

Shift-multiplexed holographic memory using the two-lambda method

George Barbastathis and Demetri Psaltis

Department of Electrical Engineering, Mail Stop 136-93, California Institute of Technology, Pasadena, California 91125

Received October 18, 1995

We present theoretical and experimental results on the application of the two-lambda method for prolonged readout of holographic memories to shift multiplexing implemented with a spherical-wave reference beam.

Key words: Shift multiplexing, two-lambda method, surface storage density. © 1996 Optical Society of America

Shift multiplexing^{1,2} is a method for holographic storage that uses nonplanar reference beams. A volume hologram recorded with a spherical-wave reference becomes Bragg mismatched when it is translated with respect to the readout beam. This effect can be used to multiplex holograms in the same recording medium.^{1,3} The required amount of translation δ between holograms is called shift selectivity and is typically of the order of a few to a few hundred micrometers, depending on the geometry, the material thickness, and the distance of the focus of the spherical wave to the hologram. Shift multiplexing is particularly convenient for holographic three-dimensional disks,⁴ because disk rotation naturally produces the required shift.

When photorefractive crystals are used to implement rewritable shift-multiplexed memories, an issue of concern is the undesired erasure of the recorded holograms during readout. Techniques have been developed for fixing holograms thermally⁵ or electrically.^{6,7} Alternatively, holograms can be sustained by periodic refreshing.⁸ Here we consider the two-lambda method, in which the recording and readout wavelengths are different.⁹ The allowable readout time is prolonged if the readout wavelength is selected such that the absorption is lower and hence erasure is reduced. Because the Bragg matching condition at the new wavelength is modified, only a portion of the angular spectrum of the recorded holograms can be Bragg matched with a single plane wave. It is possible, however, to reconstruct in its entirety a single hologram recorded with a plane wave by reading out the hologram with a spherical beam.¹⁰ In this Letter we show how the two-lambda technique applies to shift-multiplexed memories, in which holograms are recorded and read out with spherical-wave references.

We use the subscript 1 for quantities associated with the recording wavelength λ_1 and the subscript 2 for the readout wavelength λ_2 . The geometry is shown in Fig. 1. The holographic material occupies the region $-\infty < x, y < +\infty$, $|z| < L/2$. The recording reference beam (at wavelength λ_1) is a spherical wave originating from $(0, 0, -z_1)$. For now, we consider the signal beam to be a single plane-wave component propagating at angle $u_1 \approx \sin^{-1}(\lambda_1/L)$ (paraxial approximation) with respect to the hologram normal. The readout spherical wave (at wavelength λ_2) originates at $(\delta, 0, -z_2)$. Un-

der the Born, paraxial, and constant modulation depth approximations^{2,11} the diffracted field can be calculated analytically. The reconstruction is a plane wave if

$$\lambda_1 z_1 = \lambda_2 z_2. \quad (1)$$

This change in focal distance compensates for the curvature mismatch between the recording and readout reference wave fronts. The plane wave hologram is Bragg matched at the readout wavelength if it is translated relative to the readout reference by

$$\delta_B = \frac{1}{2} \left(\frac{\lambda_2}{\lambda_1} - 1 \right) u_1 z_2. \quad (2)$$

If Eqs. (1) and (2) are both satisfied, the reconstruction is a plane wave propagating at angle u_2 satisfying (in the paraxial approximation) the condition

$$u_2/u_1 = \lambda_2/\lambda_1. \quad (3)$$

Once the hologram is Bragg matched after shifting by δ_B , an additional translation by the shift Bragg selectivity

$$\delta_2 = \frac{\lambda_2 z_2}{u_2 L} = \frac{\lambda_1}{\lambda_2} \delta_1 \quad (4)$$

will make the reconstruction go away. A separate plane-wave hologram recorded at that shifted position

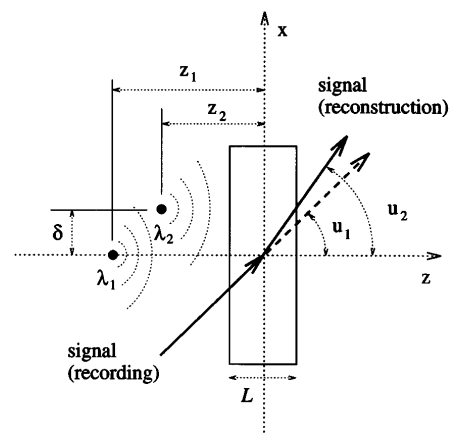


Fig. 1. Geometry for the two-lambda technique with shift multiplexing.

