{\hat{z}} = -\hbar/2$ and $S_{\hat{z}} = +\hbar/2$ are

$$\begin{aligned}
\mathcal{P}(S_{\hat{z}} = +\hbar/2) &= |\langle S_{\hat{z}} = +\hbar/2 | S_{\hat{z}} = +\hbar/2 \rangle|^2 = 1 , \\
\mathcal{P}(S_{\hat{z}} = -\hbar/2) &= |\langle S_{\hat{z}} = -\hbar/2 | S_{\hat{z}} = +\hbar/2 \rangle|^2 = 0 ,
\end{aligned}$$

respectively.

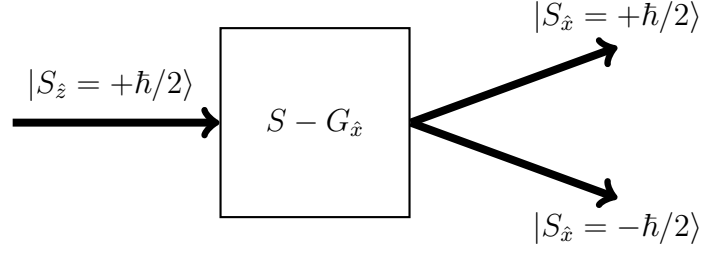


3. The first step to answer to this point is to compute $|S_{\hat{x}} = +\hbar/2\rangle$ and $|S_{\hat{x}} = -\hbar/2\rangle$. These two kets are defined by the equations

$$S_{\hat{x}} |S_{\hat{x}} = \pm \hbar/2\rangle = \pm \hbar/2 |S_{\hat{x}} = \pm \hbar/2\rangle ,$$

which give

$$\begin{aligned}
|S_{\hat{x}} = +\hbar/2\rangle &= \frac{|S_{\hat{z}} = +\hbar/2\rangle + |S_{\hat{z}} = -\hbar/2\rangle}{\sqrt{2}} , \\
|S_{\hat{x}} = -\hbar/2\rangle &= \frac{|S_{\hat{z}} = +\hbar/2\rangle - |S_{\hat{z}} = -\hbar/2\rangle}{\sqrt{2}} .
\end{aligned}$$



Finally, the probabilities that the third Stern and Gerlach experiment measures $S_x = +\hbar/2$ and $S_x = -\hbar/2$ are

$$\begin{aligned}
 \mathcal{P}(S_x = +\hbar/2) &= |\langle S_x = +\hbar/2 | S_z = +\hbar/2 \rangle|^2 \\
 &= \left| \frac{1}{\sqrt{2}} \left(\langle S_z = +\hbar/2 | + \langle S_z = -\hbar/2 | \right) | S_z = +\hbar/2 \rangle \right|^2 \\
 &= \left| \frac{1}{\sqrt{2}} \left(\langle S_z = +\hbar/2 | S_z = +\hbar/2 \rangle + \langle S_z = -\hbar/2 | S_z = +\hbar/2 \rangle \right) \right|^2 \\
 &= \frac{1}{2} ; \\
 \mathcal{P}(S_x = -\hbar/2) &= |\langle S_x = -\hbar/2 | S_z = +\hbar/2 \rangle|^2 \\
 &= \left| \frac{1}{\sqrt{2}} \left(\langle S_z = +\hbar/2 | - \langle S_z = -\hbar/2 | \right) | S_z = +\hbar/2 \rangle \right|^2 \\
 &= \left| \frac{1}{\sqrt{2}} \left(\langle S_z = +\hbar/2 | S_z = +\hbar/2 \rangle - \langle S_z = -\hbar/2 | S_z = +\hbar/2 \rangle \right) \right|^2 \\
 &= \frac{1}{2} .
 \end{aligned}$$

4. In this case, the initial state is $|S_x = +\hbar/2\rangle$, therefore, the probabilities that the fourth Stern and Gerlach experiment measures $S_z = +\hbar/2$ and $S_z = -\hbar/2$ are

$S_{\hat{z}} = -\hbar/2$ are

$$\begin{aligned}
\mathcal{P}(S_{\hat{z}} = +\hbar/2) &= |\langle S_{\hat{z}} = +\hbar/2 | S_{\hat{x}} = +\hbar/2 \rangle|^2 \\
&= \left| \frac{1}{\sqrt{2}} \langle S_{\hat{z}} = +\hbar/2 | \left(|S_{\hat{z}} = +\hbar/2\rangle + |S_{\hat{z}} = -\hbar/2\rangle \right) \right|^2 \\
&= \left| \frac{1}{\sqrt{2}} \left(\langle S_{\hat{z}} = +\hbar/2 | S_{\hat{z}} = +\hbar/2 \rangle + \langle S_{\hat{z}} = +\hbar/2 | S_{\hat{z}} = -\hbar/2 \rangle \right) \right|^2 \\
&= \frac{1}{2},
\end{aligned}$$

and

$$\begin{aligned}
\mathcal{P}(S_{\hat{z}} = -\hbar/2) &= |\langle S_{\hat{z}} = -\hbar/2 | S_{\hat{x}} = +\hbar/2 \rangle|^2 \\
&= \left| \frac{1}{\sqrt{2}} \langle S_{\hat{z}} = -\hbar/2 | \left(|S_{\hat{z}} = +\hbar/2\rangle - |S_{\hat{z}} = -\hbar/2\rangle \right) \right|^2 \\
&= \left| \frac{1}{\sqrt{2}} \left(\langle S_{\hat{z}} = -\hbar/2 | S_{\hat{z}} = +\hbar/2 \rangle - \langle S_{\hat{z}} = -\hbar/2 | S_{\hat{z}} = -\hbar/2 \rangle \right) \right|^2 \\
&= \frac{1}{2}.
\end{aligned}$$

