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### IGS: an IsoGeometric approach for Smoothing on surfaces

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#### Abstract

We propose an Isogeometric approach for smoothing on surfaces, namely estimating a function starting from noisy and discrete measurements. More precisely, we aim at estimating functions lying on a surface represented by NURBS, which are geometrical representations commonly used in industrial applications. The estimation is based on the minimization of a penalized least-square functional. The latter is equivalent to solve a 4th-order Partial Differential Equation (PDE). In this context, we use Isogeometric Analysis (IGA) for the numerical approximation of such surface PDE, leading to an IsoGeometric Smoothing (IGS) method for fitting data spatially distributed on a surface. Indeed, IGA facilitates encapsulating the exact geometrical representation of the surface in the analysis and also allows the use of at least globally  $C^1$ —continuous NURBS basis functions for which the 4th-order PDE can be solved using the standard Galerkin method. We show the performance of the proposed IGS method by means of numerical simulations and we apply it to the estimation of the pressure coefficient, and associated aerodynamic force on a winglet of the SOAR space shuttle.

**Keywords:** Functional Data Analysis; Isogeometric Analysis; Smoothing on Surfaces.

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#### 1 Introduction

The estimation of a function from a set of noisy data is a very common task, which is often tackled by minimization of a penalized least-square functional, where the penalty involves a differential operator, commonly based on second order derivatives, and occasionally on derivatives of a different order. Classical examples are offered by smoothing splines (see, e.g., Ramsay and Silverman (2005) and references therein) for estimating functions defined over real intervals, thin-plate splines (see, e.g., Wahba (1990), Wood (2006) and references therein), Soap film smoothing (Wood et al., 2008), splines over triangulations (see, e.g., Lai and Schumaker (2007)) for estimating functions defined over regions of  $\mathbb{R}^2$ , and spherical splines (see, e.g., Wahba (1981); Baramidze et al. (2006); Alfeld et al. (1996)) for estimating functions defined over spheres or spheres-like surfaces. In this context, one of the main challenges in minimizing such penalized least-square functional consists in determining a suitable finite dimensional space representing, at the discrete level, the infinite dimensional space to which the function belongs. In other words, the challenge is to find a finite dimensional problem which is tractable and whose solution is close to the solution in the infinite dimensional space. The same challenge arises when looking for the numerical solution of a PDE. This common goal has been recently exploited to develop statistical tools to deal with spatially distributed data. In this respect, Ramsay (2002) considered planar smoothing optimizing a penalized least-square functional with a regularization term involving the Laplacian, and used the Finite Element Method (FEM) (see, e.g., Quarteroni and Valli (1994)) to solve the estimation problem. Sangalli et al. (2013) generalized the method proposed by Ramsay (2002) to include space-varying covariates and to account for boundary conditions, while Wilhelm (2013) explored a generalized linear version of the method. Azzimonti et al. (2014a,b) further extended the model of Sangalli et al. (2013) to account for any elliptic differential penalization, not necessarily based on the Laplacian operator. Ettinger et al. (2012) and Dassi et al. (2015) dealt with the case of functions defined over two dimensional Riemannian manifolds, by considering a regularizing term based on the Laplace-Beltrami operator and exploited a conformal parametrization of the manifold to solve an equivalent estimation problem on  $\mathbb{R}^2$ . Duchamp and Stuetzle (2003) also explored smoothing on complex surfaces using a penalization based on the Laplace-Beltrami operator. A common feature of all these contributions is the use of FEM to estimate the underlying function corresponding to the observed data.

In this paper, we consider the estimation of functions defined over surfaces in  $\mathbb{R}^3$  starting from a discrete set of noisy observed data in points distributed on such surfaces. More precisely, we refer to sufficiently smooth functions defined over surfaces that can be represented by NURBS (Non-Uniform Rational Basis Splines). Indeed, NURBS are commonly used in Computer Aided Design (CAD) to represent most of the geometries of Engineering and industrial interest. From a more general point of view, the proposed model is a particular case of a Generalized Additive Model (GAM) (Hastie and Tibshirani, 1990) where the smooth component is defined on a surface. Hence, most of the theoretical results of GAMs can be directly applied to this model, e.g., for the quantification of uncertainty.

IGA has been first introduced by Hughes et al. (2005) with the main idea of using the same basis functions to represent the geometry and then to approximate the solution of the PDEs defined in such computational domains. This facilitates encapsulating the exact geometrical representation when performing the analysis of PDEs defined in the computational domain; see Cottrell et al. (2009). In this respect, the isogeometric paradigm facilitates the use of an exact representation of the surface, while most of the current methodologies, as FEM, generally only handle its ap-

proximation. This may induce an error on the solution due to the approximation of the geometry and may require complex meshing procedures. As mentioned, to estimate the function over the surface, starting from its discrete and noisy measurements, we minimize a least-square functional where the penalty involves the Laplace-Beltrami operator associated to the surface, analogously to Ettinger et al. (2012), Duchamp and Stuetzle (2003) and Wahba (1981). The estimation problem is tackled directly on the surface and using IGA to numerically solve the associated surface PDE. For this reason, we name the resulting method IsoGeometric Smoothing (IGS). High-order PDEs can be straightforwardly solved by NURBS-based IGA; indeed, globally  $C^k$ —continuous NURBS basis functions can be easily defined on surfaces for some  $k = 0, \ldots, p - 1$ , where p is the polynomial degree. This allows to use standard Galerkin method for numerically solving the PDE. In this respect, Dedè and Quarteroni (2015) studied the use of IGA for 2nd-order PDEs defined on surfaces, Tagliabue et al. (2014) analysed IGA for high-order PDEs, while Bartezzaghi et al. (2015) for high-order PDEs defined on surfaces. All these works showed the efficiency and accuracy of IGA for the spatial approximation of surface PDEs.

In this paper, we introduce the geometrical background to define the application framework of IGS. Then, we show how one can minimize a least-square functional involving a penalization term and hence estimate a function starting from a set of discrete and noisy measurements. We assess and compare IGS to the Thin Plate Splines (TPS) (Duchon, 1977; Wahba, 1990) by means of numerical simulations. Finally, as industrial application, we estimate the pressure coefficient and the aerodynamic force acting on the inboard winglet of the space shuttle SOAR, designed by S3, Swiss Space Systems Holding SA. S3 is a Swiss company currently developping, manufacturing, certifying, and operating a launch system for small satellites of weight inferior to 250 kg. A geometry of the SOAR suborbital shuttle for a preliminary study of aerodynamic forces is shown in Figure 1. In this application, the data are pointwise measurements of the pressure coefficient and the associated quantity of interest is the aerodynamic force. We propose a method to estimate functions defined over a NURBS surface starting from scattered noisy observations.

This paper is organized as follows. In Section 2, we briefly recall B-splines and NURBS. In Section 3, we introduce some analytical tools and the IGS method. We show results for two numerical simulations in Section 4: the first one allows to compare IGS to TPS, while the second one corresponds to a complex NURBS surface, for which planar smoothers are not suited. Finally, in Section 5, we show the results of the estimation of the pressure coefficient and aerodynamic force over the winglet of the SOAR space shuttle. Conclusions follow.



Figure 1: The SOAR shuttle highlighting the inboard winglet represented by a single NURBS patch. [Courtesy of S3, Swiss Space Systems Holding SA].

#### 2 B-Splines and NURBS

NURBS are widely used in CAD for geometrical representation of surfaces (Piegl and Tiller, 1997). We start by defining an open knot vector of degree p as a vector of values  $\Xi = \{\xi_1, \ldots, \xi_{n+p+1}\}$ , with knots  $\xi_1 \leq \ldots \leq \xi_{n+p+1}$ , where the first and the last knots are repeated p+1 times. The interior knots can be repeated at most p times and, if a knot is repeated m times, we say that its multiplicity is m. The B-Spline basis functions  $N_{i,p}$  are defined recursively using the Cox-De Boor formula. We summarize some properties of B-Spline basis functions  $N_{i,p}$  of index i and degree p, for  $i = 1, \ldots, n$ .

- The basis function  $N_{i,p}$  possesses p-m continuous derivatives across a knot of multiplicity  $1 \le m \le p$  and is  $C^{\infty}$  continuous between the knots.
- The support of the basis function  $N_{i,p}$  is compact and contained in p+1 knots spans  $[\xi_i, \xi_{i+p+1}]$ .
- The basis functions are pointwise nonnegative, i.e.  $N_{i,p}(\xi) \geq 0$ , for all  $i = 1, \ldots, n$ .
- They form a partition of the unity, that is  $\sum_{i=1}^{n} N_{i,p}(\xi) = 1$ , for all  $\xi \in [\xi_1, \xi_n]$ .

