CHAOS, CRITICALITY, AND COMPUTATION IN RECURRENT NETWORKS

MORITZ HELIAS

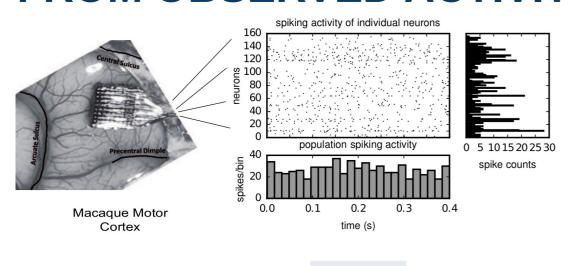
2023-06-15 LAUSANNE

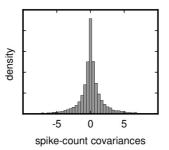
COMPUTATIONAL AND SYSTEMS NEUROSCIENCE (INM-6)
THEORETICAL NEUROSCIENCE (IAS-6)
FACULTY OF PHYSICS, RWTH AACHEN UNIVERSITY



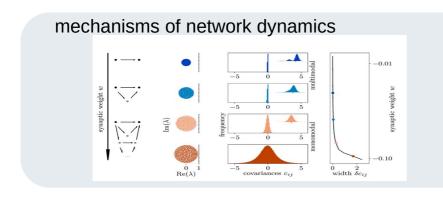


INFERENCE OF COLLECTIVE NETWORK DYNAMICS FROM OBSERVED ACTIVITY

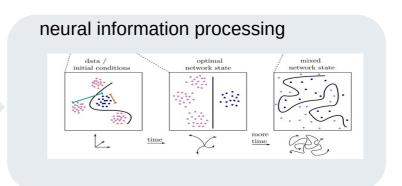




extract



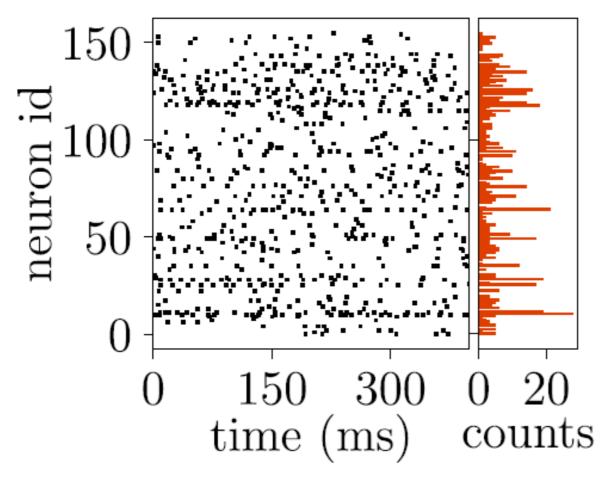
implement

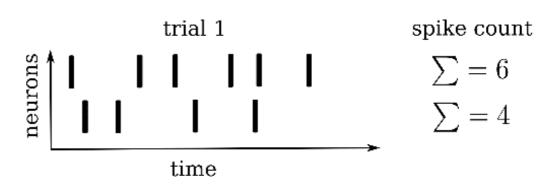


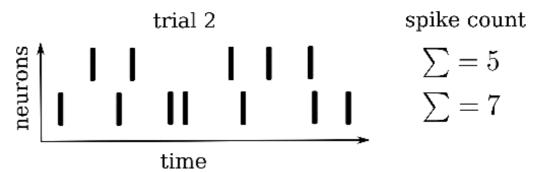




COLLECTIVE DYNAMICS - CORRELATIONS





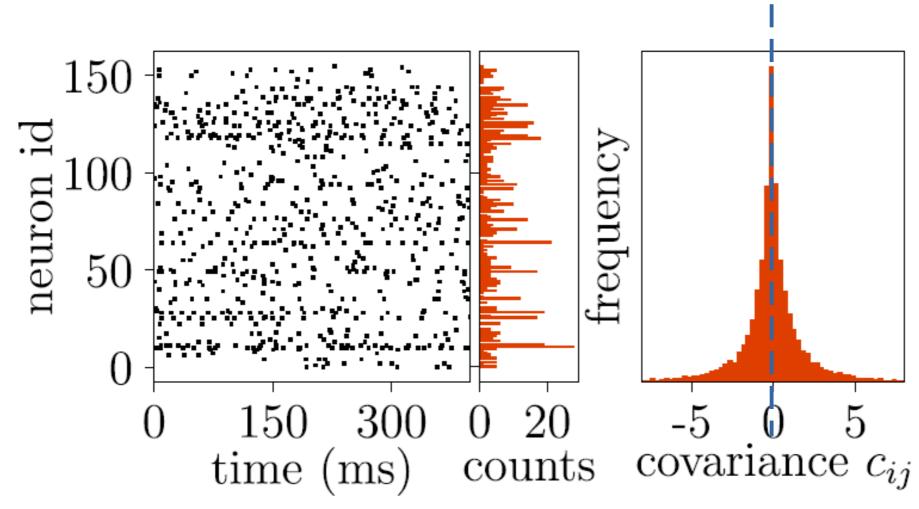


(Brochier et al. 2018)





SMALL AVERAGE CORRELATIONS

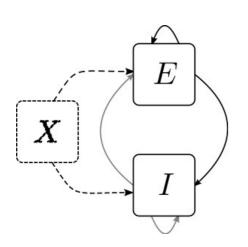


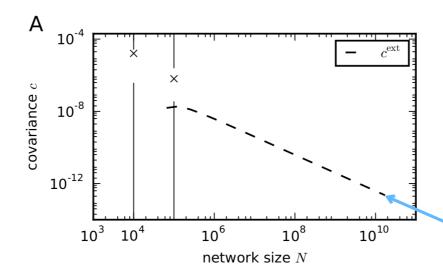
(Brochier et al. 2018)





SMALL AVERAGE CORRELATIONS – BALANCED STATE





Finite size-theory of **average** fluctuations Helias et al. 2014

$$c_{\alpha\beta} = \langle c_{ij} \rangle_{i \in \alpha, j \in \beta}$$
 $\alpha, \beta \in \{E, I, X\}$

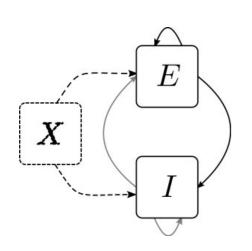
$$\begin{pmatrix} c_{EE} \\ c_{EI} \\ c_{II} \end{pmatrix} = \mathbf{c}_{\mathrm{int}}(a_E, a_I) + \mathbf{c}_{\mathrm{ext}}(a_x)$$