Table 1. Two-Lambda Equations Including Refraction and Dispersion

| | |
|--------------------------|---|
| Bragg selectivity | $\delta_1 = \frac{\lambda'_1[n_1 z'_1 - (n_1 - 1)L/2]}{u'_1 L} + \frac{\lambda'_1}{2(\text{NA})_1}$ |
| | $\delta_2 = \frac{\lambda'_2[n_2 z'_2 - (n_2 - 1)L/2]}{u'_2 L} + \frac{\lambda'_2}{2(\text{NA})_2}$ |
| Focusing condition | $\lambda'_1 \left[z'_1 - \left(1 - \frac{1}{n_1}\right) \frac{L}{2} \right] = \lambda'_2 \left[z'_2 - \left(1 - \frac{1}{n_2}\right) \frac{L}{2} \right]$ |
| Bragg matching condition | $\delta_B = \frac{1}{2} \left(\frac{n_1 \lambda'_2}{n_2 \lambda_1} - 1 \right) \frac{u'_1}{n_1} [n_2 z'_2 - (n_2 - 1)L/2] + \frac{\lambda'_2}{2(\text{NA})_2}$ |
| Reconstructed band | $\Delta u_2 = \frac{\lambda'_1}{\left(\frac{n_1 \lambda'_2}{n_2 \lambda_1} - 1 \right) u'_{1c} L}$ |

can be observed without interference. In Eq. (4), $\delta_1 = \lambda_1 z_1 / (u_1 L)$ is the shift selectivity at the recording wavelength.² Note that if $\lambda_2 > \lambda_1$, then δ_2 is smaller than δ_1 . Because holograms are recorded δ_2 apart, packing is denser than for the case in which readout is intended to be at wavelength λ_1 .

In the analysis so far we have neglected refraction and dispersion effects and the finite numerical apertures $(\text{NA})_1$ and $(\text{NA})_2$ of the spherical waves. The properly modified formulas are given in Table 1, where the primed quantities are measured in air and n_1 and n_2 are the refractive indices. In general, $(\text{NA})_2 < (\text{NA})_1$ because the required shift δ_B introduces edge effects.

We verified these theoretical predictions experimentally, using a thin Fe-doped LiNbO₃ crystal grown by Deltronics. The experimental parameters were $\lambda'_1 = 488$ nm, $\lambda'_2 = 632.8$ nm, $u'_1 \approx 40^\circ = 0.698$ rad, $n_1 = 2.3533$, $n_2 = 2.2935$, $L = 250$ μm , and $(\text{NA})_1 \approx 0.1$. The quoted values of refractive index were calculated by interpolation of the data given in Ref. 12. This experiment was done with a plane-wave signal beam. The measured values of δ_B , δ_1 , and δ_2 are plotted in Fig. 2 together with the theoretical predictions derived from the formulas of Table 1. $(\text{NA})_2$ was estimated to be 0.014. We made the measurements for the λ_2 readout by first refocusing the readout beam according to the theoretical prediction for z'_2 .

An information-bearing signal occupies a finite angular bandwidth around the carrier u_{1c} . Because δ_B [Eq. (2)] is different for each spatial frequency component, the entire hologram cannot be Bragg matched at the same time. The reconstruction consists of the portion of the angular spectrum lying in the range

$$\Delta u_2 = \pm \frac{2\lambda_1}{[(\lambda_2/\lambda_1) - 1]u_{1c}L}. \quad (5)$$

For Fourier-transform holograms, Eq. (5) implies that a slice of size $\Delta x = 2\Delta u_2 F$ (where F is the focal length of the lens used for the Fourier transformation) is reconstructed. For image-plane holograms the re-

construction is low-pass filtered with a cutoff $\Delta w = \Delta u_2 / \lambda_2$. The modified version of Eq. (5) accounting for dispersion is given in Table 1.

We used the same experimental setup to record a single Fourier-plane hologram (focal length $F = 20$ cm) of a transparency with $z'_1 = 22.5$ mm. The reconstruction at wavelength λ_1 is shown in Fig. 3(a). Two reconstructions at λ_2 (obtained by changing δ_B) are shown in Figs. 3(b) and 3(c). As predicted, only a