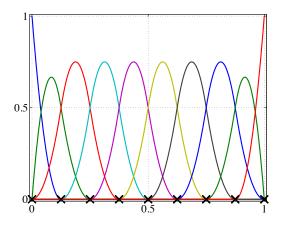
Starting from the basis function  $N_{i,p}$ , the NURBS basis functions  $R_{i,p}(\xi)$  are defined as projective transformations of B-Spline basis functions. Let  $w_1, \ldots, w_n > 0$  be positive weights, then, NURBS basis functions are defined as:

$$R_{i,p}(\xi) = \frac{w_i}{\sum_{j=1}^n w_j N_{j,p}(\xi)} N_{i,p}(\xi), \quad \forall i = 1, \dots, n.$$

A NURBS curve  $C(\xi) \in \mathbb{R}^d$ , with d = 2, 3 and control points  $\mathbf{B}_1, \dots, \mathbf{B}_n \in \mathbb{R}^d$  is defined as:

$$C(\xi) = \sum_{i=1}^{n} R_{i,p}(\xi) \mathbf{B}_{i}.$$

To define surfaces with NURBS, we resort to the tensor product rule. Given two open knot vectors  $\Xi = \{\xi_1, \ldots, \xi_{n+p+1}\}$  and  $\mathscr{H} = \{\eta_1, \ldots, \eta_{m+q+1}\}, Q_{i,p}(\xi), i = 1, \ldots, n, M_{j,q}, i = 1, \ldots, m$ 



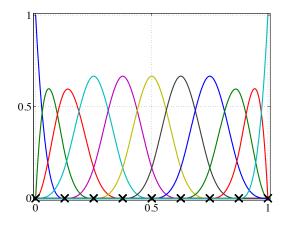


Figure 2: B-Spline basis functions for different knots vectors. Basis functions of degree p=2 and globally  $C^1$ -continuous (left) and of degree p=3 and globally  $C^2$ -continuous (right). The  $\times$  indicates the knots.

the corresponding univariate basis functions and control points  $\mathbf{B}_{i,j} \in \mathbb{R}^d$ , for i = 1, ..., n, j = 1, ..., m, a NURBS surface in  $\mathbb{R}^d$  is defined as:

$$S: (\xi_1, \xi_n) \times (\eta_1, \eta_m) \to \mathbb{R}^d, \quad S(\xi, \eta) = \sum_{i=1}^n \sum_{j=1}^m Q_{i,p}(\xi) M_{j,q}(\eta) \mathbf{B}_{i,j}.$$
 (2.1)

By redefining the bivariate NURBS basis functions as  $R_{i,p}(\xi,\eta) = Q_{j,p}(\xi)M_{k,q}(\eta)$ , after having introduced a suitable reordering of the indexes and set q = p, we rewrite (2.1) as:

$$S: \Omega \to \mathbb{R}^d, \quad S(\xi, \eta) = \sum_{i=1}^n R_{i,p}(\xi, \eta) \mathbf{B}_i,$$

where  $\Omega = (\xi_1, \xi_m) \times (\eta_1, \eta_m)$ . For an exhaustive description of NURBS, we refer the reader to Piegl and Tiller (1997).

We remember that in IGA, one uses the same basis functions to represent the surface and to approximate the solution of the PDE defined in such computational domain. In general, it is possible to enrich the NURBS basis without changing the geometry, i.e. by preserving the geometrical mapping, with the goal of obtaining a more accurate solution of the PDEs. In this respect, the h-refinement indicates a refinement of the mesh (knot insertion), while the p-refinement an order elevation of the polynomial degree of the basis functions, similarly to FEM. In addition, the so-called k- refinement is peculiar of NURBS and refers to a consecutive order elevation and a knot insertion which allow to increase the polynomial degree and continuity of basis functions. All the refinement procedures are discussed in details in Cottrell et al. (2009), while for an introduction to NURBS in the context of IGA, we refer the interested reader to Hughes et al. (2005).

#### 3 The IGS method

We describe the mathematical framework of the IGS method. First, we introduce the surface differential operators and then we introduce the IGS method for smoothing functions on surfaces.

#### 3.1 Geometrical framework

Following Dedè and Quarteroni (2015), let  $\Omega \subset \mathbb{R}^2$  be an open and bounded parametric domain of finite measure with respect to the topology of  $\mathbb{R}^2$ . Then, let  $\Sigma \subset \mathbb{R}^3$  be a compact, connected, and oriented surface, defined by a NURBS geometrical mapping  $X : \Omega \to \Sigma$  such that:

$$X: \Omega \subset \mathbb{R}^2 \to \Sigma \subset \mathbb{R}^3, \quad \mathbf{s} = (s_1, s_2) \mapsto \mathbf{x} = (x_1, x_2, x_3).$$
 (3.1)

We assume that X is sufficiently smooth, e.g.  $X \in C^1(\Omega)$ . Then, we define the Jacobian of the mapping X, denoted by  $\nabla X$  as:

$$\nabla X: \Omega \to \mathbb{R}^{3 \times 2}, \quad \mathbf{s} \mapsto \nabla X(\mathbf{s}), \quad (\nabla X)_{i,j}(\mathbf{s}) = \frac{\partial X_i}{\partial \mathbf{s}_i}(\mathbf{s}), \quad i = 1, 2, 3, \ j = 1, 2.$$

We denote by  $(\nabla X)_i$  the *i*-th column of the matrix  $\nabla X$ . The first fundamental form of the mapping X is given by:

$$\mathbf{G}: \Omega \to \mathbb{R}^{2 \times 2}, \quad \mathbf{s} \mapsto \mathbf{G}(\mathbf{s}), \quad \mathbf{G}(\mathbf{s}) = (\nabla X(\mathbf{s}))^T \nabla X(\mathbf{s})$$

and we denote with  $g(\mathbf{s})$  the square root of the determinant of the Jacobian of the mapping X:

$$g: \Omega \to \mathbb{R}, \quad \mathbf{s} \mapsto g(\mathbf{s}), \quad g(\mathbf{s}) = \sqrt{\det(\mathbf{G}(\mathbf{s}))}.$$

The Jacobian of the mapping X is assumed to be invertible almost everywhere in  $\Omega$ .

#### 3.2 Differential operators

Let  $\phi \in C^2(\Sigma)$  be a smooth function defined over the surface  $\Sigma$ , i.e.:

$$\phi: \Sigma \subset \mathbb{R}^3 \to \mathbb{R}, \quad \mathbf{x} \mapsto \phi(\mathbf{x}), \quad \phi \in C^2(\Sigma),$$

such that we can define the differential operators on the surface  $\Sigma$ . Then, the projection operator  $\mathbf{P}(\mathbf{x})$  on the tangential plane to the surface at the point  $\mathbf{x} \in \Sigma$  is given by:

$$\mathbf{P}(\mathbf{x}) = (\mathbf{I} - \mathbf{n}_{\Sigma} \otimes \mathbf{n}_{\Sigma})(\mathbf{x}) = \mathbf{I} - \mathbf{n}_{\Sigma}(\mathbf{x})\mathbf{n}_{\Sigma}(\mathbf{x})^{T}, \quad \forall \mathbf{x} \in \Sigma,$$

where  $\mathbf{n}_{\Sigma}(\mathbf{x})$  is the unit normal vector to the surface at the point  $\mathbf{x}$  and  $\mathbf{I}$  is the identity matrix<sup>1</sup>. Then, the surface gradient operator, denoted by  $\nabla_{\Sigma}$ , is defined as the projection of the gradient of a function extended in a tubular region containing  $\Sigma$ , i.e.:

$$\nabla_{\Sigma}\phi(\mathbf{x}) := \mathbf{P}(\mathbf{x})\nabla\phi(\mathbf{x}) = \nabla\phi(\mathbf{x}) - (\mathbf{n}_{\Sigma}(\mathbf{x})^{T}\nabla\phi(\mathbf{x}))\mathbf{n}_{\Sigma}(\mathbf{x}).$$

The Laplace-Beltrami operator, which is the surface Laplacian operator, is expressed as:

$$\Delta_{\Sigma}\phi(\mathbf{x}) = \nabla_{\Sigma} \cdot (\nabla_{\Sigma}\phi(\mathbf{x})) = \operatorname{trace}[\mathbf{P}(\mathbf{x}) \nabla^2 \phi \, \mathbf{P}(\mathbf{x})], \quad \forall \mathbf{x} \in \Sigma,$$

where  $\nabla_{\Sigma} \cdot \mathbf{v}(\mathbf{x}) = \operatorname{trace}(\nabla_{\Sigma} \mathbf{v}(\mathbf{x}))$ , for all  $\mathbf{v} \in C^1(\Sigma)$ , is the surface divergence and  $(\nabla^2 \phi)_{ij}(\mathbf{x}) = \frac{\partial^2 \phi}{\partial x_i \partial x_j}(\mathbf{x})$  is the Hessian matrix of  $\phi$ . These operators can be rewritten in the parametric domain  $\Omega$  using the geometrical mapping (3.1) as:

$$\nabla_{\Sigma}\phi(\mathbf{x}) = \nabla X(\mathbf{s})\mathbf{G}^{-1}(\mathbf{s})\nabla(\phi \circ X)(\mathbf{s}) \text{ and } \Delta_{\Sigma}\phi(\mathbf{x}) = \frac{1}{g(\mathbf{s})}\nabla \cdot \left[g(\mathbf{s})\mathbf{G}^{-1}(\mathbf{s})\nabla(\phi \circ X)(\mathbf{s})\right], \quad (3.2)$$

where  $\mathbf{s} = X^{-1}(\mathbf{x})$ .