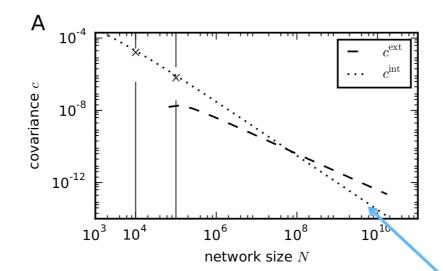
RWTHAACHEN LINIVERSITY



Renart et al. 2010

SMALL AVERAGE CORRELATIONS – BALANCED STATE





$$c_{
m ext} \propto N^{-1}$$
 $c_{
m int} \propto N^{-\frac{3}{2}}$

Finite size-theory of **average** fluctuations Helias et al. 2014

$$c_{\alpha\beta} = \langle c_{ij} \rangle_{i \in \alpha, j \in \beta}$$
 $\alpha, \beta \in \{E, I, X\}$

$$\begin{pmatrix} c_{EE} \\ c_{EI} \\ c_{II} \end{pmatrix} = \mathbf{c}_{\mathrm{int}}(a_E, a_I) + \mathbf{c}_{\mathrm{ext}}(a_x)$$

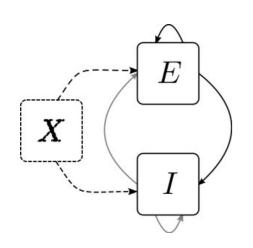
Renart et al. 2010

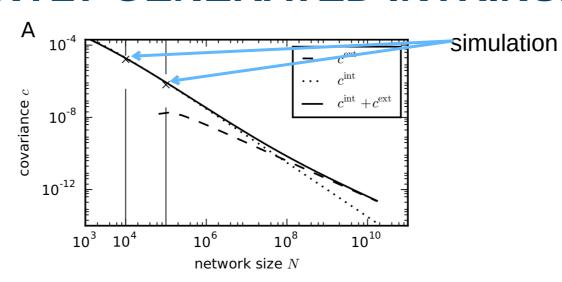




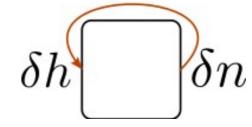
AVERAGE CORRELATIONS

- PREDOMINANTLY GENERATED INTRINSICALLY





explanation: negative feedback by inhibition



Finite size-theory of **average** fluctuations Helias et al. 2014

$$c_{\alpha\beta} = \langle c_{ij} \rangle_{i \in \alpha, j \in \beta}$$
 $\alpha, \beta \in \{E, I, X\}$

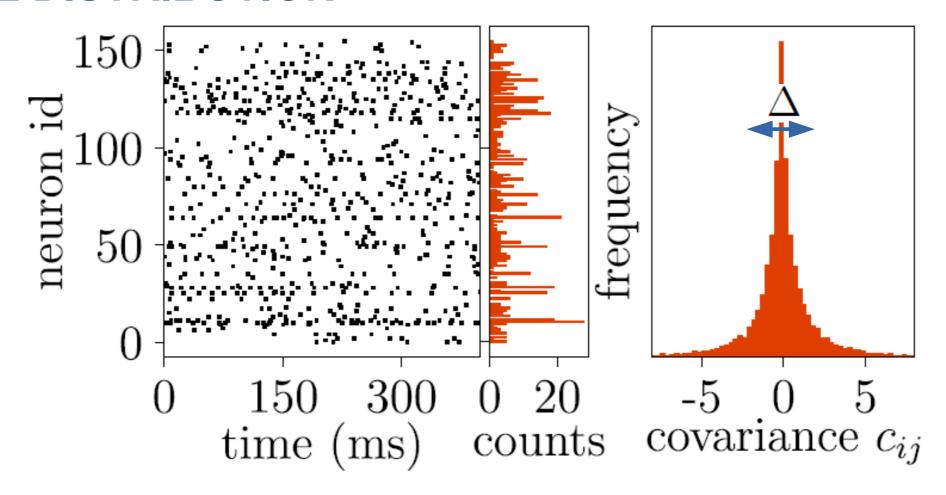
$$\left(\begin{array}{c} c_{EE} \\ c_{EI} \\ c_{II} \end{array} \right)$$
 intrinsic externally-driven $= \mathbf{c}_{\mathrm{int}}(a_E, a_I) + \mathbf{c}_{\mathrm{ext}}(a_x)$

RWTHAACHEN UNIVERSITY



Renart et al. 2010

WIDE DISTRIBUTION









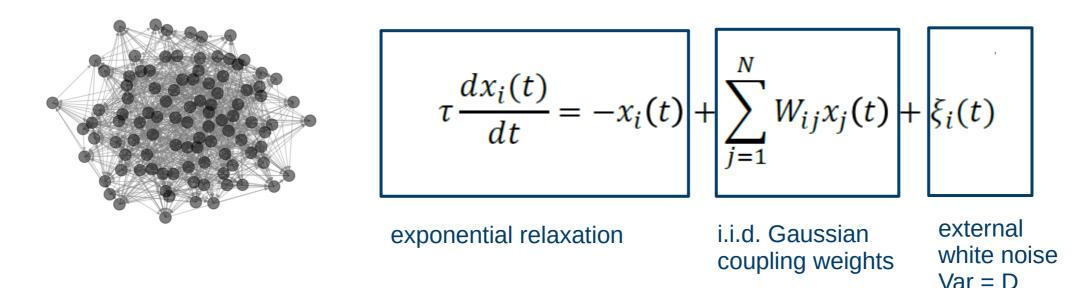
SIGNATURES OF CRITICAL STATES IN MOTOR CORTEX

DAVID DAHMEN





LINEAR NETWORK MODEL (LINEAR RESPONSE THEORY)



• Linear response theory captures fluctuations in asynchronous irregular brain states (Lindner et al. 2006, Pernice et al. 2011, Trousdale et al. 2012, Grytskyy et al. 2014)



