

General Relativity Formulary

1) Differential geometry

Metric connection (Christoffel symbols)

$$\Gamma_{\beta\gamma}^\alpha \equiv \frac{1}{2} g^{\alpha\delta} (\partial_\beta g_{\gamma\delta} + \partial_\gamma g_{\beta\delta} - \partial_\delta g_{\beta\gamma}).$$

Covariant derivative

$$\begin{aligned} V^\alpha_{;\beta} &= V^\alpha_{,\beta} + \Gamma_{\beta\gamma}^\alpha V^\gamma \\ V_{\alpha;\beta} &= V_{\alpha,\beta} - \Gamma_{\alpha\beta}^\gamma V_\gamma. \end{aligned}$$

Riemann curvature tensor

$$R^\alpha_{\beta\gamma\delta} \equiv \partial_\delta \Gamma_{\beta\gamma}^\alpha - \partial_\gamma \Gamma_{\beta\delta}^\alpha + \Gamma_{\beta\gamma}^\lambda \Gamma_{\delta\lambda}^\alpha - \Gamma_{\beta\delta}^\lambda \Gamma_{\lambda\gamma}^\alpha,$$

or, with the first index lowered

$$\begin{aligned} R_{\alpha\beta\gamma\delta} &= g_{\alpha\lambda} R^\lambda_{\beta\gamma\delta} \\ &= \frac{1}{2} (\partial_{\alpha\gamma}^2 g_{\beta\delta} + \partial_{\beta\delta}^2 g_{\alpha\gamma} - \partial_{\beta\gamma}^2 g_{\alpha\delta} - \partial_{\alpha\delta}^2 g_{\beta\gamma}) \\ &\quad + g_{\mu\nu} (\Gamma_{\alpha\gamma}^\mu \Gamma_{\beta\delta}^\nu - \Gamma_{\alpha\delta}^\mu \Gamma_{\beta\gamma}^\nu). \end{aligned}$$

Ricci tensor and Ricci scalar

$$R_{\alpha\beta} \equiv R^\gamma_{\alpha\gamma\beta}, \quad R \equiv g^{\alpha\beta} R_{\alpha\beta}.$$

Bianchi identities

$$R_{\alpha\beta\gamma\delta;\lambda} + R_{\alpha\beta\lambda\gamma;\delta} + R_{\alpha\beta\delta\lambda;\gamma} = 0.$$

2) General Relativity

Geodesic equations

$$\frac{d^2 x^\alpha}{dp^2} + \Gamma_{\beta\gamma}^\alpha \frac{dx^\beta}{dp} \frac{dx^\gamma}{dp} = 0,$$

where $p = \tau$ for a massive particle and $p = x^0$ for a massless particle.

$T_{\mu\nu}$ is the energy-momentum tensor given by

$$T_{\mu\nu} \equiv \frac{2}{\sqrt{-g}} \frac{\delta S_{\text{matter}}}{\delta g^{\mu\nu}}.$$

Variation of the determinant of the metric

$$\delta g = g g^{\mu\nu} \delta g_{\mu\nu} = -g g_{\mu\nu} \delta g^{\mu\nu}.$$

3) Energy-momentum tensor of point-particle

$$T^{\mu\nu} = \frac{m}{\sqrt{-g}} \int d\tau \frac{dz^\mu}{d\tau} \frac{dz^\nu}{d\tau} \delta^4(x - z(\tau)).$$

4) Static isotropic spacetime

$$ds^2 = B(r) dt^2 - A(r) dr^2 - r^2 (d\theta^2 + \sin^2\theta d\phi^2).$$

The corresponding Geodesic equations

$$\begin{aligned} \ddot{t} + \frac{B'}{B} \dot{t} \dot{r} &= 0 \\ \ddot{r} + \frac{1}{2} \frac{B'}{A} \dot{t}^2 + \frac{1}{2} \frac{A'}{A} \dot{r}^2 - \frac{r}{A} \dot{\theta}^2 - \frac{r \sin^2\theta}{A} \dot{\phi}^2 &= 0 \\ \ddot{\theta} + \frac{2}{r} \dot{r} \dot{\theta} - \sin \theta \cos \theta \dot{\phi}^2 &= 0 \\ \ddot{\phi} + \frac{2}{r} \dot{r} \dot{\phi} + 2 \cot \theta \dot{\theta} \dot{\phi} &= 0, \end{aligned}$$

4.1) Schwarzschild solution

$$B(r) = 1 - \frac{2GM}{r}, \quad A(r) = B^{-1}(r).$$

Motion in Schwarzschild solution

$$\begin{aligned} A(r) \left(\frac{dr}{d\phi} \right)^2 \frac{J^2}{E^2 r^4} + \frac{J^2}{E^2 r^2} - \frac{1}{B(r)} &= -\frac{m^2}{E^2} \\ d\tau^2 &= \frac{m^2 B^2(r)}{E^2} dt^2, \quad r^2 \frac{d\phi}{dt} = \frac{J}{E} B(r) \end{aligned}$$

Radial equation for a massive particle

$$\frac{1}{2} \left(\frac{dr}{d\tau} \right)^2 + \frac{1}{2} \left(1 - \frac{2GM}{r} \right) \left(1 + \frac{J^2/m^2}{r^2} \right) = \frac{E^2}{2m^2}$$

Radial equation for a massless particle ($d\lambda = B(r)dt$)

$$\frac{1}{2} \left(\frac{dr}{d\lambda} \right)^2 + \frac{1}{2} \left(1 - \frac{2GM}{r} \right) \frac{J^2}{E^2 r^2} = \frac{1}{2}$$

5) Numerical values

$$G = 6.67 \cdot 10^{-11} \text{ m}^3 \text{ kg}^{-1} \text{ s}^{-2}$$

$$M_\odot \simeq 2 \cdot 10^{30} \text{ kg}, \quad r_\odot \simeq 7 \cdot 10^5 \text{ km}$$

$$M_\oplus \simeq 6 \cdot 10^{24} \text{ kg}, \quad r_\oplus \simeq 6000 \text{ km}$$

$$M_{NS} \simeq 2M_\odot, \quad r_{NS} \simeq 12 \text{ km}$$