# Theory evaluation of LHC data for Physics beyond the Standard Model

#### Riccardo Rattazzi



## After the Higgs and Nothing Else

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#### Effective Quantum Field Theory ideology

- **♦** gauge symmetry
- **♦** field content
- ♦ local effective lagrangian

- \* renormalizability
- \* global symmetries

needn't be fundamental but just accidental low energy features The Standard Model as effective field theory

with fundamental scale  $\Lambda_{uv}^2 \gg 1 \, \text{TeV}$ 

$$\mathcal{L}_{SM} = \mathcal{L}_{kin} + gA_{\mu}\bar{F}\gamma_{\mu}F + Y_{ij}\bar{F}_{i}HF_{j} + \lambda(H^{\dagger}H)^{2}$$

d=4

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d=4

$$+ \frac{b_{ij}}{\Lambda_{UV}} L_i L_j H H$$

$$+ \frac{c_{ijkl}}{\Lambda_{UV}^2} \bar{F}_i F_j \bar{F}_k F_\ell + \frac{c_{ij}}{\Lambda_{UV}} \bar{F}_i \sigma_{\mu\nu} F_j G^{\mu\nu} + \dots$$

d>4

$$\mathcal{L}_{SM} = \mathcal{L}_{kin} + gA_{\mu}\bar{F}\gamma_{\mu}F + Y_{ij}\bar{F}_{i}HF_{j} + \lambda(H^{\dagger}H)^{2}$$

$$+ \frac{b_{ij}}{\Lambda_{UV}}L_{i}L_{j}HH$$

$$+ \frac{c_{ijkl}}{\Lambda_{UV}^{2}}\bar{F}_{i}F_{j}\bar{F}_{k}F_{\ell} + \frac{c_{ij}}{\Lambda_{UV}}\bar{F}_{i}\sigma_{\mu\nu}F_{j}G^{\mu\nu} + \dots$$

$$+ \dots$$

$$+ \dots$$

 $\Lambda_{UV} \gg {
m TeV}$  (pointlike limit) nicely accounts for 'what we see'

$$+\theta \, \tilde{G}_{\mu\nu} \tilde{G}^{\mu\nu}$$

$$\mathcal{L}_{SM} = \mathcal{L}_{kin} + gA_{\mu}\bar{F}\gamma_{\mu}F + Y_{ij}\bar{F}_{i}HF_{j} + \lambda(H^{\dagger}H)^{2}$$

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$$d > 4$$

 $\Lambda_{UV}\gg {
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$$+c\Lambda_{UV}^{2}H^{\dagger}H$$

$$+\theta \tilde{G}_{\mu\nu}\tilde{G}^{\mu\nu}$$

$$\mathcal{L}_{SM} = \mathcal{L}_{kin} + gA_{\mu}\bar{F}\gamma_{\mu}F + Y_{ij}\bar{F}_{i}HF_{j} + \lambda(H^{\dagger}H)^{2}$$

$$+ \frac{b_{ij}}{\Lambda_{UV}}L_{i}L_{j}HH$$

$$+ \frac{c_{ijkl}}{\Lambda_{UV}^{2}}\bar{F}_{i}F_{j}\bar{F}_{k}F_{\ell} + \frac{c_{ij}}{\Lambda_{UV}}\bar{F}_{i}\sigma_{\mu\nu}F_{j}G^{\mu\nu} + \dots$$

$$+ \dots$$

$$d=2$$

$$d=4$$

$$+ \frac{b_{ij}}{\Lambda_{UV}}L_{i}L_{j}HH$$

$$+ \frac{c_{ijkl}}{\Lambda_{UV}^{2}}\bar{F}_{i}F_{j}\bar{F}_{k}F_{\ell} + \frac{c_{ij}}{\Lambda_{UV}}\bar{F}_{i}\sigma_{\mu\nu}F_{j}G^{\mu\nu} + \dots$$

 $\Lambda_{UV} \gg {
m TeV}$  (pointlike limit) nicely accounts for 'what we see'

$$+ \Lambda_{UV}^4 \sqrt{g}$$
 
$$+ c \Lambda_{UV}^2 H^\dagger H$$
 
$$+ e \tilde{G}_{\mu\nu} \tilde{G}^{\mu\nu}$$
 
$$+ \theta \tilde{G}_{\mu\nu} \tilde{G}^{\mu\nu}$$
 
$$= d = 0$$

$$\mathcal{L}_{SM} = \mathcal{L}_{kin} + gA_{\mu}\bar{F}\gamma_{\mu}F + Y_{ij}\bar{F}_{i}HF_{j} + \lambda(H^{\dagger}H)^{2}$$

$$+ \frac{b_{ij}}{\Lambda_{UV}}L_{i}L_{j}HH$$

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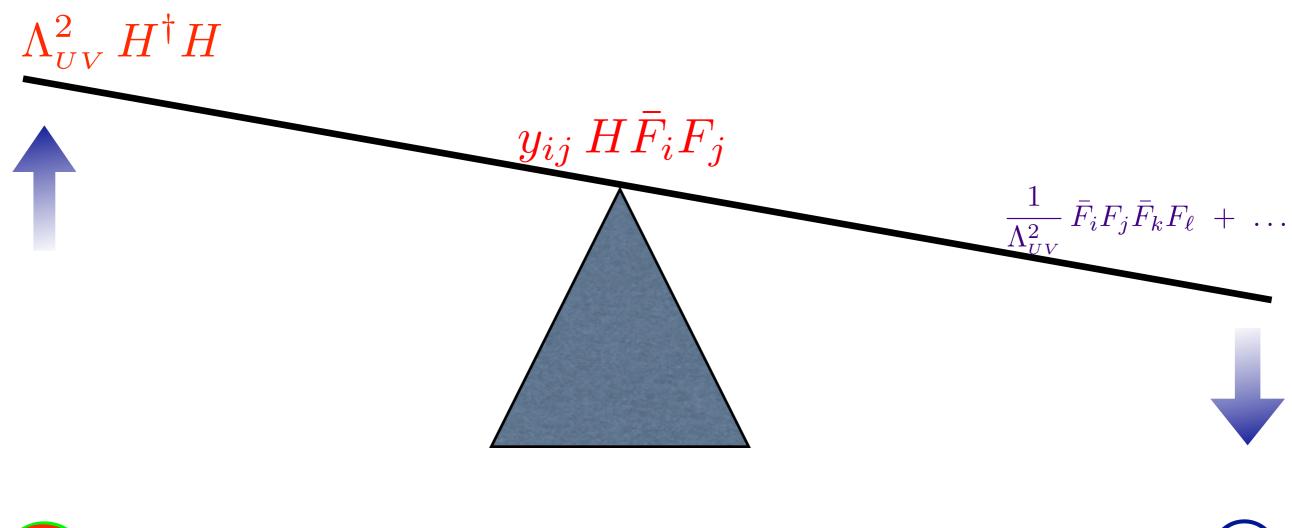
$$+ \dots$$

$$+ \dots$$

 $\Lambda_{UV} \gg {
m TeV}$  (pointlike limit) nicely accounts for 'what we see'

# Hierarchy see-saw

Standard Model up to some  $\Lambda_{UV}^2 \gg 1 \,\mathrm{TeV}$ 



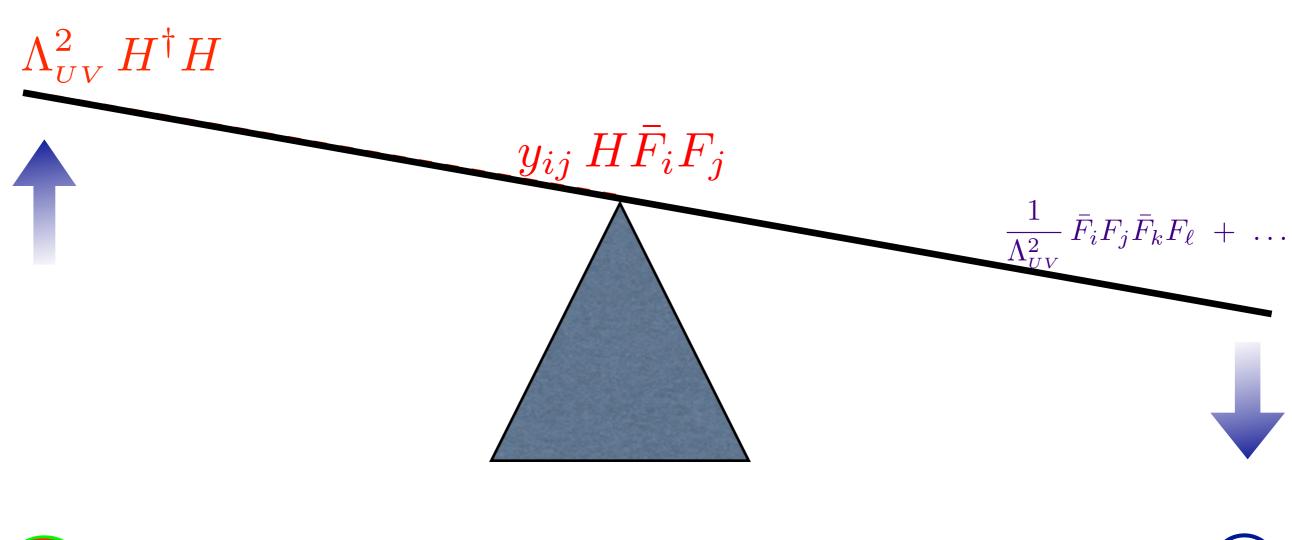






# Hierarchy see-saw

Standard Model up to some  $\Lambda_{UV}^2 \gg 1 \, {\rm TeV}$ 



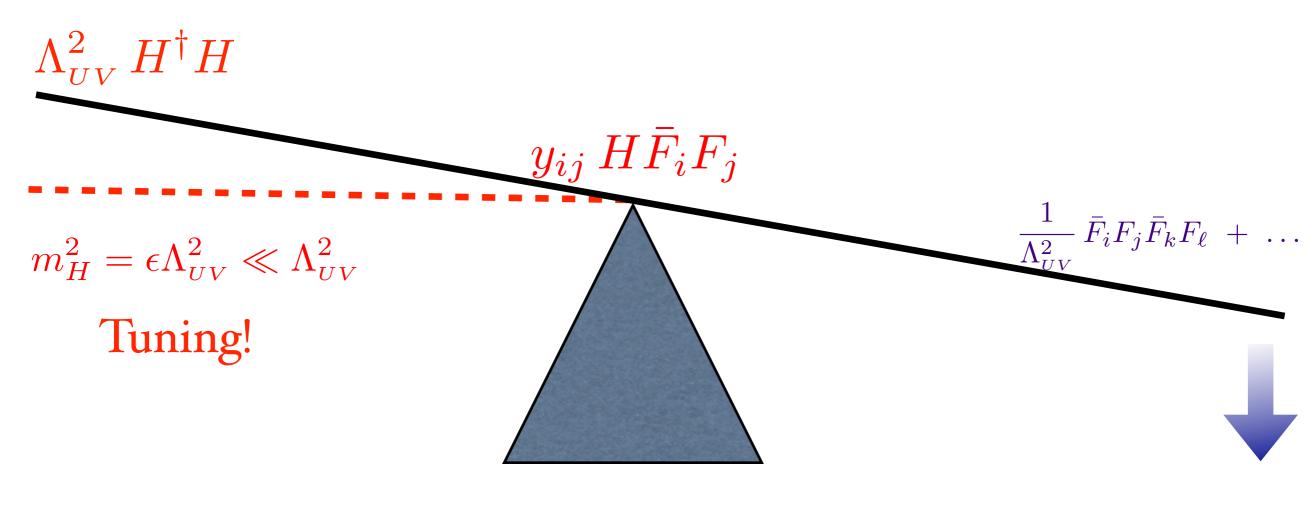






## Hierarchy see-saw

Standard Model up to some  $\Lambda_{UV}^2 \gg 1 \, \text{TeV}$ 









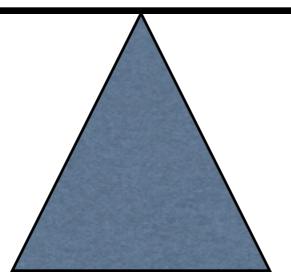
#### Natural SM:

$$\Lambda_{\scriptscriptstyle UV}^2 \lesssim 1\,{\rm TeV}$$

$$\Lambda_{\scriptscriptstyle UV}^2\,H^\dagger H$$

$$y_{ij} H \bar{F}_i F_j$$

$$\frac{1}{\Lambda_{UV}^2} \, \bar{F}_i F_j \bar{F}_k F_\ell + \dots$$









#### The two possible microphysics scenarios

- I. The SM is the correct description up to  $\Lambda_{UV} \gg TeV$ 
  - B, L and Flavor: beautifully in accord with observation
  - Hierarchy remains a mystery, probably hinting that the question was not correctly posed
    - anthropic principle
    - failure of effective field theory ideology (UV/IR connection)

- II. The SM is not the correct description already at  $\Lambda_{UV} \sim 1 \, {
  m TeV}$ 
  - In the correct theory the hierarchy problem does not even arise (naturalness)
  - What about B, L and Flavor? In all models not nearly as nice as in SM

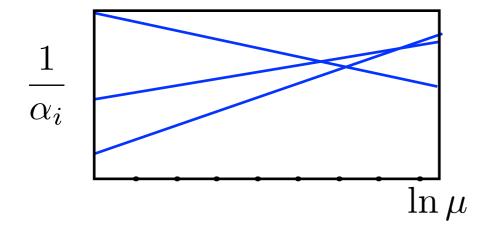
### A high scale scenario

- $\mathcal{L}^{d=4}$  experimental success (some 2- 3- $\sigma$  glitches here and there)
- $\Theta$ -QCD and Dark Matter  $\rightarrow$  high scale axion  $f_a \sim 10^{12} \; \mathrm{GeV}$

$$f_a \sim 10^{12} \text{ GeV}$$

• gauge couplings ready to unify around  $10^{15} \lesssim M \lesssim M_{Planck}$ 

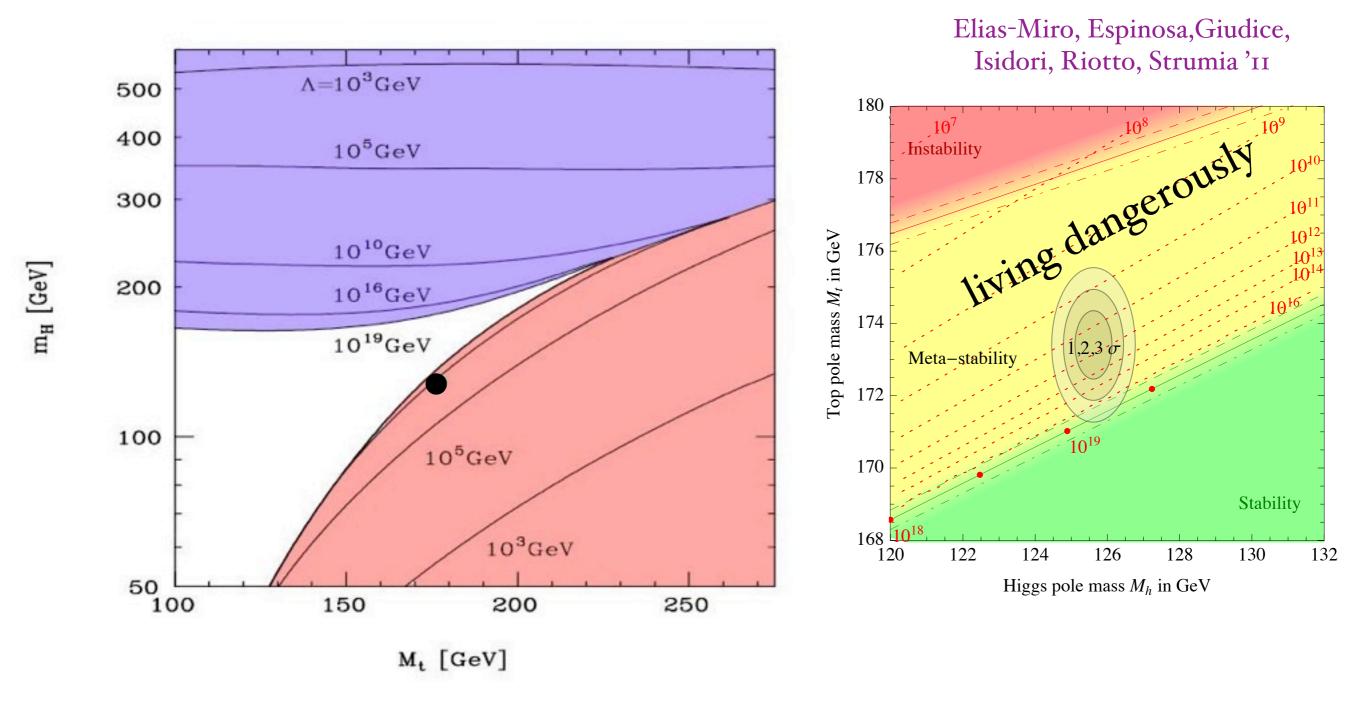
$$10^{15} \lesssim M \lesssim M_{Planck}$$



• neutrino masses

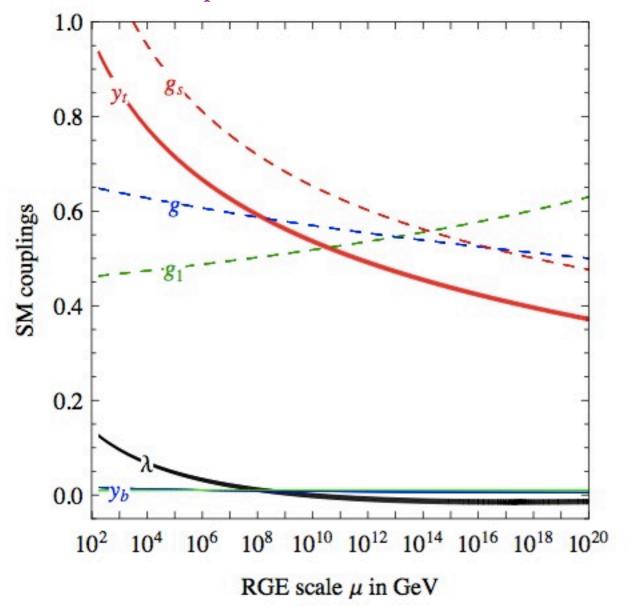
$$\frac{\ell\ell HH}{\Lambda} \longrightarrow m_{\nu} \sim \frac{v^2}{\Lambda} \qquad \Lambda \sim 10^{14} \text{ GeV}$$

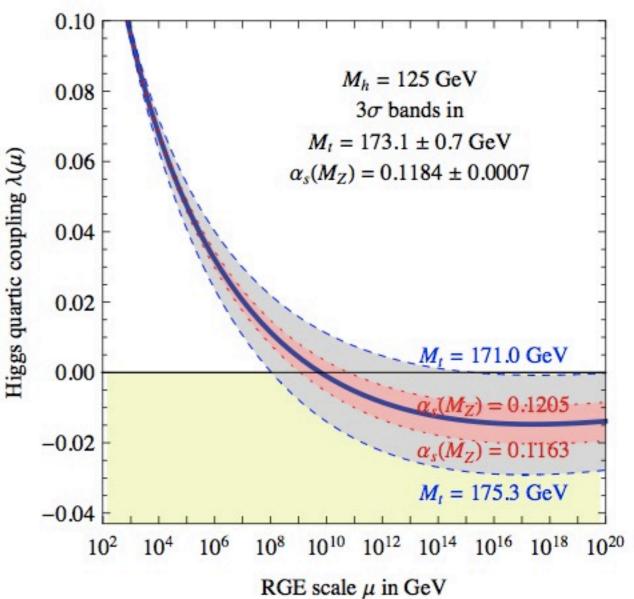
• RG-evolution of SM couplings, including  $\lambda_h$  remarkably do not require lower scales



the fact that the SM lives dangerously perhaps points to an anthropic selection of parameters though it is hard to tell

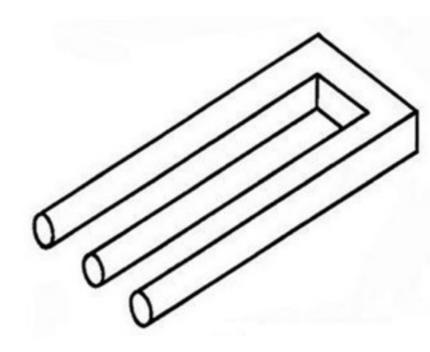
Elias-Miro, Espinosa, Giudice, Isidori, Riotto, Strumia '11