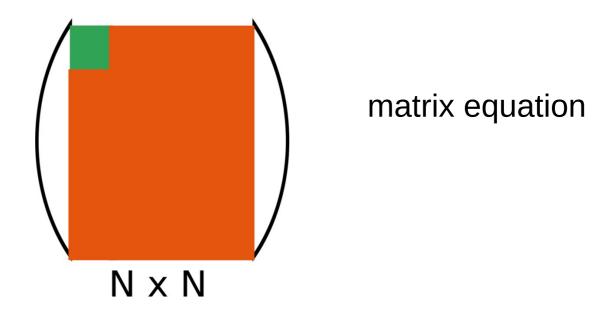


$$C = [1 - W]^{-1}D[1 - W]^{-T}$$



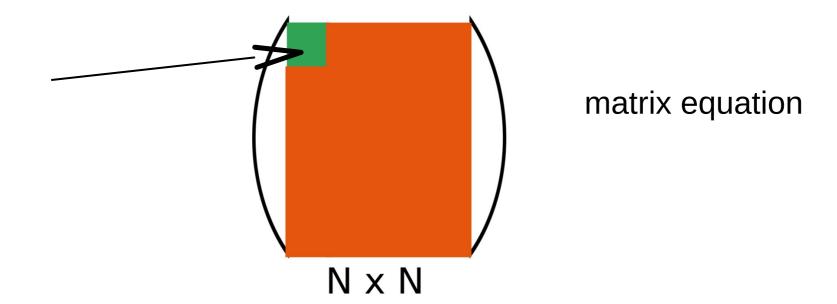
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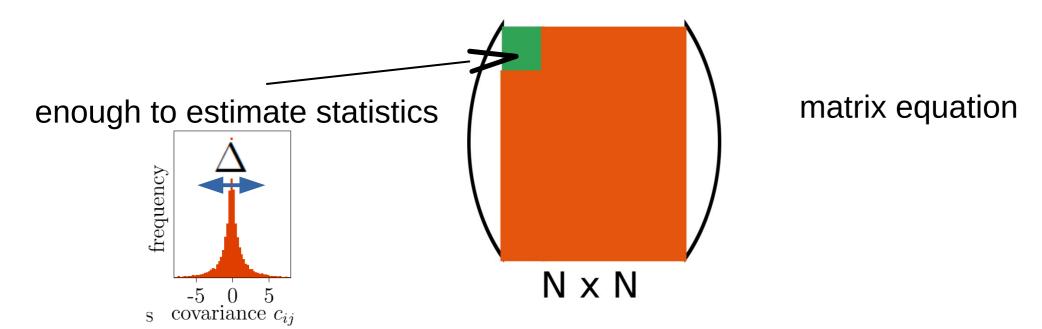


$$C = [1 - W]^{-1}D[1 - W]^{-T}$$





$$C = [1 - W]^{-1}D[1 - W]^{-T}$$





FIELD THEORETIC FORMULATION

$$C = [1 - W]^{-1}D[1 - W]^{-T}$$





FIELD THEORETIC FORMULATION

$$C = [1 - W]^{-1}D[1 - W]^{-T}$$

$$\tau \frac{dx_i(t)}{dt} = -x_i(t) + \sum_{j=1}^N W_{ij}x_j(t) + \xi_i(t)$$

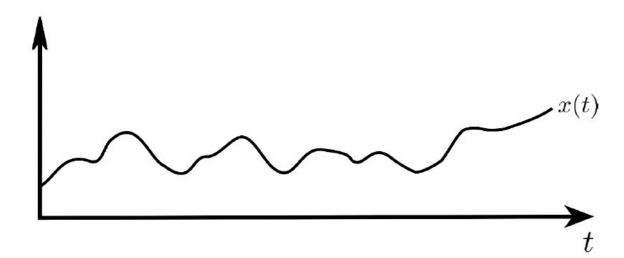




FIELD THEORETIC FORMULATION

$$C = [1 - W]^{-1}D[1 - W]^{-T}$$

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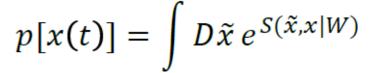
$$p[x(t)] = \int D\tilde{x} \, e^{S(\tilde{x}, x|W)}$$

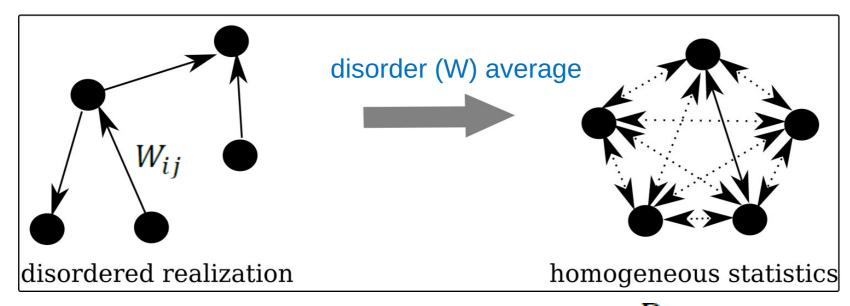
Martin Siggia Rose formalism Martin et al. 1973, DeDominicis 1975, Janssen 1976





ENSEMBLES OF NETWORKS





$$S_0(\widetilde{X}, X) = \widetilde{X}^{\mathrm{T}}(1 - \mu\{1\})X + \frac{D}{2}\widetilde{X}^{\mathrm{T}}\widetilde{X}$$

$$S_{\text{int}}(\widetilde{X}, X) = \frac{\sigma^2}{2N} \widetilde{X}^{\text{T}} \widetilde{X} X^{\text{T}} X$$





BEYOND MEAN-FIELD THEORY

$$S_0(\widetilde{\boldsymbol{X}}, \boldsymbol{X}) = \widetilde{\boldsymbol{X}}^{\mathrm{T}}(1 - \mu\{\mathbf{1}\})\boldsymbol{X} + \frac{D}{2}\widetilde{\boldsymbol{X}}^{\mathrm{T}}\widetilde{\boldsymbol{X}}$$

$$S_{\text{int}}(\widetilde{X}, X) = \frac{\sigma^2}{2N} \widetilde{X}^T \widetilde{X} X^T X$$
mean + fluctuation corrections