We compute  $\mathbf{n}_{\Sigma}(\mathbf{x})$  as  $\mathbf{n}_{\Sigma}(\mathbf{x}) = \frac{[\nabla X]_1(\mathbf{s}) \times [\nabla X]_2(\mathbf{s})}{\|[\nabla X]_1(\mathbf{s}) \times [\nabla X]_2(\mathbf{s})\|_2}$ , where  $\mathbf{s} = X^{-1}(\mathbf{x})$ .

#### 3.3 Mathematical model

Let us consider N points  $\mathbf{p}_1, \dots, \mathbf{p}_N$  located on  $\Sigma$  for which the observed values are  $y_1, \dots, y_N$  respectively. We assume that:

$$y_i = f(\mathbf{p}_i) + \varepsilon_i, \quad \forall \ i = 1, \dots, N,$$
 (3.3)

where f is a sufficiently smooth function defined on  $\Sigma$  that we aim at estimating and  $\varepsilon_i$  are independent observational errors of zero mean and constant variance  $\sigma^2$ . Given a positive smoothing parameter  $\lambda$ , we aim at minimizing the following parameter dependent functional:

$$J_{\lambda}(v) = \sum_{i=1}^{N} (y_i - v(\mathbf{p}_i))^2 + \lambda \int_{\Sigma} (\Delta_{\Sigma} v)^2 d\Sigma = \|\mathbf{y} - \mathbf{v}_N\|_2^2 + \lambda \int_{\Sigma} (\Delta_{\Sigma} v)^2 d\Sigma, \tag{3.4}$$

where  $\mathbf{y} = (y_1, \dots, y_N)^T$  and  $\mathbf{v}_N = (v(\mathbf{p}_1), \dots, v(\mathbf{p}_N))^T$  is the vector of evaluation of the general function v at the points  $\mathbf{p}_1, \dots, \mathbf{p}_N$ . This functional is the same considered by Ettinger et al. (2012) and by Duchamp and Stuetzle (2003) that uses a finite element representation, and by Wahba (1981), that considers functions on spheres and proposes spherical splines. The functional (3.4) involves the surface Laplace-Beltrami operator, which is, roughly speaking, a measure of the curvature of the function related to the surface. The use of the Laplace-Beltrami operator ensures that the regularization is invariant to rigid transformations of the coordinate system, which is desirable both for theoretical and practical reasons. If the observational errors are normally distributed, the functional (3.4) can be viewed as a negative rescaled Gaussian penalized log-likelihood. Hence, in such case, minimizing (3.4) is equivalent to maximizing a penalized log-likelihood. Then, for a given positive parameter  $\lambda$ , the estimation problem is:

find 
$$\hat{f} \in \mathcal{F}: J_{\lambda}(\hat{f}) \leq J_{\lambda}(v), \quad \forall v \in \mathcal{F},$$
 (3.5)

where  $\mathcal{F}$  is a suitable functional space to be defined to ensure the well-posedness of the problem and  $\hat{f}$  is the estimate of f in the functional space  $\mathcal{F}$ , which should lie at least in  $H^2(\Sigma)$ , i.e.  $\mathcal{F} \subset H^2(\Sigma)$ . Indeed, since  $\hat{f} \in H^2(\Sigma) \subset C^0(\Sigma)$ , the evaluation of  $\hat{f}$  at the points  $\mathbf{p}_1, \ldots, \mathbf{p}_N$  is well defined. Since the problem (3.5) will be later associated to a PDE, some essential boundary conditions (Brezis, 1999) should be specified in relation with the choice of  $\mathcal{F}$ . In the simpler context of functions defined over planar domains, and with the penalizing terms involving linear second order elliptic operators, the well-posedness of problem (3.4) is extensively discussed in Azzimonti et al. (2014a) under different kind of boundary conditions (Dirichlet, Neumann, or Robin (see, e.g., Brezis (1999)). In the following, we prove the well-posedness of the estimation problem (3.5) in the particular case of homogeneous boundary conditions using the Lax-Milgram theorem, following Ramsay (2002) and Sangalli et al. (2013). Let  $H_0^2(\Sigma)$  be defined as:

$$H_0^2(\Sigma) := \left\{ v \in H^2(\Sigma) : \nabla_{\Sigma} v \cdot \mathbf{n} = 0 \text{ and } v = 0 \text{ on } \partial \Sigma \right\}, \tag{3.6}$$

where **n** denotes the outward directed unit vector normal to  $\partial \Sigma$ , the boundary of  $\Sigma$ . Then, one can characterize the solution of the minimization problem (3.5) and ensure the existence and the uniqueness of the solution in the case where  $\hat{f}$  is assumed to lie in  $\mathcal{F} = H_0^2(\Sigma)$ . With this aim, we recall the Lax-Milgram Theorem (see, e.g., Quarteroni and Valli (1994)):

**Theorem 3.1** (Lax-Milgram). Let  $\mathcal{F}$  be a Hilbert space,  $G(\cdot, \cdot): \mathcal{F} \times \mathcal{F} \to \mathbb{R}$  a continuous and coercive bilinear form and  $F: \mathcal{F} \to \mathbb{R}$  a linear and continuous functional. Then, there exists a unique solution of the following problem:

find 
$$u \in \mathcal{F} : G(v, u) = F(v), \forall v \in \mathcal{F}.$$

Moreover, if  $G(\cdot, \cdot)$  is symmetric, then  $u \in \mathcal{F}$  is the unique minimizer in  $\mathcal{F}$  of the functional  $J: \mathcal{F} \to \mathbb{R}$ , defined as

$$J(v) = G(v, v) - 2F(v).$$

**Lemma 3.1.** Let  $v, u \in \mathcal{F} = H_0^2(\Sigma)$  be two functions and  $\lambda > 0$  a positive smoothing parameter. Let  $G(\cdot, \cdot)$  and  $F(\cdot)$  be defined as:

$$G(v,u) = \langle \mathbf{v}_N, \mathbf{u}_N \rangle + \lambda \int_{\Sigma} (\Delta_{\Sigma} v)(\Delta_{\Sigma} u) \ d\Sigma, \qquad and \qquad F(v) = \langle \mathbf{y}, \mathbf{v}_N \rangle, \tag{3.7}$$

where  $\mathbf{v}_N = (v(\mathbf{p}_1), \dots, v(\mathbf{p}_N))^T$  and  $\mathbf{u}_N = (u(\mathbf{p}_1), \dots, u(\mathbf{p}_N))^T$  for some distinct points  $\mathbf{p}_1, \dots, \mathbf{p}_N$  on  $\Sigma$ . Then, the bilinear form  $G(\cdot, \cdot)$  is coercive, continuous and symmetric and  $F(\cdot)$  is linear and continuous in  $H_0^2(\Sigma)$ .