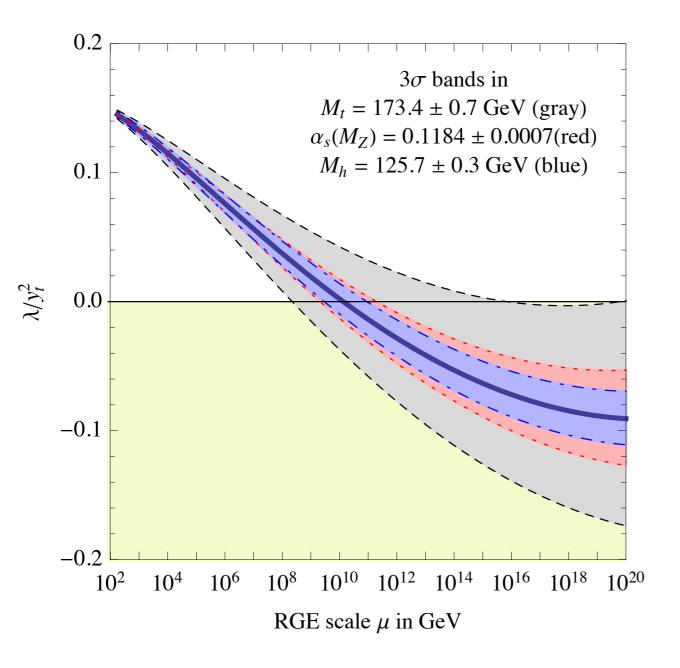


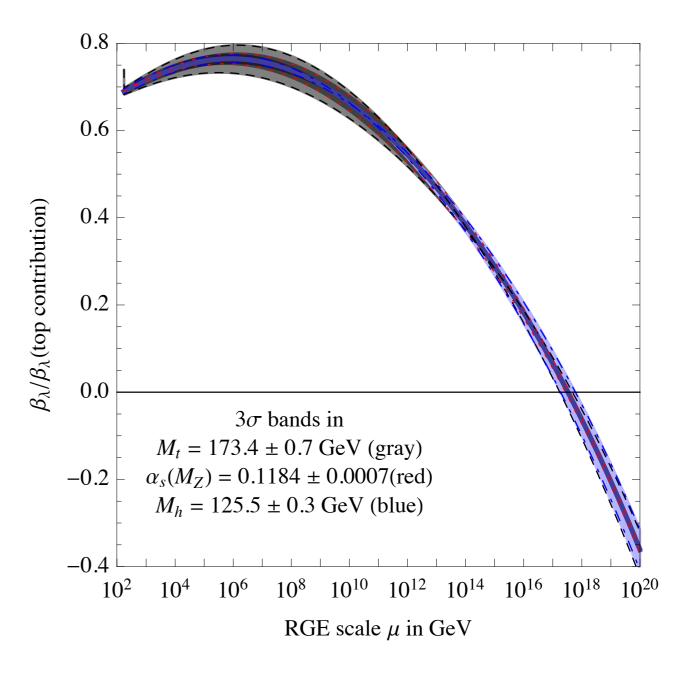


Elias-Miro, Espinosa, Giudice, Isidori, Riotto, Strumia '11 0.10 1.0  $M_h = 125 \text{ GeV}$ 0.08  $3\sigma$  bands in 0.8  $M_t = 173.1 \pm 0.7 \text{ GeV}$ Higgs quartic coupling  $\lambda(\mu)$ 0.06  $\alpha_s(M_Z) = 0.1184 \pm 0.0007$ 0.6 SM couplings 0.04 0.02 0.4  $M_t = 171.0 \text{ GeV}$ 0.00 0.2 -0.02 $\alpha_s(M_Z) = 0.1163$ 0.0 -yb  $M_t = 175.3 \text{ GeV}$ -0.04 $10^{8} \ 10^{10} \ 10^{12} \ 10^{14} \ 10^{16} \ 10^{18} \ 10^{20}$  $10^4 10^6$  $10^{10} \ 10^{12} \ 10^{14} \ 10^{16} \ 10^{18} \ 10^{20}$  $10^{2}$ 104  $10^{6}$ RGE scale  $\mu$  in GeV RGE scale  $\mu$  in GeV

but watch your eyes







#### Strumia's courtesy

The two natural scenarios for electroweak symmetry breaking

Elementary Higgs exists but a symmetry *protects* its mass

Supersymmetric Models

No elementary Higgs exists

Technicolor,

Composite Light Higgs

(and its holograms)

Large Extra Dimensions: exciting, fantastic, great, but not very plausible without extra mass scale separation

Plus a list of not even wrong scenarios...

Flavor?

Supersymmetry: the existence of scalar matter fields introduces a myriad of  $d \le 4$  terms violating F, B and L

$$\mathcal{L}^{d\leq 4} = m_{ij}^2 \tilde{Q}_i^{\dagger} \tilde{Q}_j + A_{ij} Y_{ij}^D \tilde{Q}_i \tilde{D}_j H_d + \lambda_{ijk} \tilde{U}_i D_j D_k + \dots$$

Naive Composite Higgs (TC): the Yukawa themselves are d > 4

$$Y_{ij} H \bar{Q}_L^i Q_R^j \to Y_{ij} \frac{1}{\Lambda_F^2} (\bar{\Psi} \Psi) Q_L^i Q_R^j$$
  $m_{ij} = Y_{ij} \frac{v_F^3}{\Lambda_F^2}$ 

 $\Lambda_F$  must be not too far above weak scale: expect unwanted FCNC

In all natural models, extra assumptions (often clever) are needed to meet flavor physics constraints

## Approaches

Symmetry

pick a subgroup of  $U(3)_Q \times U(3)_U \times U(3)_D \times U(3)_L \times U(3)_E$ pick a set of spurions to break it construct a lagrangian using the selection rules

Dynamics

mass mixing hierarchy from radiative corrections
flavor from geography in extra-dim
flavor from partial compositeness

holography

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### What naturalness demands



supersoft models: Higgs mass parameter fully saturated by IR physics

- •dirac gauginos in supersymmetry
- •general composite Higgs

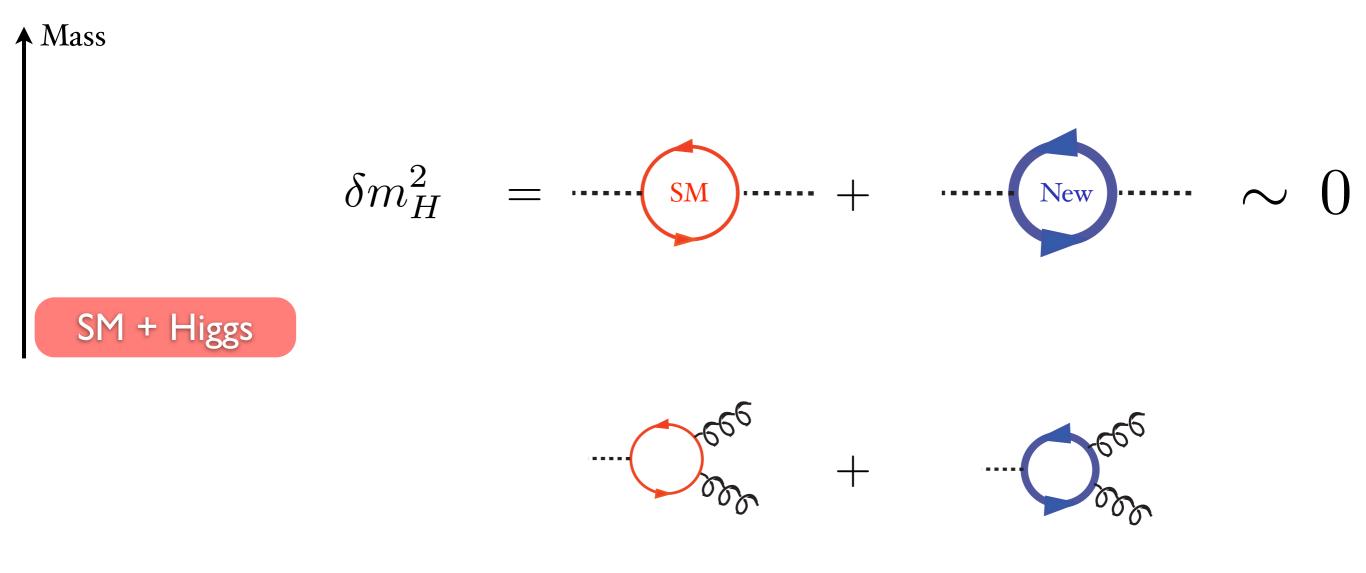
$$\delta m_H^2 \; = \; \textbf{---} \frac{3\lambda_t^2}{8\pi^2} \, M^2 \qquad m_h = 125 \, \text{GeV} \quad \left(\frac{400 \, \text{GeV}}{M}\right)^2$$



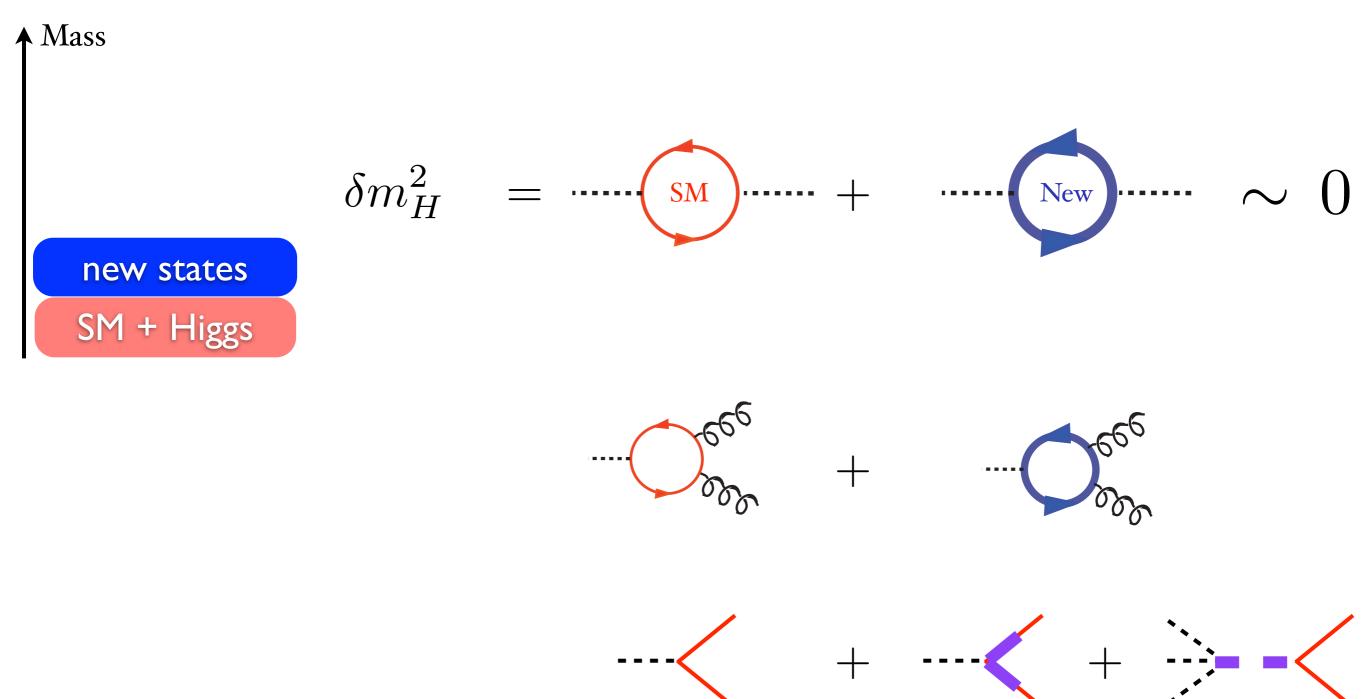
soft models: Higgs mass logarithmically sensitive to UV