Result:

variance of entries of W spectral radius of connectivity W

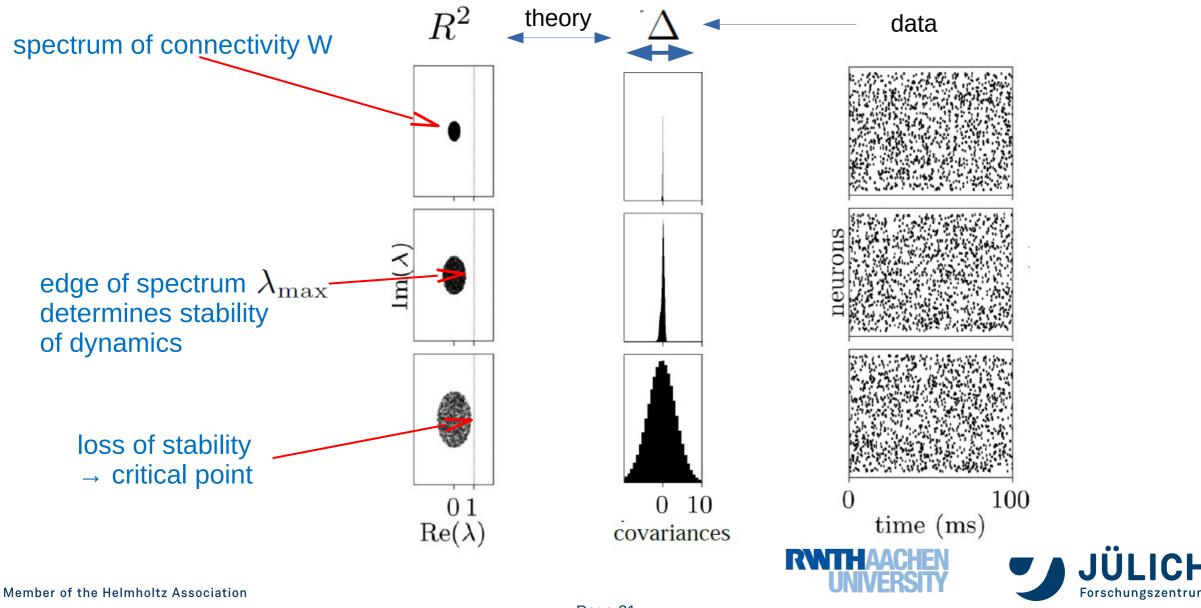
$$R^2 = 1 - \sqrt{\frac{1}{1 + N \, \Delta}}$$
 number of neurons