*Proof.* First, we note that  $|\cdot|_{H^2(\Sigma)}$ , defined as  $|v|_{H^2(\Sigma)} = \int_{\Sigma} (\Delta_{\Sigma} v)^2 d\Sigma$ , is equivalent to the norm  $\|\cdot\|_{H^2(\Sigma)}$  in  $H_0^2(\Sigma)$ . This implies that there exists a constant  $C_{0,\Omega} > 0$  such that  $|v|_{H^2(\Sigma)} \ge C_{0,\Omega} \|v\|_{H^2(\Sigma)}$ ,  $\forall v \in H_0^2(\Sigma)$  (see, e.g., Quarteroni (2015)). Then, we have:

$$G(v,v) = \langle \mathbf{v}_N, \mathbf{v}_N \rangle + \lambda \int_{\Sigma} (\Delta_{\Sigma} v)^2 d\Sigma = \|\mathbf{v}_N\|_2^2 + \lambda |v|_{H^2(\Sigma)}^2 \ge \lambda C_{0,\Omega} \|v\|_{H^2(\Sigma)}^2, \quad \forall v \in H_0^2(\Sigma),$$

and so  $G(\cdot, \cdot)$  is coercive. We now show the continuity of  $G(\cdot, \cdot)$ . Since  $H^2(\Sigma) \subset C^0(\Sigma)$ , there exists a constant  $C_{1,\Omega,N}$  such that  $\|\mathbf{v}_N\|_2 \leq C_{1,\Omega,N} \|v\|_{H^2(\Sigma)}$ ,  $\forall v \in H^2(\Sigma)$ . We have:

$$G(v, u) = \langle \mathbf{v}_{N}, \mathbf{u}_{N} \rangle + \lambda \int_{\Sigma} (\Delta_{\Sigma} v)(\Delta_{\Sigma} u) \leq \|\mathbf{v}_{N}\|_{2} \|\mathbf{u}_{N}\|_{2} + \lambda |v|_{H^{2}(\Sigma)} |u|_{H^{2}(\Sigma)}$$

$$\leq C_{1,\Omega,N}^{2} \|v\|_{H^{2}(\Sigma)} \|u\|_{H^{2}(\Sigma)} + \lambda |v|_{H^{2}(\Sigma)} |u|_{H^{2}(\Sigma)}$$

$$\leq \max\{C_{1,\Omega,N}^{2}, \lambda\} \|v\|_{H^{2}(\Sigma)} \|u\|_{H^{2}(\Sigma)}, \quad \forall v, u \in H_{0}^{2}(\Sigma).$$

Then, the bilinear form  $G(\cdot, \cdot)$  is also continuous and its symmetry is obvious. Finally, we have:

$$|F(v)| = |\langle \mathbf{y}, \mathbf{v}_N \rangle| \le ||\mathbf{y}||_2 ||\mathbf{v}_N||_2 \le C_{1,\Omega,N} ||\mathbf{y}||_2 ||v||_{H^2(\Sigma)}, \quad \forall v \in H_0^2(\Sigma),$$

which proves the continuity of the linear form F and concludes the proof.

**Proposition 3.1.** Let  $\mathcal{F} = H_0^2(\Sigma)$  and  $\lambda > 0$ . Then, the solution of problem (3.5) exists and is unique. Moreover, problem (3.5) is equivalent to:

find 
$$\hat{f} \in \mathcal{F} : \langle \mathbf{v}_N, \mathbf{\hat{f}}_N \rangle + \lambda \int_{\Sigma} \Delta_{\Sigma} v \, \Delta_{\Sigma} \hat{f} \, d\Sigma = \langle \mathbf{y}, \mathbf{v}_N \rangle, \quad \forall v \in \mathcal{F}.$$
 (3.8)

*Proof.* The functional  $J_{\lambda}(v)$  (3.4) can be rewritten as:

$$J_{\lambda}(v) = \|\mathbf{y} - \mathbf{v}_{N}\|_{2}^{2} + \lambda \int_{\Sigma} (\Delta_{\Sigma} f)^{2} d\Sigma = \|\mathbf{y}\|_{2}^{2} - 2 \langle \mathbf{y}, \mathbf{v}_{N} \rangle + \|\mathbf{v}_{N}\|_{2}^{2} + \lambda \int_{\Sigma} (\Delta_{\Sigma} v)^{2} d\Sigma.$$

By defining  $\tilde{J}_{\lambda}(v)$  as

$$\tilde{J}_{\lambda}(v) = \|\mathbf{v}_N\|_2^2 + \lambda \int_{\Sigma} (\Delta_{\Sigma} v)^2 d\Sigma - 2 \langle \mathbf{y}, \mathbf{v}_N \rangle,$$

we have that

$$\underset{v \in \mathcal{F}}{\operatorname{arg \ min}} \ J_{\lambda}(v) \equiv \underset{v \in \mathcal{F}}{\operatorname{arg \ min}} \ \tilde{J}_{\lambda}(v).$$

From the definitions of G(v, u) and F(v) of (3.7), the functional  $\tilde{J}_{\lambda}(v)$  can be written as:

$$\tilde{J}_{\lambda}(v) = G(v, v) - 2F(v).$$

Thanks to the Lax-Milgram Theorem, it is then sufficient to use Lemma 3.1 to establish the well-posedness of problem (3.5). Moreover, since  $G(\cdot, \cdot)$  is symmetric, problem (3.5) is equivalent to problem (3.8).

Setting a priori the essential boundary conditions  $\nabla_{\Sigma} \hat{f} \cdot \mathbf{n} = 0$  and  $\hat{f} = 0$  on  $\partial \Sigma$  following (3.6) may be an inadequate choice in several applications, especially when the behaviour of the function f at the boundaries is not known a priori. In such cases, it may be more convenient to consider instead natural boundary conditions, that is:

$$\begin{cases} \nabla_{\Sigma}(\Delta_{\Sigma} \hat{f}) \cdot \mathbf{n} = 0 & \text{on } \partial \Sigma, \\ \Delta_{\Sigma} \hat{f} = 0 & \text{on } \partial \Sigma. \end{cases}$$
(3.9)

In the case that the boundary conditions (3.9) are set, we are unable to show the well-posedness of problem (3.8) with  $\mathcal{F} = H^2(\Sigma)$ . Nevertheless, numerical experience indicates that it still yields a numerically stable problem. We can also note that the usual planar smoothers, such as TPS, also implicitly use natural boundary conditions.

#### 3.4 Numerical approximation: IGS

Let  $\hat{f}^h$  be a finite dimensional approximation of  $\hat{f}$  obtained by means of IGS. Let  $\{\psi_1, \ldots, \psi_{N^h}\}$  be a basis of a discrete function space  $\mathcal{F}^h \subset \mathcal{F} \subseteq H^2(\Sigma)$  of dimension  $N^h$ . In the finite dimensional space  $\mathcal{F}^h$ , problem (3.8) reads:

find 
$$\hat{f}^h \in \mathcal{F}^h : \left\langle \mathbf{v}_N^h, \hat{\mathbf{f}}_N^h \right\rangle + \lambda \int_{\Sigma} \Delta_{\Sigma} v^h \, \Delta_{\Sigma} \hat{f}^h \, d\Sigma = \left\langle \mathbf{y}, \mathbf{v}_N^h \right\rangle, \quad \forall v^h \in \mathcal{F}^h, \quad (3.10)$$

where  $\mathbf{v}_N^h := (v^h(\mathbf{p}_1), \dots, v^h(\mathbf{p}_N))^T$  and  $\hat{\mathbf{f}}_N^h := (\hat{f}^h(\mathbf{p}_1), \dots, \hat{f}^h(\mathbf{p}_N))^T$ . Since  $\mathcal{F}^h$  is finite dimensional, problem (3.10) is equivalent to:

find 
$$\hat{f}^h$$
:  $\langle \boldsymbol{\psi}_{iN}, \hat{\mathbf{f}}_N^h \rangle + \lambda \int_{\Sigma} \Delta_{\Sigma} \psi_i \, \Delta_{\Sigma} \hat{f}^h \, d\Sigma = \langle \mathbf{y}, \boldsymbol{\psi}_{iN} \rangle, \quad \forall i = 1 \dots, N^h,$  (3.11)

where  $\psi_{iN} := (\psi_i(\mathbf{p}_1), \dots, \psi_i(\mathbf{p}_N))^T$ . Let us define the  $N^h \times N^h$  matrix  $\mathbf{R}$  as  $(\mathbf{R})_{ij} = \int_{\Sigma} \Delta_{\Sigma} \psi_i \, \Delta_{\Sigma} \psi_j \, d\Sigma$ , and the  $N \times N^h$  matrix  $\mathbf{\Psi}$  as  $(\mathbf{\Psi})_{ij} = \psi_j(\mathbf{p}_i)$ . Since  $\hat{f}^h$  belongs to  $\mathcal{F}^h$ , it can be written as a linear combination of the basis functions:

$$\hat{f}^h(\mathbf{x}) = \sum_{i=1}^{N^h} \hat{f}_i \ \psi_i(\mathbf{x}), \quad \forall \mathbf{x} \in \Sigma,$$

or compactly as  $\hat{f}^h(\mathbf{x}) = \boldsymbol{\psi}^T(\mathbf{x}) \hat{\mathbf{f}}^h$ , where  $\hat{\mathbf{f}}^h := (\hat{f}_1, \dots, \hat{f}_{N^h})^T$  and  $\boldsymbol{\psi}(\mathbf{x}) = (\psi_1(\mathbf{x}), \dots, \psi_{N^h}(\mathbf{x}))^T$ . Then, we have:

$$\mathbf{\hat{y}} := \mathbf{\hat{f}}_N^h = \mathbf{\Psi}\mathbf{\hat{f}}^h.$$

Problem (3.11) in matrix form reads as:

find 
$$\hat{\mathbf{f}}^h \in \mathbb{R}^{N^h}$$
:  $\mathbf{A}\hat{\mathbf{f}}^h = \mathbf{\Psi}^T \mathbf{v}$ .