•MSSM and its extensions with high scale mediation

$$m_h^2 \sim m_Z^2 \sim m_{\tilde{t}}^2$$
  $m_h = 125\,\mathrm{GeV}$  tuning  $\left(\frac{125\,\mathrm{GeV}}{\tilde{m}}\right)^2$ 



The more natural the theory the more the Higgs rates deviate from SM



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#### Supersymmetry

- •higgs couplings
- •EWPT  $(\epsilon_1, \epsilon_2, \epsilon_3)$
- $\bullet b \rightarrow s \gamma$ 
  - ...

$$\frac{\delta \mathcal{O}}{\mathcal{O}_{SM}} \sim \frac{m_{weak}^2}{\tilde{m}^2} \longrightarrow O(\alpha_t/4\pi)$$
super-soft

Compositeness

$$\frac{\delta \mathcal{O}}{\mathcal{O}_{SM}} \sim \frac{v^2}{f^2}$$

Little Higgs with T-parity

•EWPT 
$$(\epsilon_1, \epsilon_2, \epsilon_3)$$

- $b \rightarrow s \gamma$
- direct searches (susy)

imply 
$$\frac{\delta \mathcal{O}}{\mathcal{O}_{SM}} < 1 \qquad (10\%)$$

in most cases, O(1) deviations in Higgs rates were already disfavored

# Perspective on Supersymmetry before and after LHC 7/8

once upon a time, expectation in the less clever models was

$$Z$$
  $\qquad \qquad = \qquad \tilde{t}$   $\qquad \qquad \tilde{\ell}$   $\qquad \qquad \tilde{\ell}$   $\qquad \qquad \tilde{\chi}^+ \quad \tilde{\chi}^0$ 

LEP/Tevatron

$$m_{\tilde{l}}, m_{\chi^+} \gtrsim 100 \, \mathrm{GeV}$$

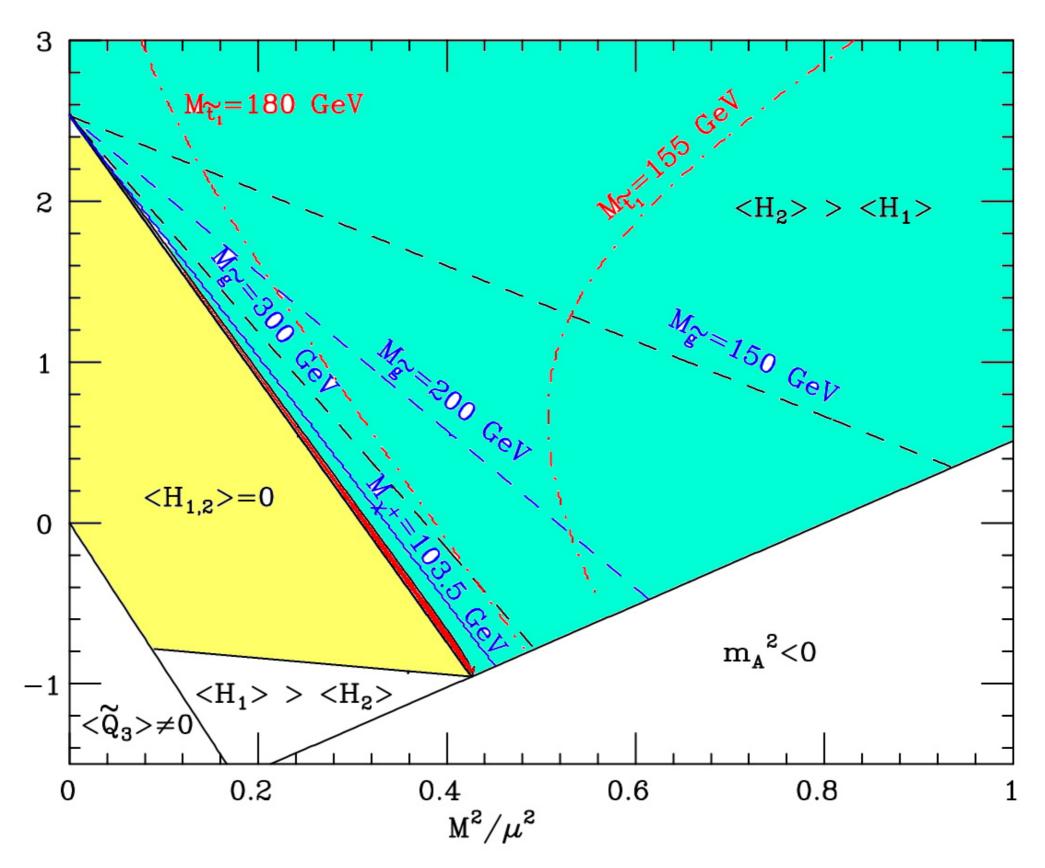
$$m_{\tilde{q}}, m_{\tilde{g}} \gtrsim 300 \, \mathrm{GeV}$$

 $m_h > 114.4 \,\mathrm{GeV}$ 

MSSM

 $1 \, \mathrm{TeV} \quad ---- \quad \tilde{t} \quad ---- \quad \tilde{v}^0$ 

Z ———



Giudice, Rattazzi '06

 $\begin{array}{ccc} & m_{\tilde{l}},\,m_{\chi^+}\,\gtrsim\,100\,{\rm GeV}\\ \\ {\rm LEP/Tevatron} & m_{\tilde{q}},\,m_{\tilde{g}}\,\gtrsim\,300\,{\rm GeV} \end{array}$ 

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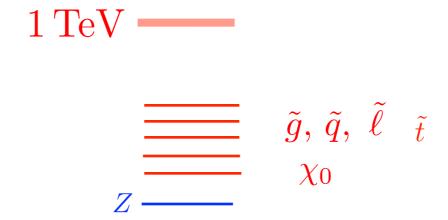
Before LHC: natural and simple spectrum possible within NMSSM

#### After LHC

- ullet situation only slightly worse with  $m_h \simeq 125\,\mathrm{GeV}$  Hall, Pinner, Ruderman '11
- GUT perturbativity borderline, but needn't worry too much

see Barbieri et al 2013

• direct searches have however eliminated the natural region of the most straighforward NMSSM scenario (in particular flavor universal sfermion masses)



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$$1 \text{ TeV} = \begin{bmatrix} \tilde{g}, \tilde{q}, \tilde{\ell} & \tilde{t} \\ \chi_0 & \end{bmatrix}$$

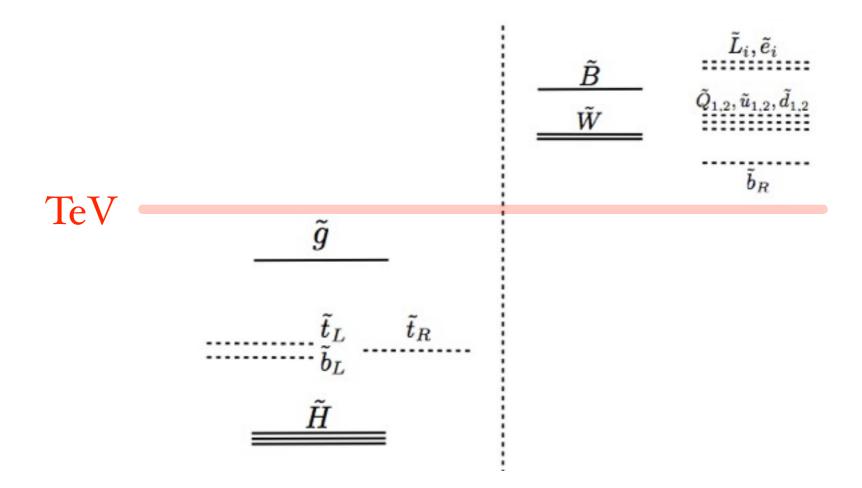
Z ———

MSSM: simplest scenarios were under pressure already before LHC

NMSSM: simplest scenarios came under pressure with LHC

What about more 'structured' models, those that stick to naturalness like a mussel to her reef?

#### not-so-un-Natural SUSY

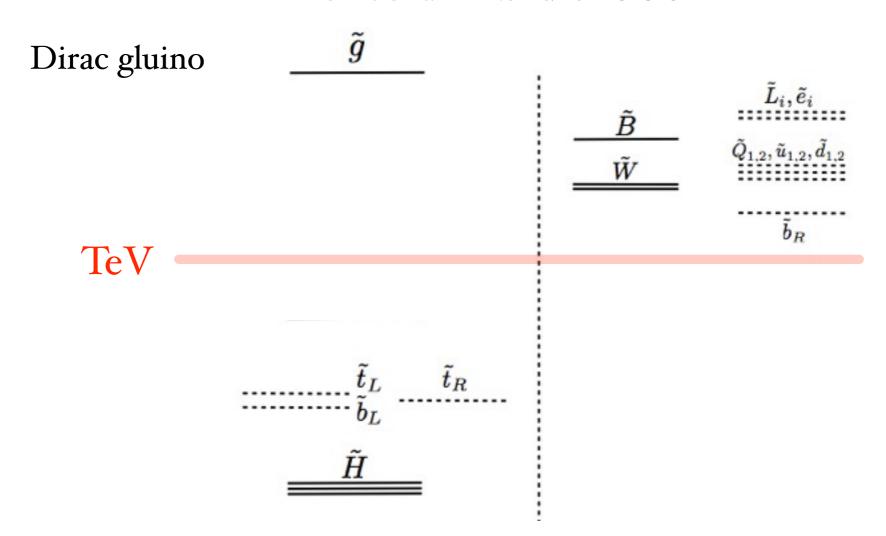


 $m_{\tilde{t}} \lesssim 500 \, \mathrm{GeV}\,$  some other physics ( ex NMSSM) takes care of  $\, m_h \,$ 

consider all possible ways to suppress the signal:

Dirac gluino, RPV, compressed spectrum,...