width of distribution of correlations

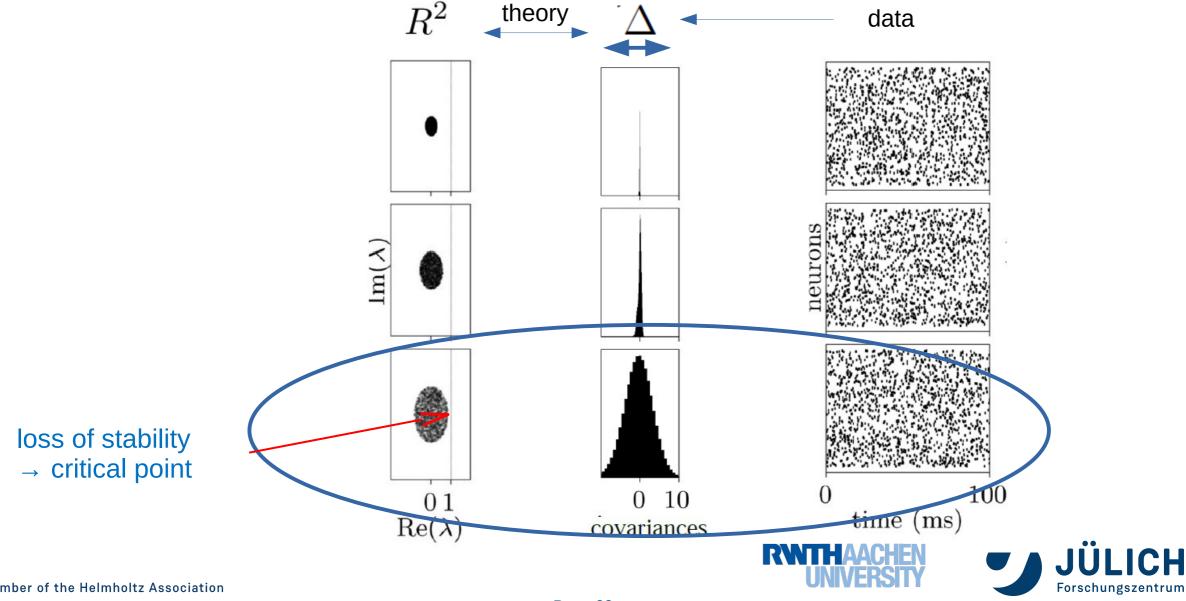




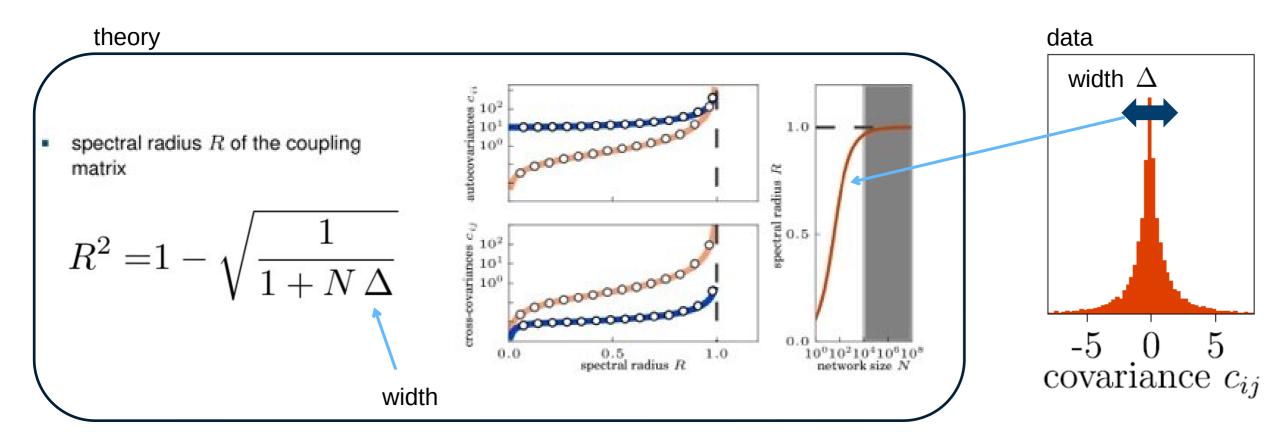
LARGE WIDTH IMPLIES CRITICALITY



LARGE WIDTH IMPLIES CRITICALITY



MOTOR CORTEX NEARLY UNSTABLE

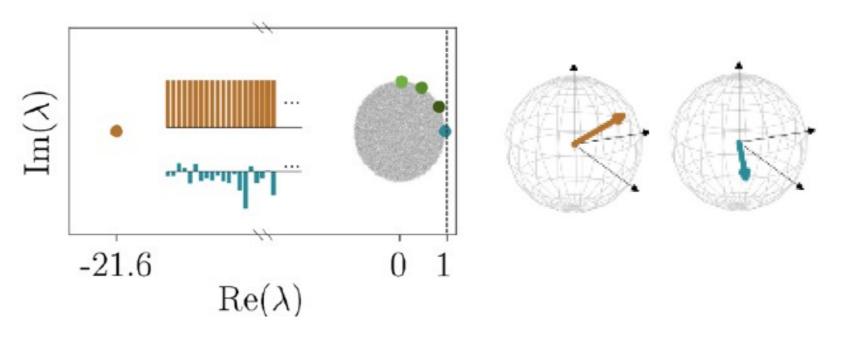


Motor cortex is operating close to critical point of linear instability R=1!



DYNAMICAL AND FUNCTIONAL CONSEQUENCES

- RICH REPERTOIRE OF DYNAMICAL MODES



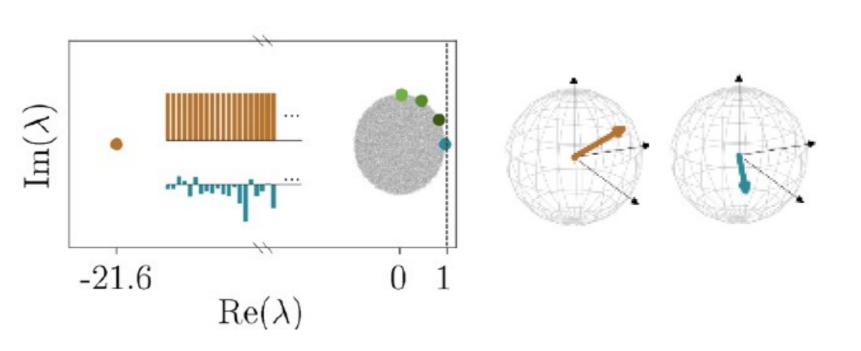
Dahmen et al., Second type of criticality in the brain uncovers rich multiple-neuron dynamics, PNAS, 2019



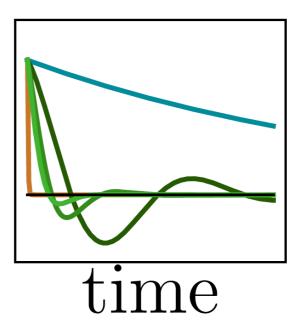


DYNAMICAL AND FUNCTIONAL CONSEQUENCES

- RICH REPERTOIRE OF DYNAMICAL MODES



$$v_{\alpha}(t) \sim e^{-t/\frac{\tau}{1-\lambda_{\alpha}}}$$



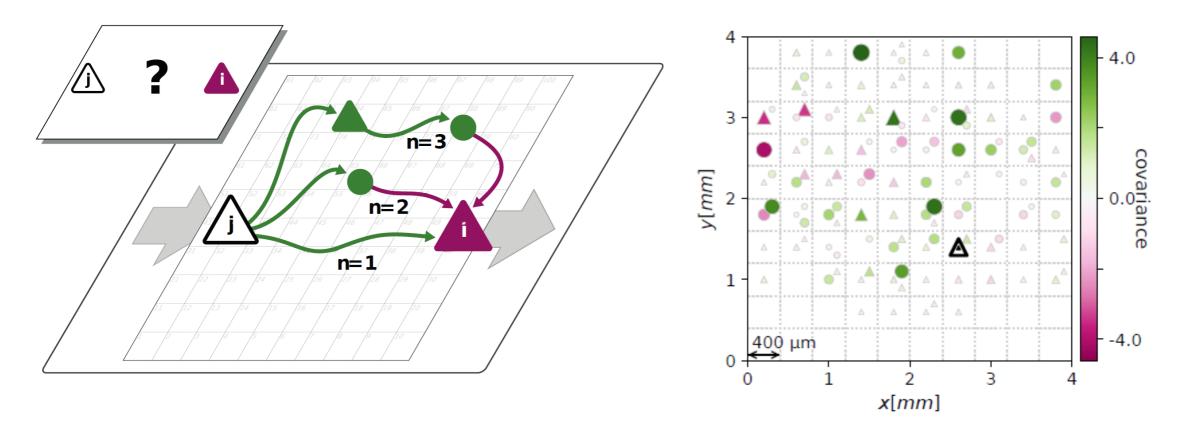
Dahmen et al., Second type of criticality in the brain uncovers rich multiple-neuron dynamics, PNAS, 2019





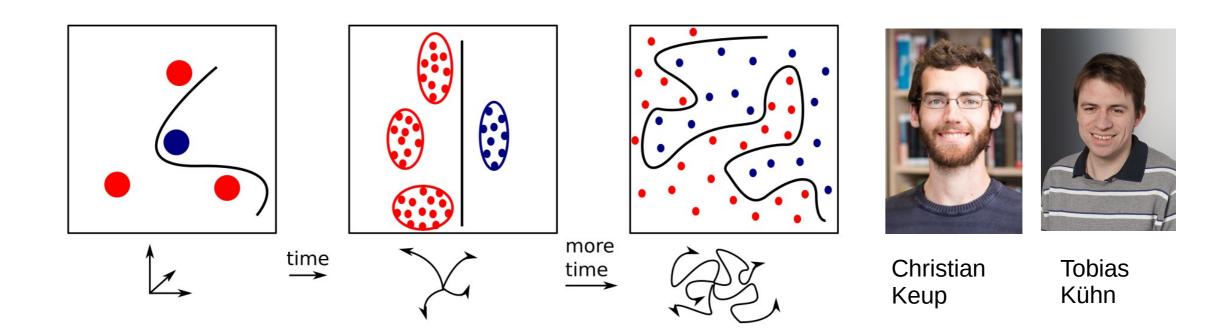
DYNAMICAL AND FUNCTIONAL CONSEQUENCES

-LONG-RANGE INTERACTIONS DESPITE SHORT-RANGE CONNECTIONS



Dahmen et al., Long-range coordination patterns in cortex change with behavioral context, elife, 2022