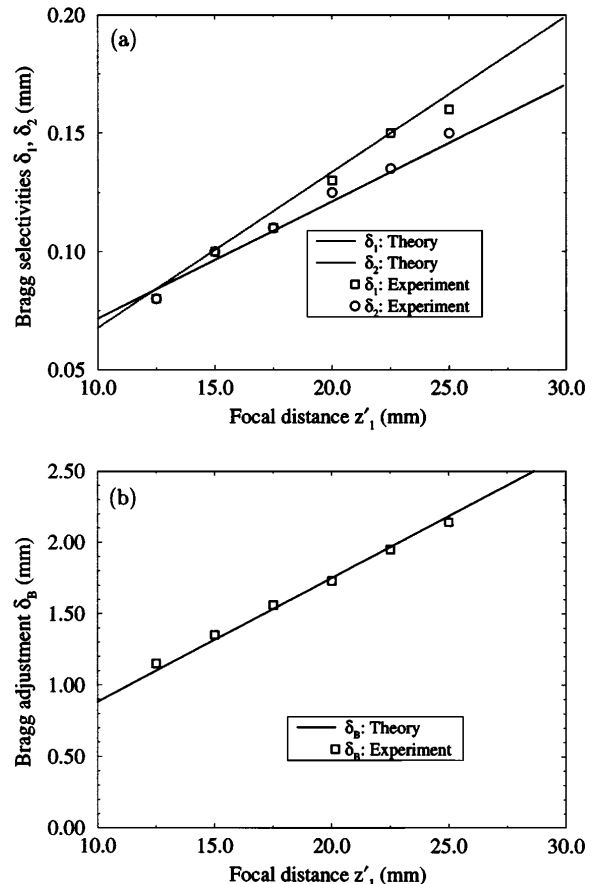


Fig. 2. Experimental results for the Bragg matching and selectivity properties of the two-lambda method applied to shift-multiplexed holograms.

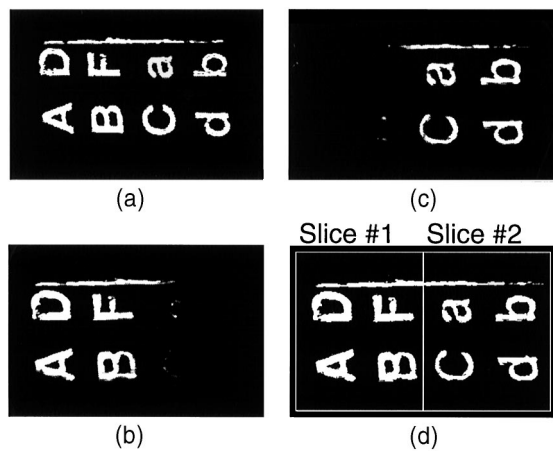


Fig. 3. Hologram reconstructions obtained with the two-lambda method.

slice of the stored image is obtained at one time. The width of the slice is 0.7 mm, in agreement with Eq. (5). If we continuously shift the medium, a sliding window of the stored image will appear on the CCD. In addition, the reconstruction as viewed through this sliding window is also shifting because of the motion of the medium. A time-delay-and-integrate detector array can compensate for this motion and integrate the response to produce a complete, unblurred image. Such a detector was simulated in software to produce the complete reconstruction shown in Fig. 3(d). Alternatively, we can set the size of the recorded images in the x direction equal to the width of the reconstructed slice. Either method results in the complete reconstruction of the stored images at the expense of reduced storage density.

The surface storage density of a shift-multiplexed memory can be expressed as

$$\mathcal{D} = (N_{px}N_{py})/(s_1\delta_2), \quad (6)$$

where N_{px} and N_{py} are the number of spatial light modulator pixels in the x and y directions, respectively, and s_1 is the transverse size of the signal beam. For Fourier-plane recording we can relate N_{px} to the geometry of the holographic system by imposing the requirement that the angular bandwidth of the stored holograms (and hence the hologram thickness) be matched to the width of the allowable reconstructed angles (the size of the spatial light modulator). Substituting Eqs. (5) and (6), we obtain

$$\mathcal{D} = \frac{2N_{py}}{\left(\frac{\lambda_2}{\lambda_1} - 1\right)\lambda_1 z_1}. \quad (7)$$

Equation (7) is valid only if the spatial light modulator's size is large enough to accommodate the entire strip. Notice that when $\lambda_2 = \lambda_1$ the strip is infinite, and therefore Eq. (7) does not hold. As an example, if $N_{py} = 1000$, $z_1' = 1.3$ mm, $\lambda_1' = 488$ nm, and $\lambda_2' = 633$ nm, then $\mathcal{D} = 10.6$ bits/ μm^2 . A similar derivation can be carried out for the density of the image-plane geometry, and it leads to the same expression as Eq. (7).

We have presented the theory of the two-lambda technique applied to spherical reference wave holograms and verified the results experimentally. Shift multiplexing is also achieved with a one- or two-dimensional fan of plane waves (array method¹). We are currently working on the application of two-lambda readout to array-multiplexed holograms.

This research was supported by the U.S. Air Force Office of Scientific Research under grant F49620-92-J-0400. We thank Geoffrey Burr, Ernest Chuang, Allen Pu, Jean-Jacques Drolet, and Fai Mok for helpful discussions and Yayun Liu for technical assistance.

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