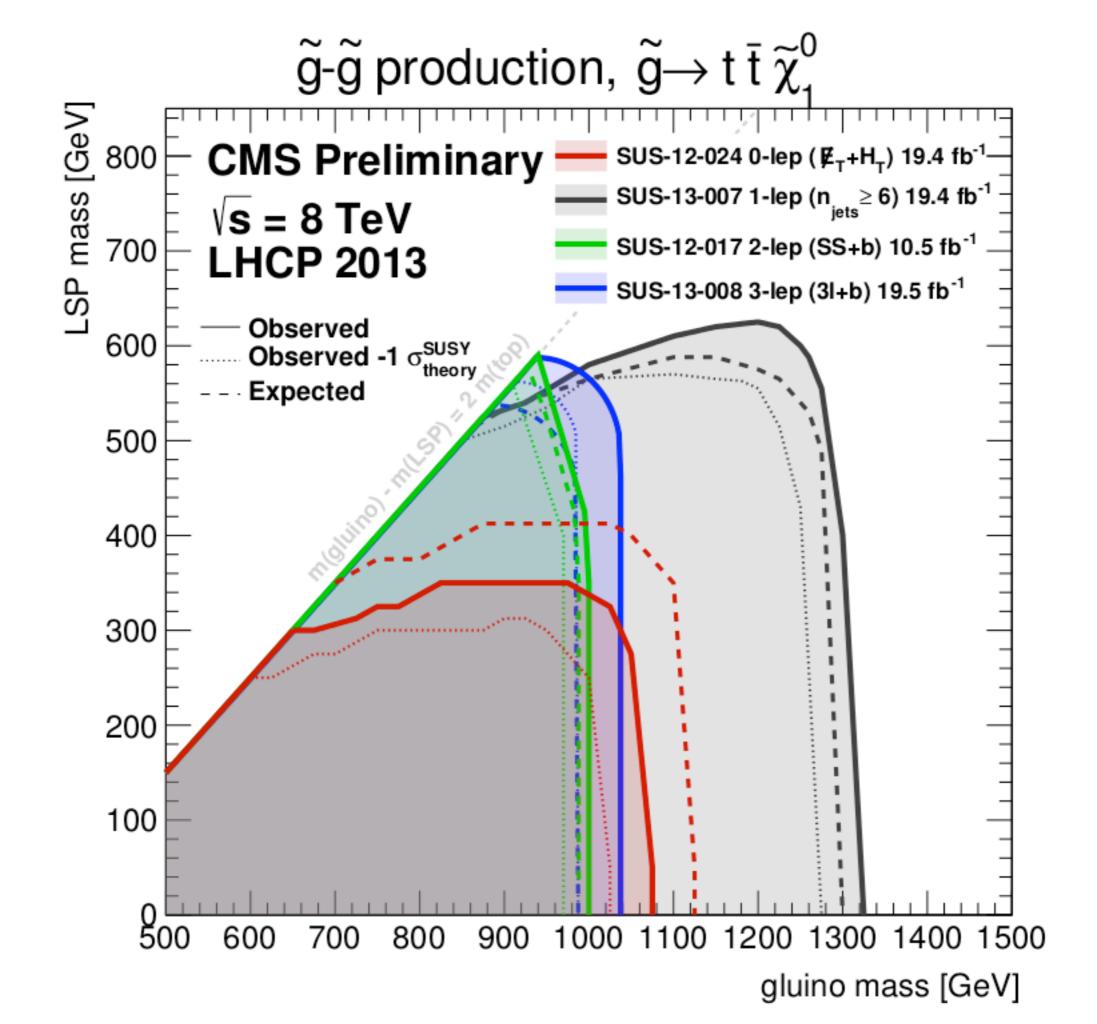
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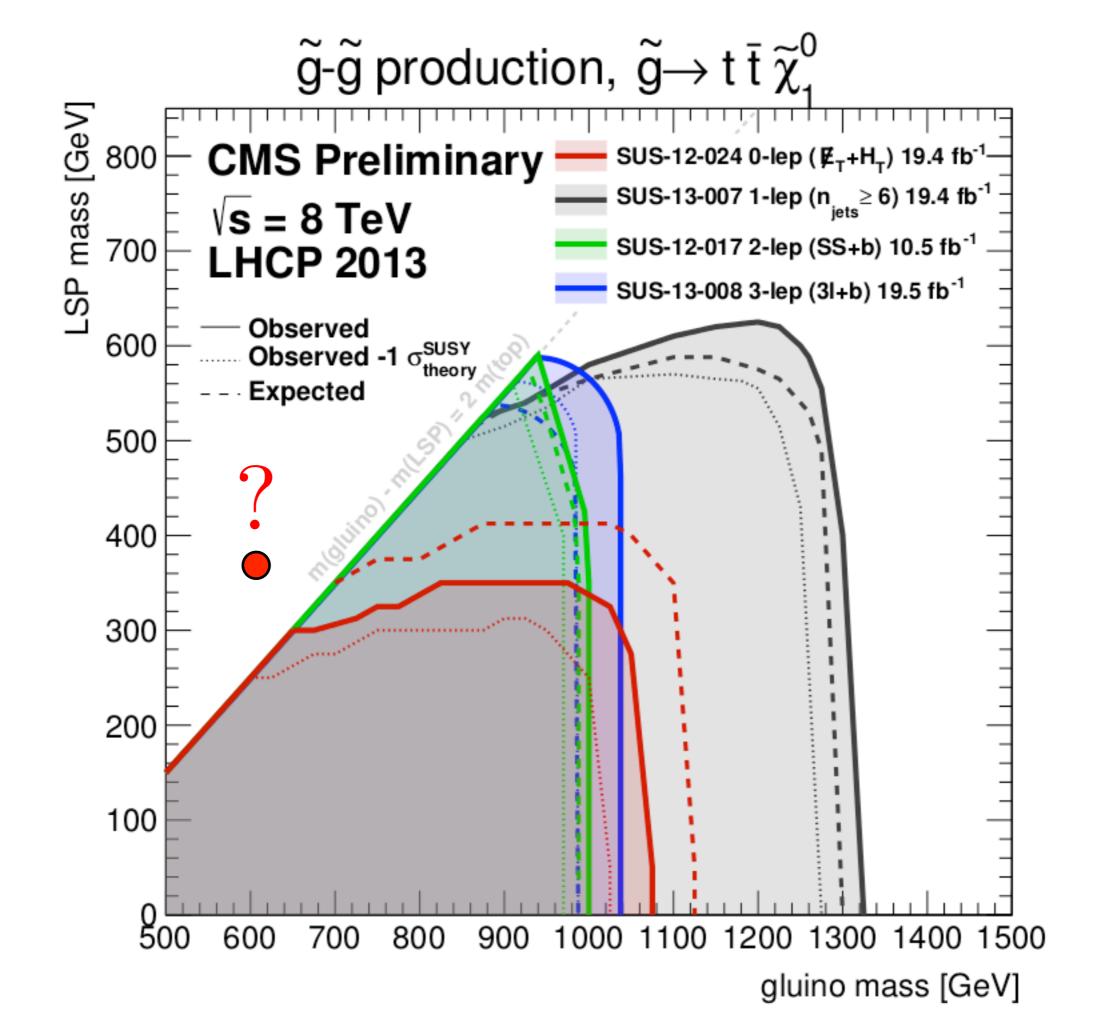


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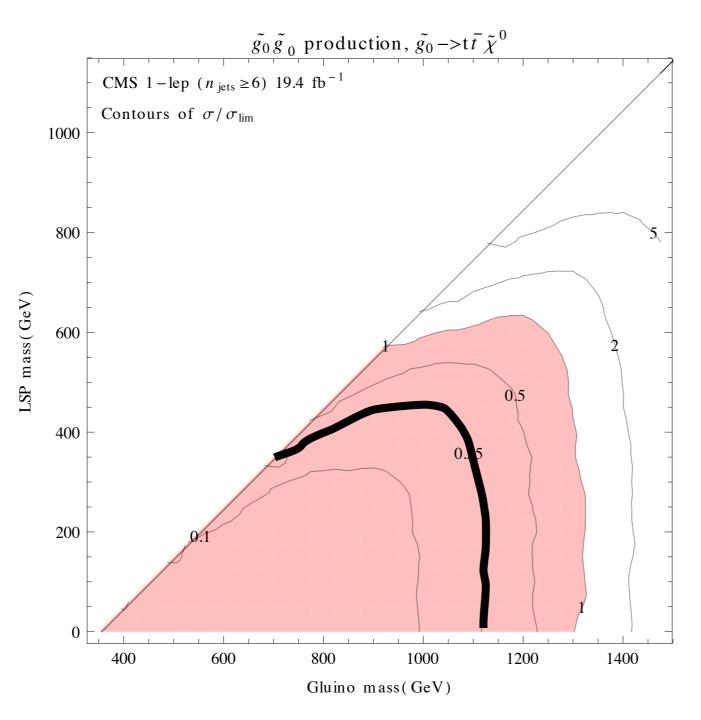
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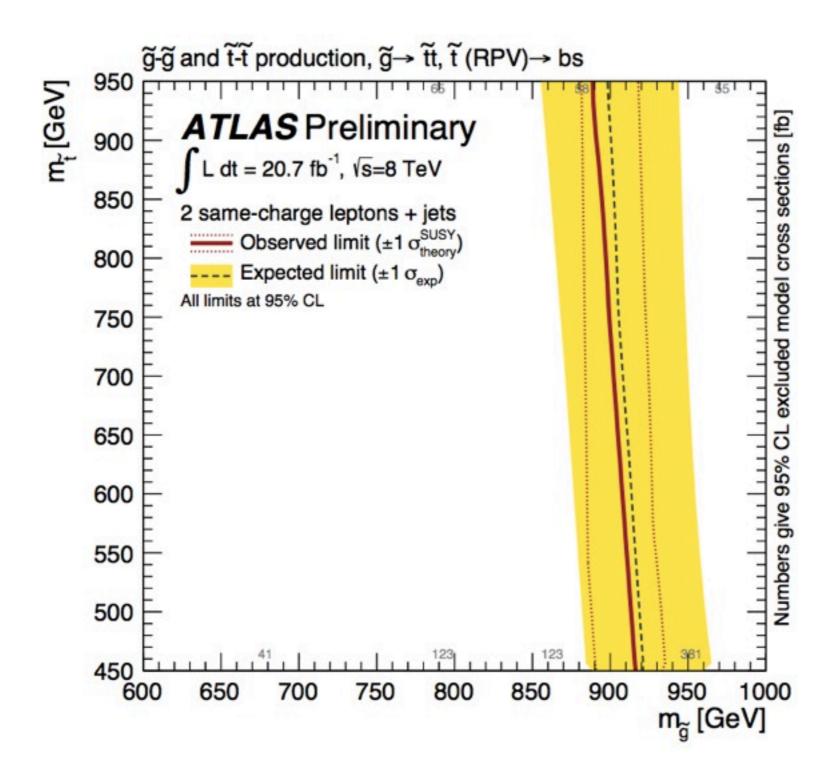
# notice Br(4t) < 0.25 in relevant scenario: 15-20% reduction of bound