TRANSIENT CHAOTIC DIMENSIONALITY EXPANSION

CHRISTIAN KEUP, TOBIAS KÜHN, DAVID DAHMEN





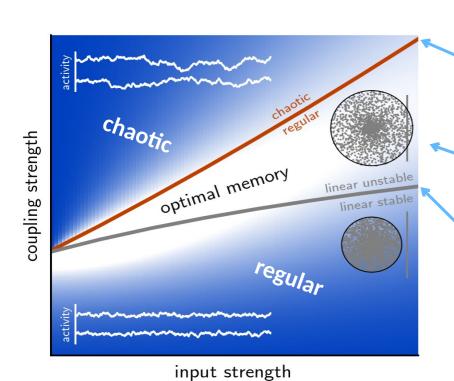
DRIVEN RANDOM RATE NETWORKS

- OPTIMAL MEMORY CLOSE TO CRITICALITY

coupling input

nonlinear network:

$$\tau \frac{dx_i(t)}{dt} = -x_i(t) + \sum_{j=1}^N J_{ij}\phi(x_j(t)) + \xi_i(t)$$



$$g^2\langle\phi^2\rangle > \langle x^2\rangle$$

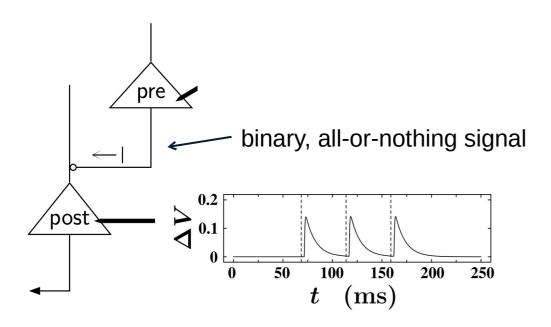
dynamical state between loss of linear stability and onset of chaos has optimal memory

linear instability
$$R^2 = g^2 \langle \phi'^2 \rangle > 1$$

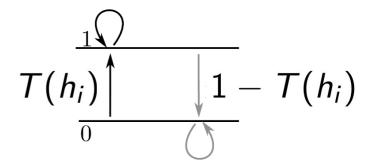
Schuecker et al., Optimal Sequence Memory in Driven Random Networks, PRX, 2018

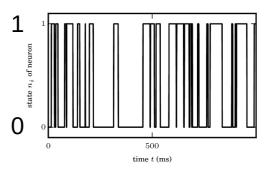
SPIKING INTERACTION: ABSTRACTION AS BINARY

Taking into account discrete coupling





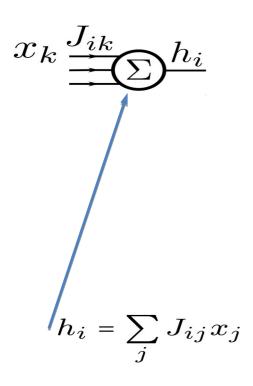




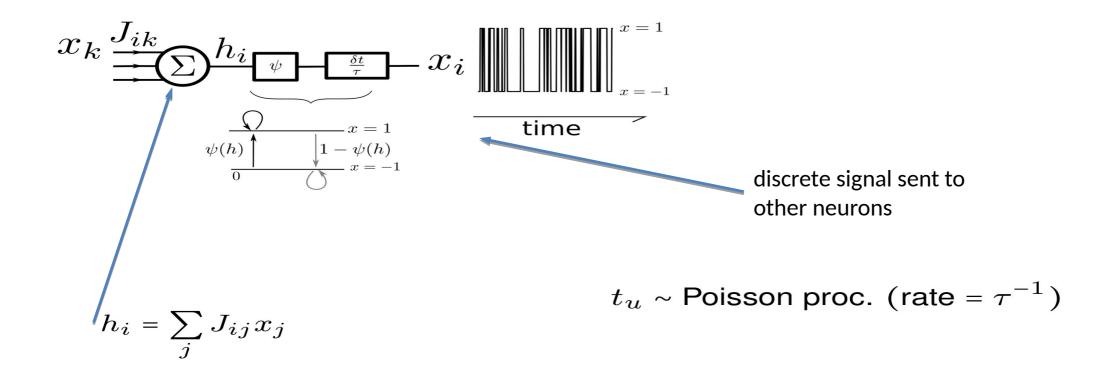


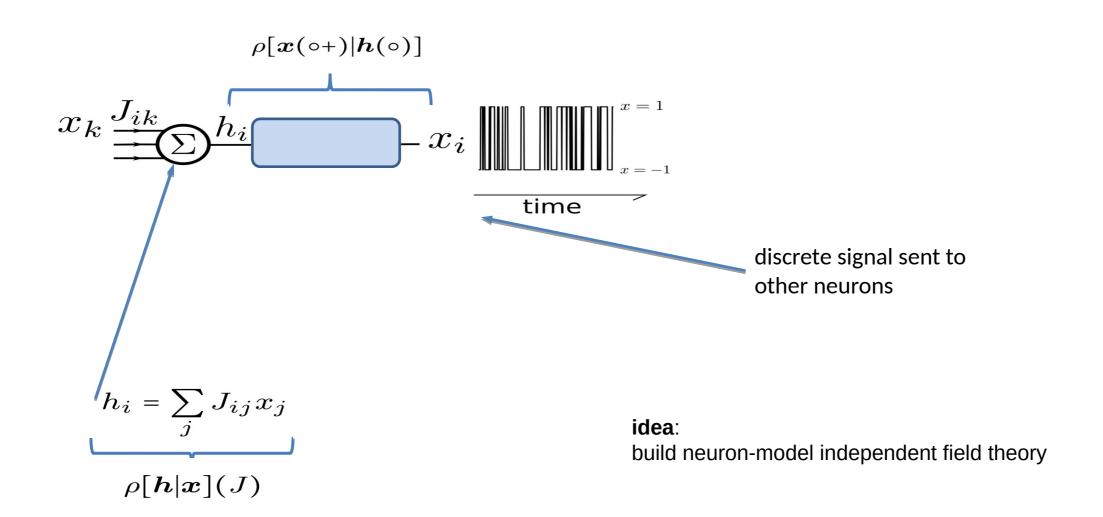


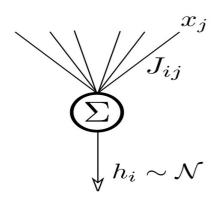
DISCRETE COUPLING: BINARY NEURON



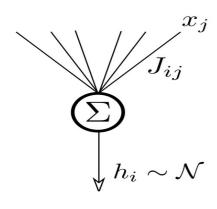
DISCRETE COUPLING: BINARY NEURON







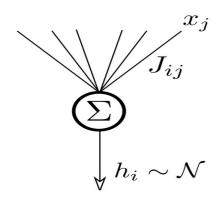
$$\rho[\boldsymbol{h}|\boldsymbol{x}](J) = \delta[\boldsymbol{h} - \boldsymbol{J}\boldsymbol{x}]$$



$$\rho[\boldsymbol{h}|\boldsymbol{x}](J) = \delta[\boldsymbol{h} - \boldsymbol{J}\boldsymbol{x}]$$

$$= \int \mathcal{D}\hat{\boldsymbol{h}} \exp(\hat{\boldsymbol{h}}^{T}\boldsymbol{h}) \exp(-\hat{\boldsymbol{h}}^{T}\boldsymbol{J}\boldsymbol{x}).$$