where  $\mathbf{A} := (\mathbf{\Psi}^T \mathbf{\Psi} + \lambda \mathbf{R})$ . Then, the explicit form of  $\mathbf{\hat{f}}^h$  is given by:

$$\mathbf{\hat{f}}^h = \mathbf{A}^{-1} \mathbf{\Psi}^T \mathbf{y} = (\mathbf{\Psi}^T \mathbf{\Psi} + \lambda \mathbf{R})^{-1} \mathbf{\Psi}^T \mathbf{y}$$
(3.12)

We see from (3.12) that the estimator  $\hat{\mathbf{f}}^h$  has the typical form of a penalized least-square estimator. Since  $\hat{\mathbf{y}} = \mathbf{\Psi}\hat{\mathbf{f}}^h$ , we finally get the evaluation of the function  $\hat{f}^h$  in the points  $\{\mathbf{p}_1, \dots, \mathbf{p}_N\}$  as:

$$\hat{\mathbf{y}} = \mathbf{\Psi} (\mathbf{\Psi}^T \mathbf{\Psi} + \lambda \mathbf{R})^{-1} \mathbf{\Psi}^T \mathbf{y}. \tag{3.13}$$

The smoothing matrix, that maps the observed data values  $\mathbf{y}$  in the fitted data values  $\hat{\mathbf{y}}$ , is given by:

$$\mathbf{S} = \mathbf{\Psi}(\mathbf{\Psi}^T \mathbf{\Psi} + \lambda \mathbf{R})^{-1} \mathbf{\Psi}^T. \tag{3.14}$$

The trace of S is a measure of the equivalent degrees of freedom of the estimator (see Buja et al. (1989)).

We use IGA to solve the minimization problem (3.5), for which, we define:

$$\mathcal{F}^h = \operatorname{span}\left\{R_{i,p}(\xi,\eta), i = 1,\dots,N^h\right\},\,$$

where  $R_{i,p}$  are the NURBS basis functions used to build  $\Sigma$ , eventually after the application of some h-, p- or k-refinement procedure, as described in Section 2. NURBS allow to define basis functions which are globally  $C^1-$ continuous on  $\Sigma$ . As consequence, one can approximate problem (3.5) with the standard Galerkin method, since  $\mathcal{F}^h \subset H^2(\Sigma)$ ; see Bartezzaghi et al. (2015) and Tagliabue et al. (2014). In this manner, we obtain a method, which we name IGS, allowing to perform smoothing on surfaces by means of NURBS-based IGA. This also allows encapsulating the original description of the geometry of the surface in the analysis.

The smoothing parameter  $\lambda$  may be chosen by minimization of a generalized cross-validation criterion (GCV), defined as:

$$GCV(\lambda) = \frac{N}{[N - \operatorname{trace}(\mathbf{S}(\lambda))]^2} \|\hat{\mathbf{y}}(\lambda) - \mathbf{y}\|^2;$$
(3.15)

see Craven and Wahba (1978). Here, we use trace  $(\mathbf{S}(\lambda))$  as a measure of the equivalent degrees of freedom (EDF) of the model (Buja et al., 1989). In order to solve the optimization problem corresponding to GCV minimization, we use a BFGS quasi-Newton method (see, e.g., Nocedal and Wright (1999)) with a sufficiently small tolerance.

#### 3.5 Distributional properties and quantification of uncertainty

For a given smoothing parameter  $\lambda$ , the estimate  $\hat{\mathbf{f}}^h$  is a linear transformation of the observations  $\mathbf{y}$ , as shown in (3.12). Moreover, we have  $\mathbb{E}[\mathbf{y}] = \mathbf{f}_N = (f(\mathbf{p}_1), \dots, f(\mathbf{p}_n))^T$  and  $\text{var}(\mathbf{y}) = \sigma^2 \mathbf{I}$ . Then, from (3.12), we get:

$$\mathbb{E}[\mathbf{\hat{f}}^h] = \mathbb{E}[(\mathbf{\Psi}^T \mathbf{\Psi} + \lambda \mathbf{R})^{-1} \mathbf{\Psi}^T \mathbf{y}] = \mathbf{A}^{-1} \mathbf{\Psi}^T \mathbf{f}_N.$$

Similarly, we can directly express the variance as:

$$\operatorname{var}(\mathbf{\hat{f}}^h) = \operatorname{var}(\mathbf{A}^{-1}\mathbf{\Psi}^T\mathbf{y}) = \sigma^2 \mathbf{A}^{-1}\mathbf{\Psi}^T\mathbf{\Psi}\mathbf{A}^{-1}, \tag{3.16}$$

where we used the fact that the matrix **A** is symmetric and hence  $\mathbf{A}^{-T} = \mathbf{A}^{-1}$ . In particular, under the assumption of normality of the errors, we have:

$$\hat{\mathbf{f}}^h \sim \mathcal{N}(\mathbf{A}^{-1}\mathbf{\Psi}^T\mathbf{f}_N, \sigma^2\mathbf{A}^{-1}\mathbf{\Psi}^T\mathbf{\Psi}\mathbf{A}^{-1}).$$

Then, we can also express the expectation and the variance of the fitted values  $\hat{\mathbf{y}}$  explicitly as:

$$\mathbb{E}\left[\hat{\mathbf{y}}\right] = \mathbf{S} \; \mathbf{f}_N,$$

and:

$$var(\mathbf{\hat{y}}) = \sigma^2 \mathbf{S} \mathbf{S}^T = \sigma^2 \ \mathbf{S}^2,$$

respectively, since the smoothing matrix **S** is symmetric. In practice, the error variance  $\sigma^2$  must be estimated from the data. Following Hastie and Tibshirani (1990) we can estimate  $\sigma^2$  by:

$$\hat{\sigma}^2 = \frac{\|\hat{\mathbf{y}} - \mathbf{y}\|^2}{N - \text{trace}(\mathbf{S})}.$$
(3.17)

Given an additional point  $\mathbf{p}_{N+1}$  on  $\Sigma$ , the predicted value of the function f is given by:

$$\hat{f}^h(\mathbf{p}_{N+1}) = \sum_{i=1}^{N^h} \hat{f}_i \ \psi_i(\mathbf{p}_{N+1}) = \boldsymbol{\psi}(\mathbf{p}_{N+1})^T \mathbf{\hat{f}}^h;$$

its expected value and variance are given by

$$\operatorname{var}(\hat{f}^h(\mathbf{p}_{N+1})) = \boldsymbol{\psi}(\mathbf{p}_{N+1})^T \operatorname{var}(\mathbf{\hat{f}}^h) \boldsymbol{\psi}(\mathbf{p}_{N+1}) = \sigma^2 \boldsymbol{\psi}(\mathbf{p}_{N+1})^T \mathbf{A}^{-1} \boldsymbol{\Psi}^T \boldsymbol{\Psi} \mathbf{A}^{-1} \boldsymbol{\psi}(\mathbf{p}_{N+1}).$$

An estimate of  $\operatorname{var}(\hat{f}^h(\mathbf{p}_{N+1}))$ ,  $\widehat{\operatorname{var}}(\hat{f}^h(\mathbf{p}_{N+1}))$  is then given by:

$$\widehat{\operatorname{var}}(\widehat{f}^h(\mathbf{p}_{N+1}) = \widehat{\sigma}^2 \psi(\mathbf{p}_{N+1})^T \mathbf{A}^{-1} \mathbf{\Psi}^T \mathbf{\Psi} \mathbf{A}^{-1} \psi(\mathbf{p}_{N+1}).$$

These results fully characterize the estimates in the case of Gaussian noise. In such case, one can derive confidence bands on the estimated functions and thus quantifying the uncertainty of the estimations of any predicted value. If the Gaussian assumption does not hold, this confidence interval should be used with caution and the confidence level is only approximated. More generally, the quantification of uncertainty has been widely studied in the context of generalized additive models (see, e.g., Wood (2006) and references therein). A Bayesian approach to uncertainty quantification for this class of models is also possible; see Marra and Wood (2012).

#### 4 Numerical simulations

In order to assess the IGS methodology, we consider two simulations on surfaces represented by NURBS for which the function f to be estimated is given a priori. We aim at showing different properties of IGS in different settings. In the first simulation, the configuration is such that any method used for planar smoothing such as Thin Plate Splines (TPS) (Duchon, 1977; Wahba, 1990) can be efficiently used and thus compared to IGS. In the second simulation, the setting is such that methods for two dimensional smoothing are not appropriate, while IGS can be straightforwardly used.