Mahbubani '13

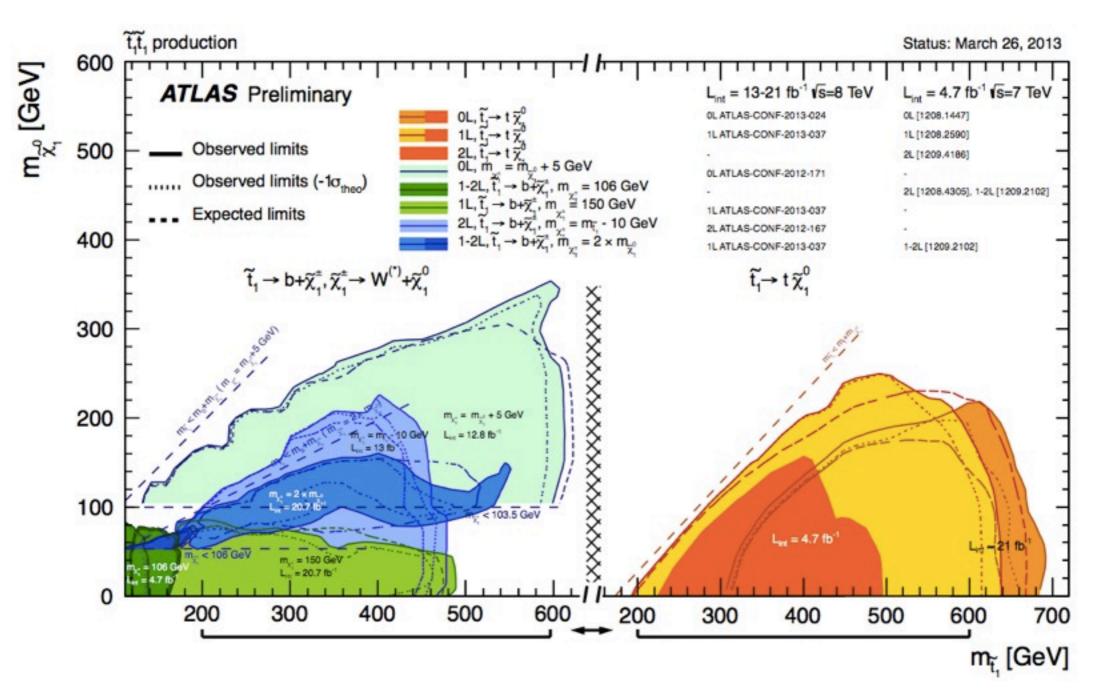
however the other channels (tttb, ttbb) add to same signal expect compensation, in the end bound should not change much

### what about some help from baryonic RPV?



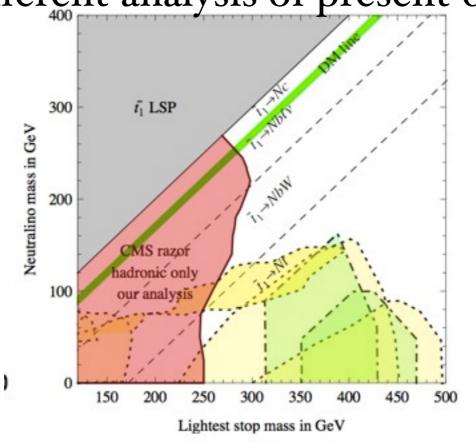
some relaxation but not much

# forget about majorana gluino and go to dirac: supersoft stop masses zoom on stops and higgsinos



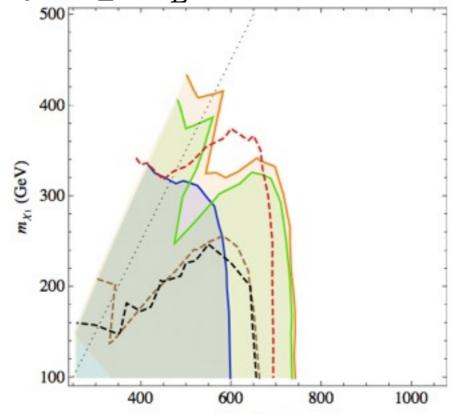
impressive but significant natural regions with squashed spectra remain

probably these regions can already be more significantly constrained by different analysis of present data



Delgado, Giudice, Isidori, Pierini, Strumia '13

also addition of  $\tilde{t}_R$ ,  $\tilde{t}_L$ ,  $\tilde{b}_L$  production should be considered

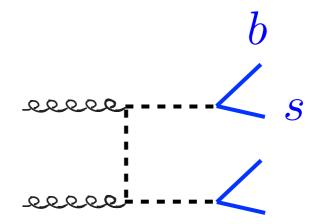


m: (GeV)

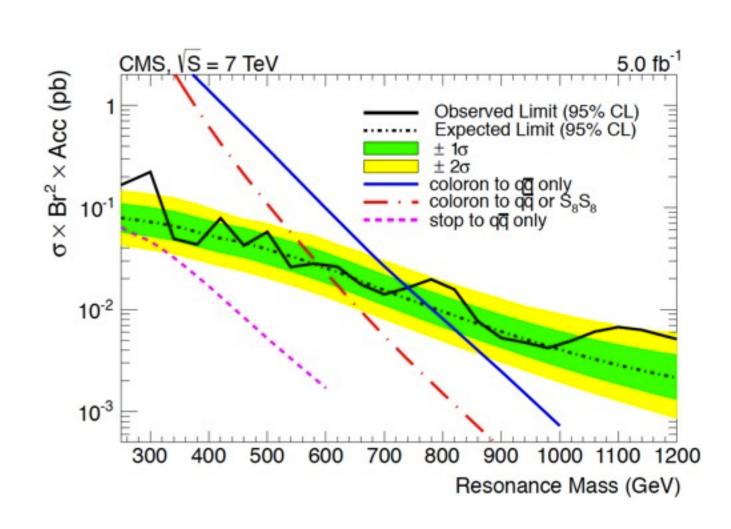
Kribs, Martin, Menon '13

#### what about RPV decaying

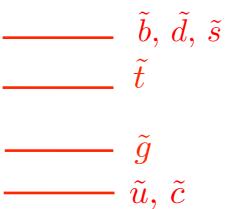
$$ilde{t}_R$$
 ,  $ilde{t}_L$  ,  $ilde{b}_L$  ?

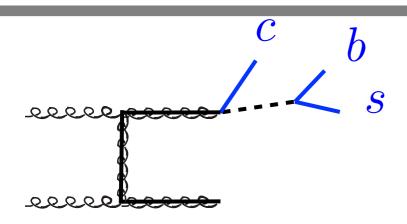


CMS 7 TeV data below 250 Gev wait for upcoming 8 TeV



what about topless gluino decays?

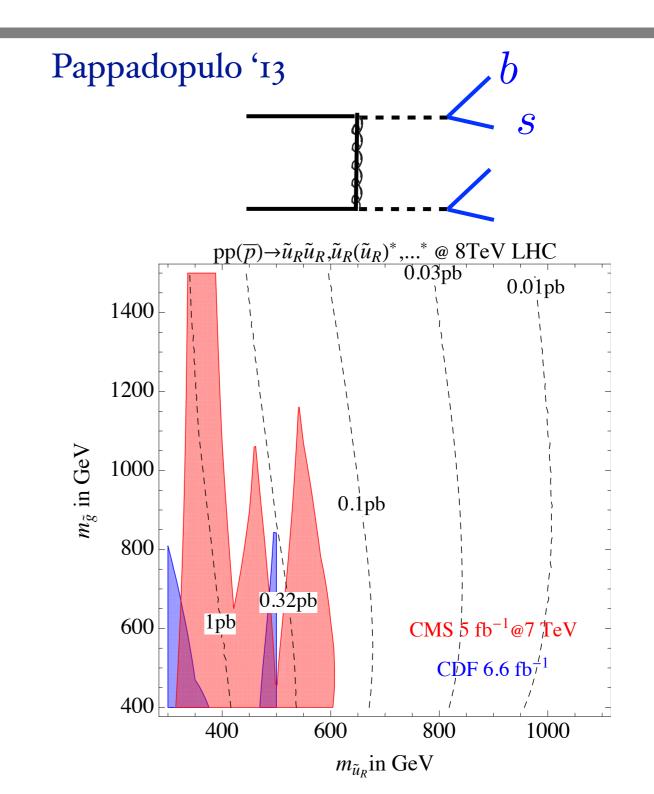




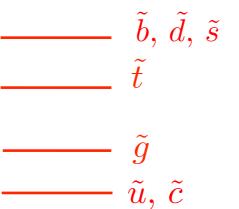
ATLAS, 7 TeV 4.6 fb<sup>-1</sup>  $m_{\tilde{g}} > 666 \, \mathrm{GeV}$ 

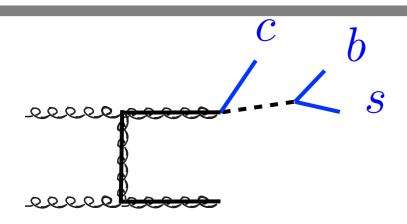
but bound relaxed if bino enters decays chain

 $\tilde{g}\tilde{g} \to 10 \mathrm{jets}$ 



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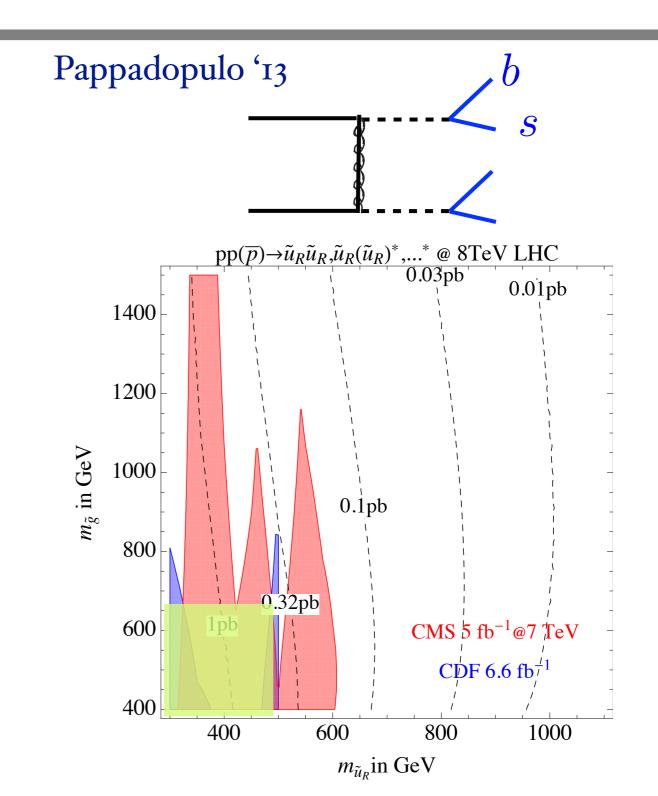