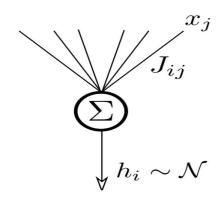
linear J in exponent

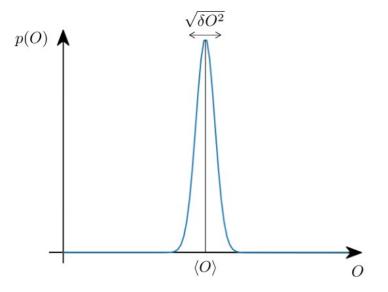


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linear J in exponent



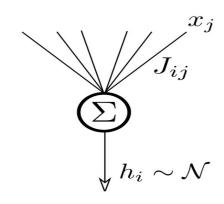


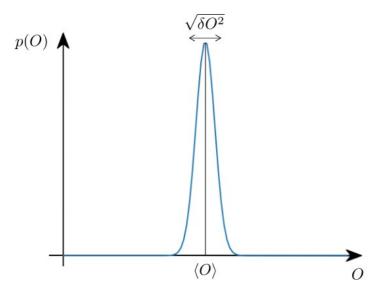
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MODEL-INDEPENDENT FIELD THEORY

linear J in exponent





instantaneous synaptic coupling

$$\rho[\boldsymbol{h}|\boldsymbol{x}](J) = \delta[\boldsymbol{h} - \boldsymbol{J}\boldsymbol{x}]$$

$$= \int \mathcal{D}\hat{\boldsymbol{h}} \exp(\hat{\boldsymbol{h}}^{\mathrm{T}}\boldsymbol{h}) \exp(-\hat{\boldsymbol{h}}^{\mathrm{T}}\boldsymbol{J}\boldsymbol{x}).$$

only term affected: interaction

$$= \langle \exp(-\hat{h}^{T} J x) \rangle_{J_{ij} \stackrel{\text{i.i.d.}}{\sim} \mathcal{N}(\frac{\bar{g}}{N}, \frac{g^{2}}{N})}$$

$$= \exp\left(-\frac{\bar{g}}{N} \hat{h}^{T} \mathcal{R} + \frac{g^{2}}{2N} \hat{h}^{T} \mathcal{Q} \hat{h}\right)$$

auxiliary fields

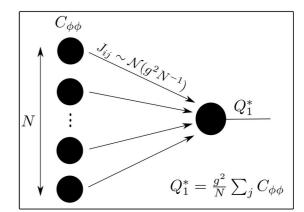
$$\mathcal{R}(t) = \frac{\bar{g}}{N} \sum_{j=1}^{N} x_j(t)$$

$$\mathcal{Q}(t,s) = \frac{g^2}{N} \sum_{j=1}^{N} x_j(t) x_j(s)$$

Macroscopic field theory

• auxiliary fields and conjugate fields $(\mathcal{R}, \mathcal{Q}, \hat{\mathcal{R}}, \hat{\mathcal{Q}}) \sim e^{NS[\mathcal{R}, \mathcal{Q}, \hat{\mathcal{R}}, \hat{\mathcal{Q}}]}$

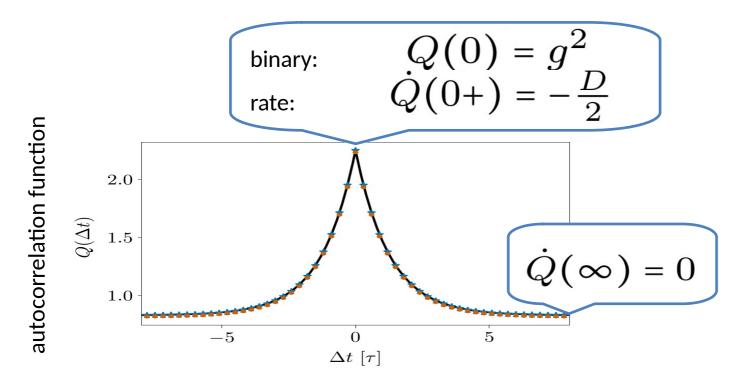
• saddle point approximation $\delta S/\delta R \stackrel{!}{=} 0, \ldots \rightarrow$



 $R(t) = \bar{g} \langle x(t) \rangle_{S(R,Q)}$ $Q(t,s) = g^2 \langle x(t)x(s) \rangle_{S(R,Q)}$

mean input to a neuron variance of input

Continuous and discrete coupling: same DMFT



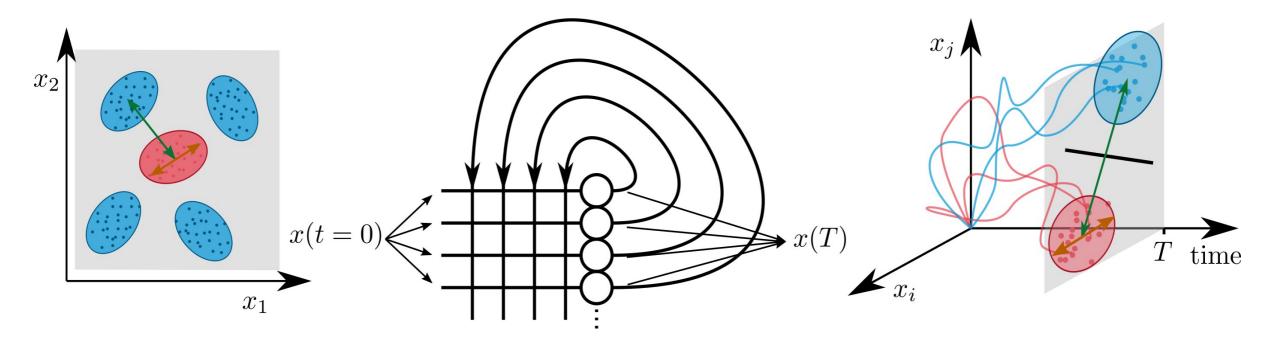
same dynamical e.o.m.