#### 4.1 Simulation 1

We consider a quarter of cylinder  $\Sigma$  defined as  $\Sigma = \{x^2 + y^2 = 1, 0 < z < 2, x > 0, y > 0\}$ . Using cylindrical coordinates, it is parametrized by the following mapping:

$$X: \Omega = (0,2) \times \left(0, \frac{\pi}{2}\right) \to \Sigma, \quad X(s_1, s_2) = (\cos(s_1), \sin(s_1), s_2).$$

This is an isometric mapping (Stoker, 1989). That means that the mapping preserves lengths and angles, and thus the area. In other words, the parametric domain  $\Omega$  is not distorted by the mapping. This kind of mapping allows to indifferently work on the parametric domain  $\Omega$  or directly on the surface  $\Sigma$ . Thus, the smoothing can be performed on the parametric domain, which is planar, namely using any traditional method for planar smoothing. In particular, in the following, we shall compare the proposed technique to TPS.

The surface  $\Sigma$  is exactly representable using NURBS basis functions of degree 2 or higher which are at least globally  $C^1$ —continuous. We are interested in recovering the function:

$$f(x, y, z) = \sin\left(\frac{5\pi}{2} \left[ xy^2 - y\left(\frac{z}{2} - 1\right)^2 + x^2\left(\frac{z}{2} - 1\right) \right] + \frac{\pi}{3} \right),$$

from noisy and discrete observations. The quarter of cylinder  $\Sigma$  and the function f are shown in Figure 3. The function f is evaluated in N=100 points  $\{\mathbf{p}_i\}_{i=1}^N$  located on a lattice on  $\Sigma$  and is affected by independent Gaussian observational errors  $\varepsilon_i$  of zero mean and standard deviation  $\sigma=0.125$ . We generate the data  $y_i=f(\mathbf{p}_i)+\epsilon_i$  for  $i=1,\ldots,N$  and the simulation is repeated M=100 times. The dimension of the IGA space  $\mathcal{F}^h$  varies between  $N^h=49$  and  $N^h=121$  and we use different NURBS basis functions, namely globally  $C^1-$ ,  $C^2-$ , and  $C^3-$  continuous of degrees p=2,3, and 4 respectively, obtained from k-refinement. The smoothing parameter  $\lambda$  is chosen at each simulation repetition and for each basis setting considered, by minimizing GCV criterion in (3.15).

The first two estimated functions  $\hat{f}^h$  over the M ones are displayed in Figure 4. We observe that these are qualitatively good estimates of f. In general, the number of basis functions  $N^h$  must be chosen carefully. Indeed, this choice should depend both on the complexity of the function to be estimated f and of the number of data points N available. When the number of basis functions  $N^h$  is small, IGS is not able to capture the behaviour of the function f, as it would be the case with any other smoother. On the contrary, when the number of basis functions  $N^h$  is high, we see that there is a larger variability in the estimated functions  $\hat{f}^h$ . Indeed, when the number of basis used is too high, GCV criterion can lead to overfitting, that is the estimated function incorporates noise. However, we see in Figure 4 that the estimates are not very sensitive to this choice. Finally,

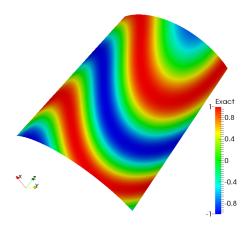


Figure 3: Simulation 1. Surface  $\Sigma$  (quarter of cylinder) and exact function f.

$N^h$	25	36	49	64	81	100	121	144	169	196
$K_{\infty}(\mathbf{A})$	46.1	47.3	48.2	46.5	43.9	80.9	210	431	355	307

Table 1: Simulation 1. Mean condition number  $K_{\infty}(\mathbf{A})$  for the matrix  $\mathbf{A}$ , with smoothing parameter chosen with GCV, p=2 and for different number of basis functions  $N^h$ .

we remark that the minimum number of basis functions is basically dictated by the number of functions used to represent the surface with NURBS.

We notice, following Section 3.3, that we used natural boundary conditions, for which we have not formally proved the well-posedness of the problem. We report in Table 1 the mean condition number  $K_{\infty}$  for the matrix A of (3.12), with  $\lambda$  chosen with the GCV criterion and for p=2 only, since results for p=3 and p=4 are similar. As a matter of fact, the system (3.12) results to be well-conditioned in all our numerical experiments. In order to better assess the IGS method, we use the empirical mean function  $\widehat{f}^h = \frac{1}{M} \sum_{i=1}^M \widehat{f}^h_i(\mathbf{x})$  and the associated empirical variance function  $\frac{1}{M} \sum_{i=1}^M (\widehat{f}^h_i(\mathbf{x}) - \widehat{f}^h(\mathbf{x}))^2$ , where  $\widehat{f}^h_i$  denotes the *i*-th estimated function. These are both shown in Figure 5, in the case of globally  $C^1$ -continuous basis functions of degree p=2 and with  $N^h=81$  basis functions. The estimates appears to have a negligible bias and a relatively small variance.

We compare our methodology with a widely used smoothing technique, TPS, using the cylindrical coordinates. This method is implemented for instance in the R package mgcv (Wood, 2015). We use different number of basis functions  $N^h$  for a comparison, selecting the smoothing parameter at each simulation repetition and for each number of basis considered via GCV. As a criterion for the comparison, we use the mean squared error (MSE) of the estimator, computed as:

$$MSE = \frac{1}{L} \sum_{j=1}^{L} \left( \hat{f}^h(\mathbf{l}_j) - f(\mathbf{l}_j) \right)^2,$$

where  $\mathbf{l}_1, \dots, \mathbf{l}_L$  is a lattice of  $150 \times 200$  evaluation points on  $\Sigma$ . We compute the MSE for  $N^h = 49, 64, 81, 100$ , and 121 and degree p = 2, 3, and 4 for the IGS and for  $N^h = 40, 50, \dots, 100$  for TPS. The comparisons of MSE in terms of the equivalent degree of freedom is shown in Figure 6. In the case of IGS, we see that the equivalent degrees of freedom is always less than the number of basis functions, hence meaning that the estimated functions are always penalized. We also compare in

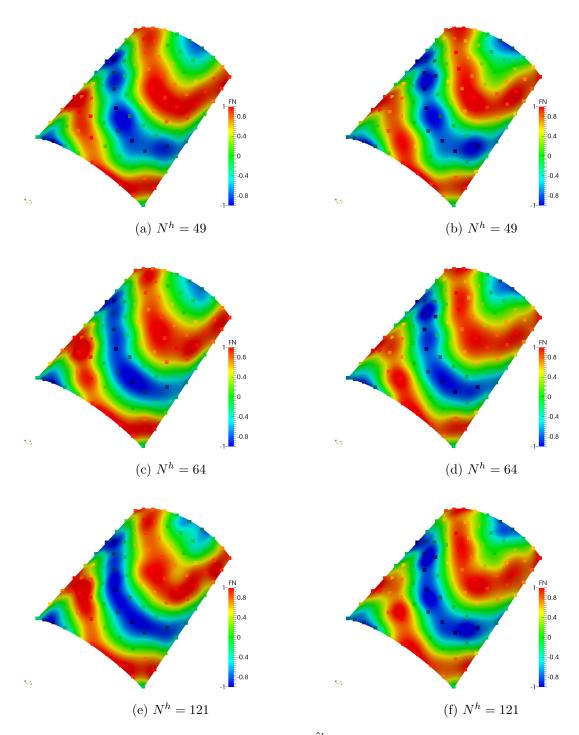


Figure 4: Simulation 1. Estimated functions  $\hat{f}^h$  in the first (left) and second (right) simulation repetition, out of M=100 repetitions, using  $N^h=49$  (top),  $N^h=64$  (middle) and  $N^h=121$  (bottom) number of basis functions, globally  $C^1$ -continuous of degree p=2. The corresponding measured values  $\{y_i\}_{i=1}^N$  are displayed on the same color scale as the true function and estimates.

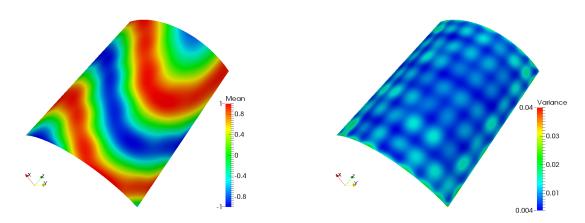


Figure 5: Simulation 1. Empirical mean function  $\overline{\hat{f}^h}$  (left) and empirical variance of the function  $\hat{f}^h$  (right) over the M=100 simulation repetitions, for NURBS basis functions of degree p=2, globally  $C^1$ — continuous and of dimension  $N^h=81$ .