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$$m_{\tilde{g}} > 666 \, \mathrm{GeV}$$

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ATLAS & CMS are hunting down Supersymmetry in nooks and crannies

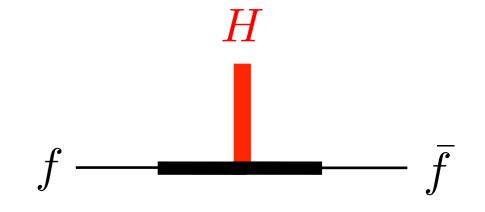
Looking forward to next run and to even more clever analyses

Too early tough to fully overthrow naturalness in favor
of scenarios like, for instance, mini-split SUSY
but pressure is clearly building up

anyway, wait to hear Nima

## Compositeness

Flavor: only option is partial compositeness



without any additional symmetry

$$m_* \gtrsim 10 - 20 \,\mathrm{TeV}$$

$$m_* \gtrsim 40 \, \text{TeV}$$

$$\mu \to e \gamma$$

$$m_* \gtrsim 150 \,\mathrm{TeV}$$

with combination of SU(2)'s and SU(3) can bring scale down to ~TeV

Redi '12 Babrbieri, Buttazzo, Sala, Straub, Tesi '12

$$V(H) = \begin{array}{c} t_L \\ \longleftarrow \\ \text{partners} \end{array} \quad (B_{-1/3}, \quad T_{2/3}, \quad X_{5/3}, \quad \Xi_{8/3}, \ldots)$$

## expectations from naturalness

generic 
$$V \equiv V(H/f)$$



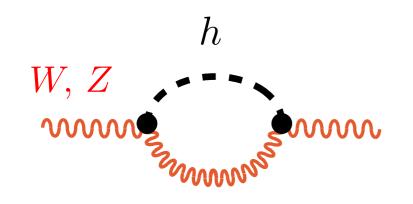
$$\frac{v^2}{f^2} = O(1)$$

$$m_h \simeq 125 \, {\rm GeV}$$



$$\left(\frac{400\,\text{GeV}}{m_T}\right)^2 = O(1)$$

#### Electroweak Precision Tests



$$\Delta \widehat{S}, \, \Delta \widehat{T}$$

$$W,Z$$
  $S$   $\Delta \widehat{S},\Delta \widehat{T}$   $v^2 \lesssim 0.05$  optimist  $f^2 \lesssim 0.05$  pessimist

$$W,Z$$
  $V$   $\Delta \widehat{S} \sim \frac{m_W^2}{m_V^2}$   $m_V^2 \gtrsim 2-3\,\mathrm{TeV}$ 

$$\Delta \widehat{S} \sim \frac{m_W^2}{m_V^2}$$

$$m_V^2 \gtrsim 2 - 3 \, {\rm TeV}$$

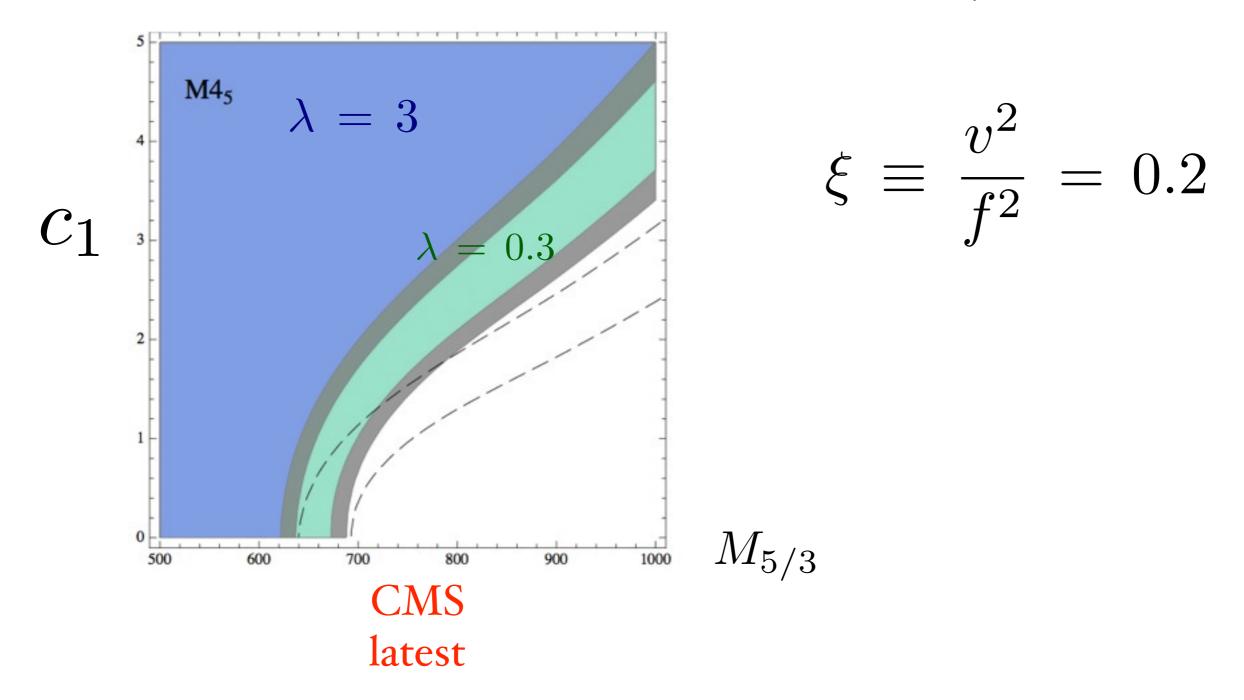
EWPT already imply some tuning

technical naturalness demands structural complexity (  $m_V \gg m_T$  )

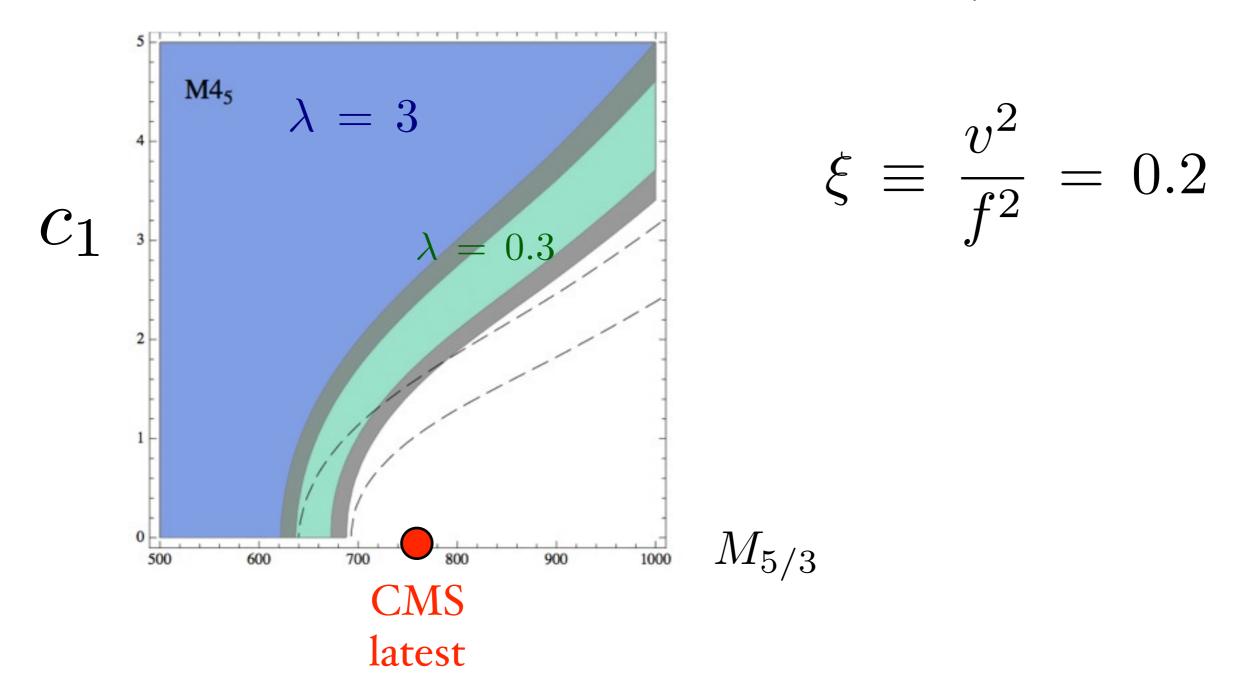
## Higgs couplings

$$\frac{\delta g}{g_{\scriptscriptstyle SM}} \sim \frac{v^2}{f^2} \lesssim 0.2$$
 from EWPT

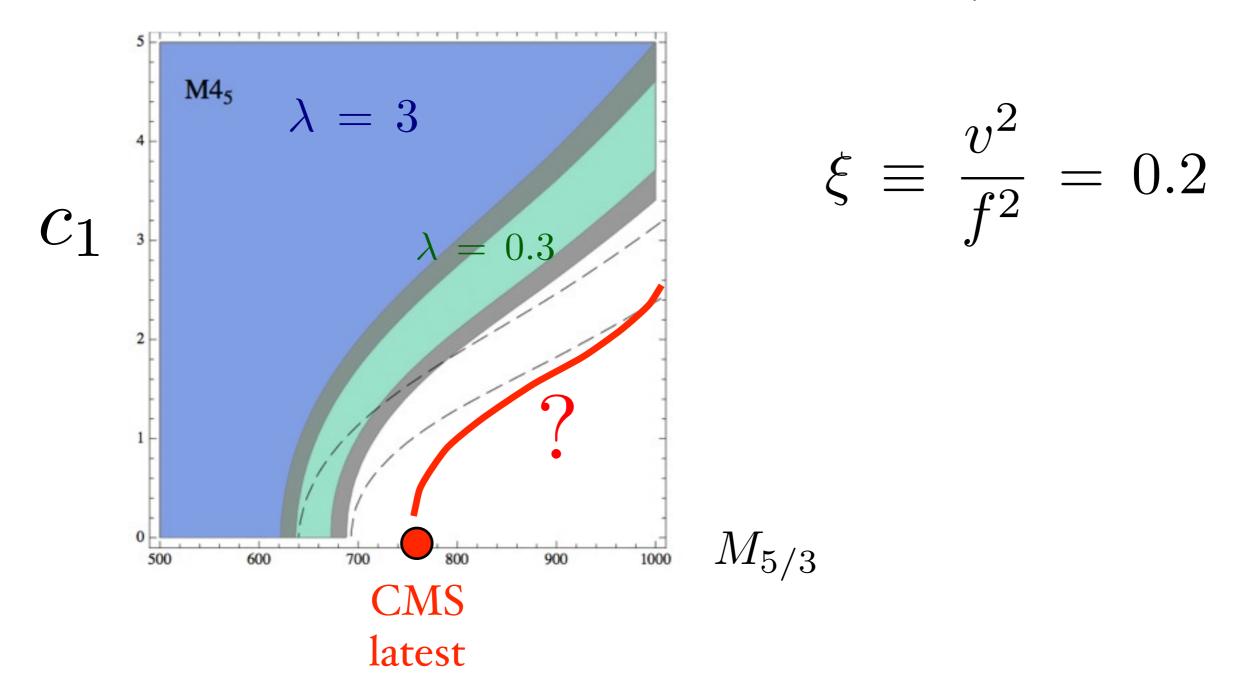
wait for more integrated luminosity to break new grounds



this is a very significant direct 'test' of naturalness but result not fully unexpected in view of LEP/etc...