Figure 7 the best setting of each methodology and for each number of basis functions, that is the lowest median MSE. Finally, Figures 6 and 7 show that IGS is comparable to TPS in terms of performance, in a setting where the latter technique may be applied.

#### 4.2 Simulation 2

We consider the surface  $\Sigma$  reported in Figure 8 which is represented in terms of NURBS basis functions of degree p=2 starting from the knot vector  $\Xi=\{0,0,0,1,1,1\}$  along both the parametric directions and control points  $\mathbf{P}_1=(1,0,0)^T$ ,  $\mathbf{P}_2=(0,0.75,0)^T$ ,  $\mathbf{P}_3=(0,1,0)^T$ ,  $\mathbf{P}_4=(1,0,1)^T$ ,  $\mathbf{P}_5=(1,1,0.5)^T$ ,  $\mathbf{P}_6=(0,1,1)^T$ ,  $\mathbf{P}_7=(0.25,0,1)^T$ ,  $\mathbf{P}_8=(0.25,0.25,0.25)^T$ , and  $\mathbf{P}_9=(0.25,0,1)^T$ , and with the corresponding weights vector  $\mathbf{w}=(1,1/\sqrt{2},1,1/\sqrt{2},1/2,1/\sqrt{2},1,1/\sqrt{2},1)^T$ . As reported in Figure 8, we consider the exact function to be estimated:

$$f(x, y, z) = (2 (x \sin(\pi y) + \cos(\pi x)y) - 1) \cos(\frac{5}{4}\pi z),$$

which is evaluated in N=100 points  $\{\mathbf{p}_i\}_{i=1}^N$  located on  $\Sigma$ ; the N observations are affected by independent Gaussian observational errors of zero mean and standard deviation  $\sigma=0.125$ . The sampling of the data  $y_i=f(\mathbf{p}_i)+\epsilon_i,\ i=1,\ldots,N$ , is repeated M=50 times. The sampling locations  $\{\mathbf{p}_i\}_{i=1}^N$  remain the same for all the repetitions. We consider three NURBS spaces  $\mathcal{F}^h$  of dimensions  $N^h=100$ , of globally  $C^1-$ ,  $C^2-$  and  $C^3-$  continuous NURBS basis functions of degrees p=2,3, and 4, respectively. These are obtained by k-refinement of the original NURBS basis.

In Figure 9 we report the estimated functions  $\hat{f}^h$  for the last two simulation repetitions, by considering NURBS spaces of basis functions of degree p=3 and globally  $C^2$ -continuous on  $\Sigma$ . Figure 10 highlights the empirical mean function  $\hat{f}^h$  and the corresponding empirical variance function over the M=50 simulation repetitions, obtained for the same NURBS basis functions. The results obtained for the degrees p=2 and 4 are very similar and are not reported here for the sake of brevity. Our estimation has a negligible bias and a small variance, as shown in Figure 10.

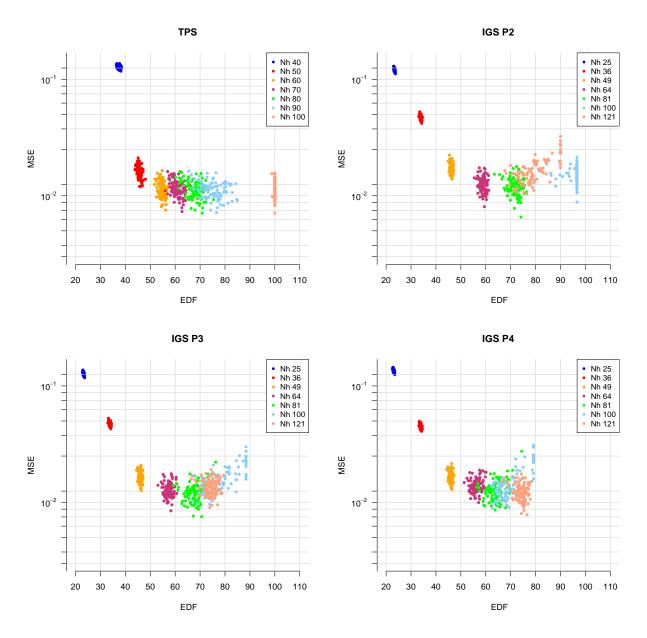


Figure 6: Simulation 1. Comparisons of the MSE for IGS and TPS (top-left), in terms of the equivalent degrees of freedom (EDF). IGS uses globally  $C^1-$ ,  $C^2-$  and  $C^3-$ continuous basis functions of degrees p=2 (top-right), 3 (bottom-left) and 4 (bottom-right), respectively.

The behaviour of the function f seems to be very well captured by our estimator  $\hat{f}^h$ , as shown in Figure 9. In this setting, we increase the regularity of the basis functions without changing the number of basis functions  $N^h$ . The estimated functions are thus globally  $C^{p-1}$ —continuous. We then compare the mean integrated squared errors (MISE) of IGS for different regularity and degrees of NURBS basis functions. The MISE of any estimated function  $\hat{f}$  is defined as:

$$\mathrm{MISE}(\hat{f}) = \frac{1}{|\Sigma|} \int_{\Sigma} \left( \hat{f} - f \right)^2 \ d\Sigma.$$

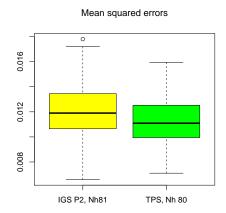


Figure 7: Simulation 1. Comparisons of the MSE for the best setting of IGS and TPS.

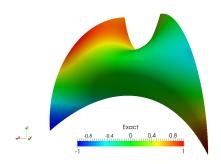


Figure 8: Simulation 2. Surface  $\Sigma$  and exact function f

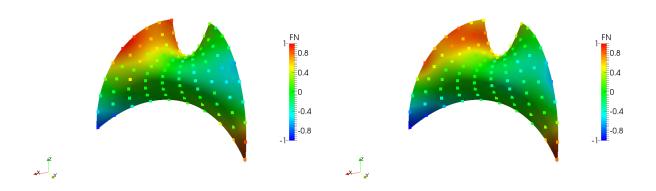


Figure 9: Simulation 2. Functions  $\hat{f}^h$  estimated for the last two repetitions, for N=100 points; positions of the points  $\{\mathbf{p}_i\}_{i=1}^N$  and the corresponding measured values  $\{y_i\}_{i=1}^N$  for NURBS basis functions of degrees p=3 and globally  $C^2$ -continuous on  $\Sigma$ .

As we observe in Figure 11, the quality of the estimation is not significantly affected by the regularity of the basis functions.



Figure 10: Simulation 2. Empirical mean function  $\widehat{f}^h$  (left) and empirical variance of the function  $\widehat{f}^h$  (right) for N=100 and M=50 obtained with NURBS basis functions of degrees p=3 and globally  $C^2$ -continuous on  $\Sigma$ .

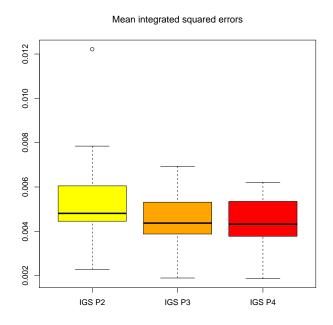


Figure 11: Simulation 2. Boxplots of the MISE for IGS with degrees p=2,3, and 4, globally  $C^1-, C^2-$  and  $C^3-$ continuous basis functions.

#### 5 Estimation of aerodynamic force on the SOAR's winglet

We aim at estimating the pressure coefficient field  $C_p$  and the corresponding aerodynamic force on the inboard winglet  $\Sigma$  of the SOAR shuttle shown in Figure 1. The pressure coefficient is a dimensionless field related to the pressure field. It describes the relative pressure over the surface of the winglet and is given by:

$$C_p(\mathbf{x}) = \frac{p(\mathbf{x}) - p_{\infty}}{\frac{\rho_{\infty}}{2} v_{\infty}^2},$$

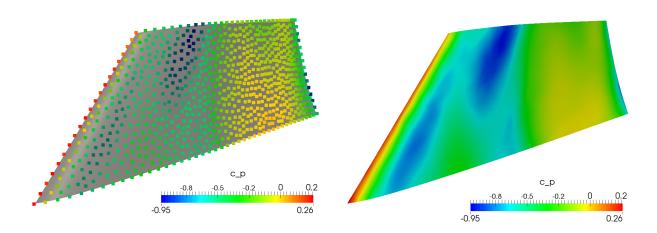


Figure 12: Original data (left) and smoothed estimated pressure coefficient  $\hat{C}_p^h$ , with N=824 data points and  $N^h=584$  basis functions (right).

where  $p(\mathbf{x})$  is the pressure at point  $\mathbf{x}$ ,  $v_{\infty}$  and  $p_{\infty}$  are the far field wind speed and pressure, with  $\rho_{\infty}$  the air density. The flow regime of the free stream is subsonic, namely the Mach number is 0.7, for which we can reasonably assume that the pressure coefficient field over the winglet remains sufficiently smooth, since transonic effects are marginal. The force  $\mathbf{F}$  acting on the winglet and due to the contribution of the pressure (see, e.g., Anderson (2010)) is given by:

$$\mathbf{F} = \int_{\Sigma} (p - p_{\infty}) \mathbf{n}_{\Sigma} d\Sigma = \frac{\rho_{\infty}}{2} v_{\infty}^{2} \int_{\Sigma} C_{p} \mathbf{n}_{\Sigma} d\Sigma,$$

where  $\mathbf{n}_{\Sigma}$  is the unit normal vector to the surface  $\Sigma$ . The final quantity of interest of the estimation is the aerodynamic force  $\mathbf{F}$ . The pressure coefficient  $C_p$  is measured in N=824 data points on the surface  $\Sigma$  for which the sampled data are depicted in Figure 12, which shows the data after application of an affine transformation<sup>2</sup>, that modifies the values of about 3% at most. These data are derived from Computational Fluid Dynamics (CFD) simulations and represent a preliminary study prior experimental campaign in a wind tunnel. The inboard winglet is represented by a single NURBS patch surface built with  $N^h=584$  degrees of freedom. The geometry is built with a NURBS basis of degree p=5 and functions globally  $C^2$ -continuous. Thus, the minimum number of basis functions representing the function space  $\mathcal{F}^h$  for IGS is  $N^h=584$ .

First, we estimate the pressure coefficient field  $\hat{C}_p^h$  with all the available points. We use IGS to estimate the pressure coefficient field and then the aerodynamic force. By using IGS with N=824 and  $N^h=584$ , for NURBS basis functions of degree 5 and globally  $C^2$ -continuous basis functions, we obtain the  $\hat{C}_p^h$  field reported in Figure 12 (right).

We now assess the quality of IGS under the hypothesis that less data points than those effectively available can be used. Indeed, in industrial applications, the experimental measurements can be quite expensive and it could be impractical to have a large number of sampling points (in this case, pressure probes), and the pressure field can be measurable in fewer points than with numerical simulations. To assess the validity of IGS when few data points are available, we compute the aerodynamic force  $\hat{\mathbf{F}}^h$  estimated from a subset of points over the winglet. We compare the results

<sup>&</sup>lt;sup>2</sup>The displayed data have been modified for copyright reasons with respect to the original ones provided by S3, Swiss Space Systems Holding SA.

$N^h$	N	difference of direction (in degrees)	$\ \mathbf{\hat{F}}^h - \mathbf{\hat{F}}_{\mathrm{ref}}^h\ /\ \mathbf{\hat{F}}_{\mathrm{ref}}^h\ (\mathrm{in}~\%)$
584	103	1.007	3.293
584	206	0.816	1.036
584	412	0.084	0.745
584	618	0.001	0.003
989	412	0.278	0.864
989	618	0.061	0.205
989	824	0.003	0.117
2079	824	0.0113	0.296

Table 2: Relative difference between vector of forces with respect to the reference force  $\hat{\mathbf{F}}_{\text{ref}}^h$ , estimated in the case where  $N^h = 584$  and N = 824, in degrees (left) and in norm (right). Direction (in degrees) and relative modulus change (in %) of the estimated aerodynamic force  $\hat{\mathbf{F}}^h$  obtained with different points N and basis functions  $N^h$  with respect to the reference force  $\hat{\mathbf{F}}_{\text{ref}}^h$ .

in terms of the aerodynamic force, which is the quantity of interest, by using fewer data points and different number of basis functions with respect to the reference setting corresponding to  $N^h = 584$  and N = 824. We compare the forces in terms of the relative difference of the modules (in %), that is  $\|\hat{\mathbf{F}}^h - \hat{\mathbf{F}}^h_{\text{ref}}\|_2 / \|\hat{\mathbf{F}}^h_{\text{ref}}\|_2$  and of the direction (angle in degrees). The results are reported in Table 2. We first use all the N = 824 available points depicted in the Figure 12 (left) to estimate the pressure coefficient field. The estimation of the reference force is given by  $\hat{\mathbf{F}}^h_{\text{ref}} = (0.5080, 49.40, -26.03)^T$  kN, which is in line with the expectations.

We then use only N=618,412,206, and 103 points and NURBS spaces of dimensions  $N^h=584,989$ , and 2079. In Figure 13, we report the estimated pressure coefficient fields on the surface computed using IGS. The choice of the smoothing parameter is done via minimization of the GCV criterion. The estimated functions are very similar when the number of points is relatively large (N=824,618, and 412), as shown in Figure 13. The relative change in the estimated force is quite small, as we can observe on Table 2. IGS is able to accurately estimate the force even with a small subset of the original data.

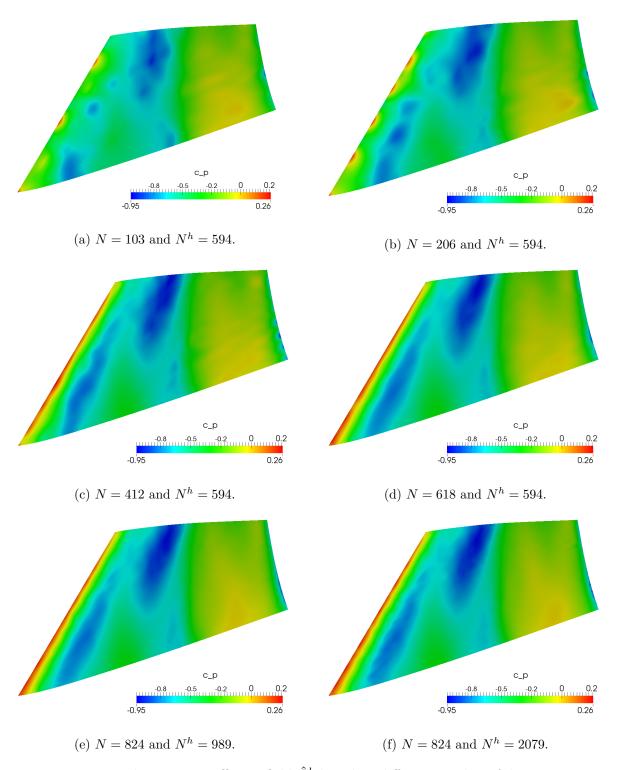


Figure 13: Estimated pressure coefficient field  $\hat{C}_p^h$  based on different number of data points N and number of basis functions  $N^h$ .

#### 6 Conclusion and discussion

In this paper, we proposed a methodology, IGS, to deal with functional data defined on surfaces described by NURBS, specifically to reconstruct the functions from noisy observations. The proposed IGS method has the potential of being widely applicable, in particular in industrial contexts where geometries are commonly defined by NURBS. Simulations indicate that IGS is comparable to other widely used methods, such as TPS, in cases where the latter is applicable. Moreover, IGS avoids the use of complex meshing procedures since the geometry of the surface is directly used as a data of the problem. IGS is also computationally efficient; this is due to the fact that the NURBS basis functions have compact support and thus the matrices involved in the computations are usually very sparse. These are very convenient features, especially if the number of basis functions is relatively large (e.g. in the case of a complex surface) and/or when there are many data points.

A differential penalization is equivalent to the assumption that the function to be estimated must be close to the kernel of the penalization operator (Green and Silverman, 1993). In the case where the observed physical phenomenon is driven by a known PDE, it would be convenient to use a penalization operator related to the PDE (Azzimonti et al., 2014b). Since IGS is based on IGA, it would be easy to implement other kind of penalization. IGS also allows the application of different kind of boundary conditions, which can be useful in practical applications (Wood et al., 2008; Sangalli et al., 2013; Azzimonti et al., 2014a). In addition, IGS can straightforwardly deal with data in one, two, or even three dimensions. Indeed, all the results presented here can be extended to any lower dimensional manifold defined by NURBS. Moreover, as shown in the application, functionals of an estimated field, as the aerodynamic force, can be easily computed. In addition, IGS is very flexible and allows local refinement of the basis functions, which can be needed when the distribution of the data points is not uniform. Finally, we remember that it is also possible to extend this model to take account of spatially varying covariates on the manifold, using a generalized additive framework.

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