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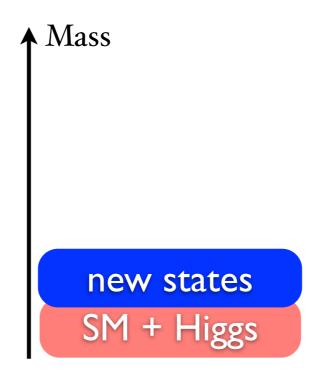
#### Vector Resonances

$$\int_{ar{q}}^{Q} V = rac{g_W^2}{g_V} \ll g_W \qquad \bigvee_{ar{w}_L}^{V \otimes v_W} = g_V$$

$$\sigma \sim 10 \, \text{fb} \, \left(\frac{3}{g_V}\right)^2 \left(\frac{2 \, \text{TeV}}{m_V}\right)^6$$

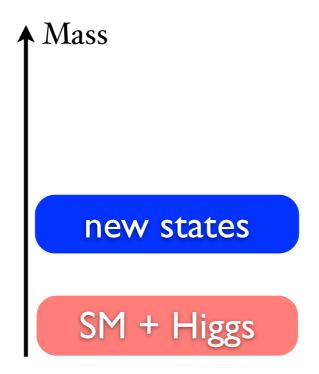
- resonances couple 'superweakly' to light fermions
- significantly different from weakly coupled W' and Z'
- LHC bound still below 2 TeV for  $g_V > 3$
- wait for LHC13 to break new ground

# Higgzoology



Can use effective lagrangian to describe deviations from SM

= simple parametrization encompassing a large class of models

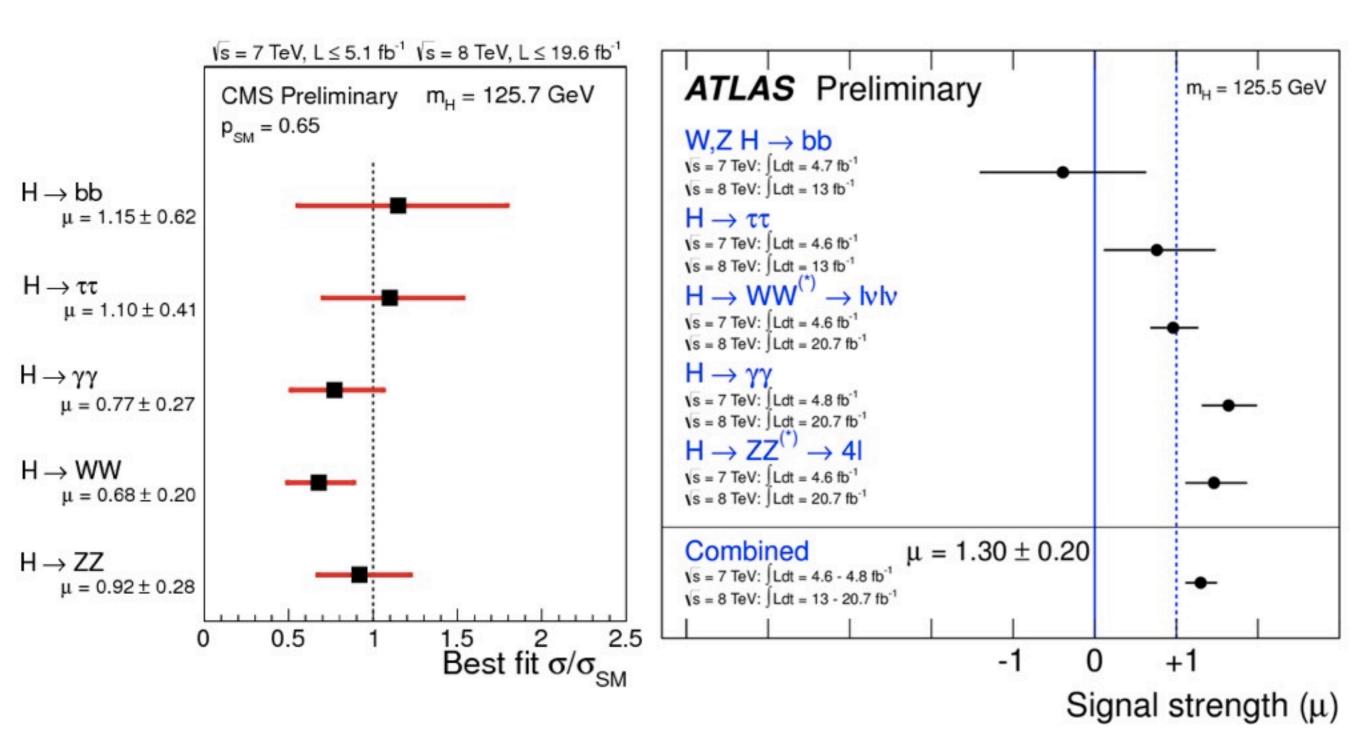


Can use effective lagrangian to describe deviations from SM

= simple parametrization encompassing a large class of models

$$c_V imes rac{2m_V^2}{v}$$
  $c_i imes rac{m_i}{v}$   $c_i imes rac{m_i}{v}$ 

$$c_V>1$$
 — only if  $\exists$  scalar of electric charge 2  $c_b>1,\ c_t<1$  — MSSM  $c_b<1,\ c_t>1$  — NMSSM dominated quartic  $c_V,c_f<1$   $c_g,c_\gamma\sim0$  — Composite Higgs



No clear trend of deviation from SM Compatible with SM within 30%



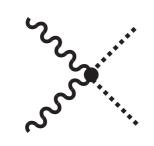


### The precision frontier

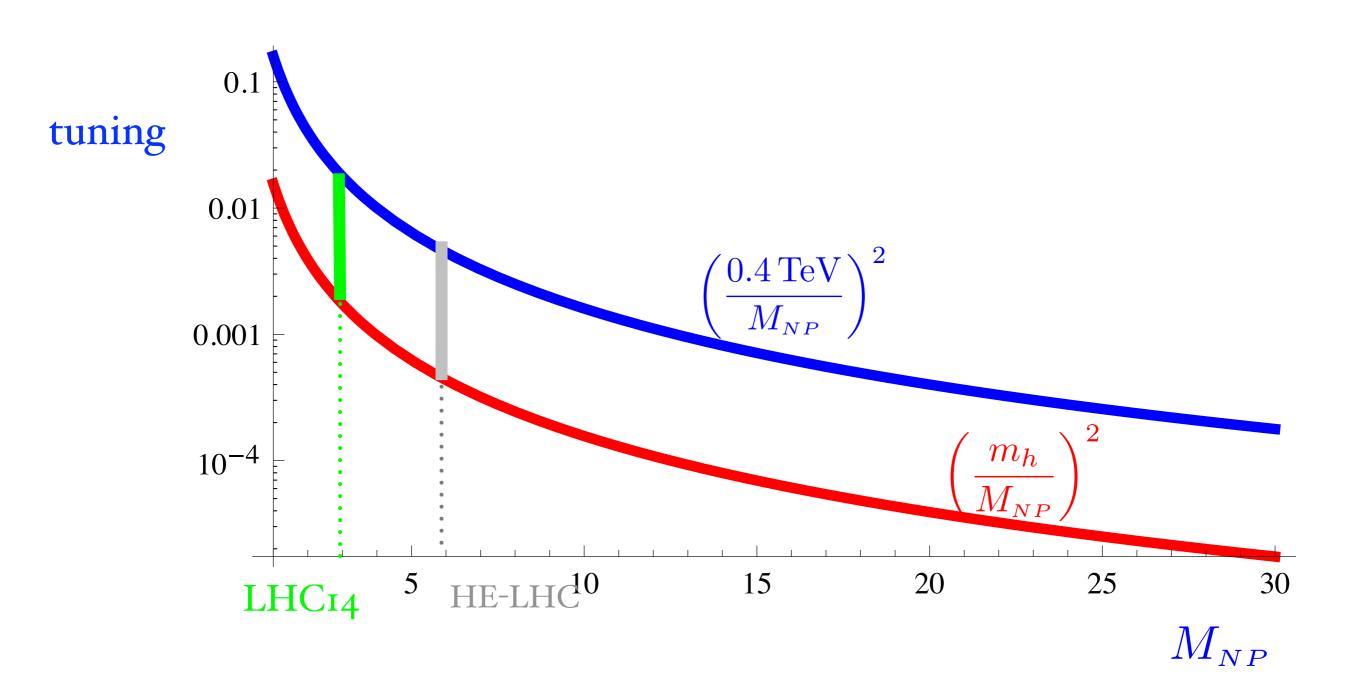
single Higgs measures 
$$\frac{g_{NP}^2 v^2}{M_{NP}^2} \lesssim$$
 fine tuning

precision Higgs physics breaks ground in the test of naturalness though it can only give us indirect clues

semidirect clues



### The energy frontier



# Summary

- discovery of Higgs boson with rates in agreement with SM within the present precision of 30% is far from unexpected: it is basically a corollary of results of LEP/Tevatron/B-factories
- direct searches are directly pushing several scenarios (especially in SUSY) into 1% tuning grounds, though that is basically were the simplest (and maybe nicest) models were already expected by indirect reasoning (ex, MSSM with flavor universal soft terms)
- LHC searches are now also putting significant pressure on cleverly natural models, though regions with moderate tuning are not ruled out yet
- Refined analyses & LHC13 can break grounds on those regions and perform a comprehensive test of naturalness
- •HL-LHC will break grounds in EW physics by testing Higgs rates at, so it seems, 5%