$$\tau^{2}\ddot{Q}\left(\Delta t\right) = -V'_{Q(0)}\left(Q\left(\Delta t\right)\right).$$

same activity statistics (mean and fluctuations)

CLASSIFICATION OF INPUT PATTERNS

Reservoir computing setup

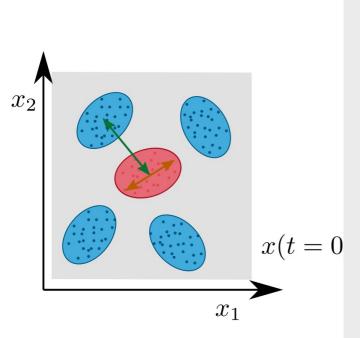


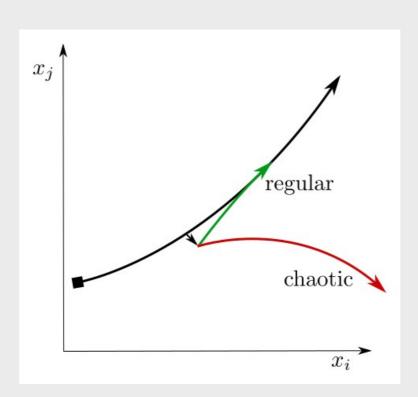


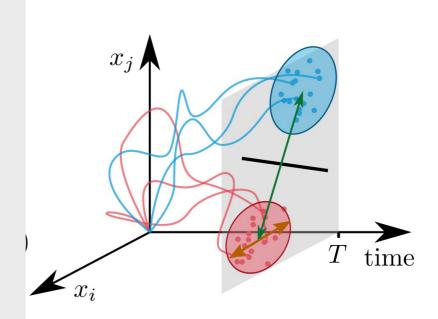


CLASSIFICATION OF INPUT PATTERNS

Reservoir computing setup



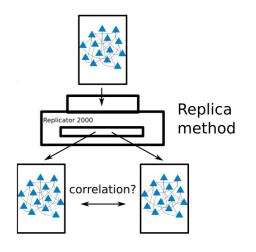




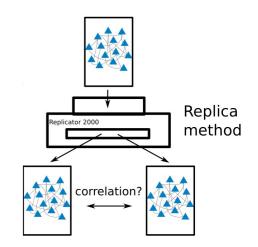




CHAOS AS CORRELATION TRANSMISSION

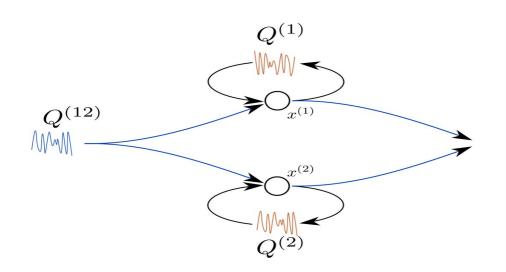


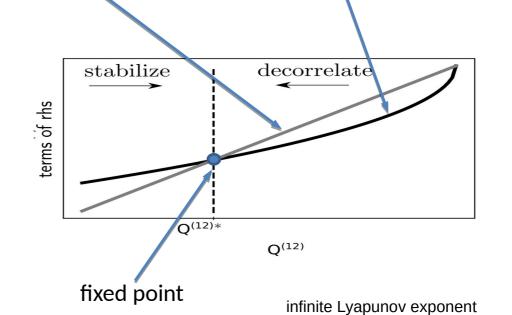
CHAOS AS CORRELATION TRANSMISSION



• correlation between replicas $Q^{(12)} = \frac{g^2}{N} \langle x^{(1)T} x^{(2)} \rangle$

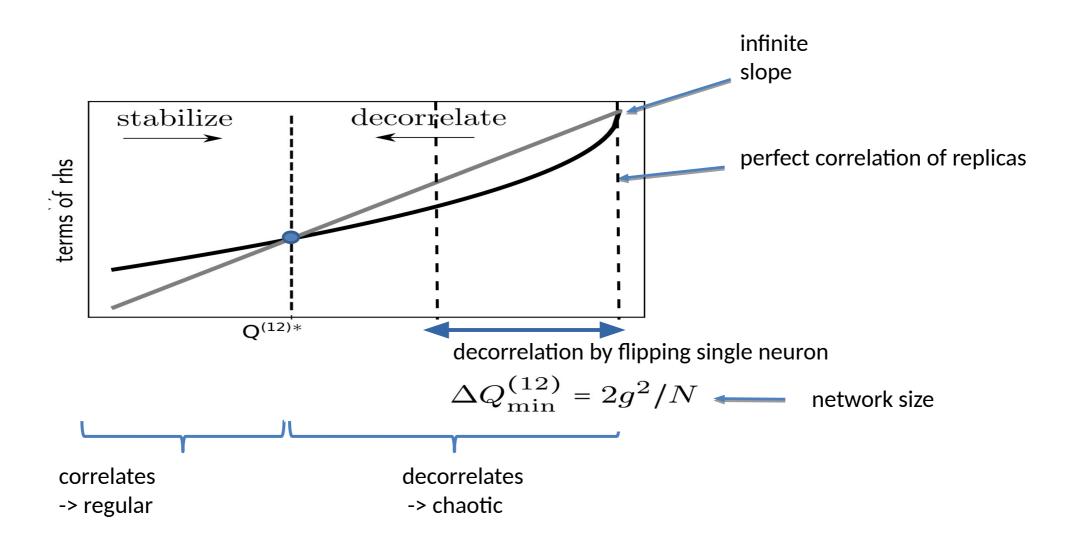
$$\tau \frac{d}{dt} Q^{(12)} (t) = - \underbrace{Q^{(12)} (t)}_{\text{correlation between outputs of replicas}} + \underbrace{g^2 \left(1 - \left\langle \left| \phi \left(h^{(1)} \right) - \phi \left(h^{(2)} \right) \right| \right) \right)}_{\text{correlation between inputs}}$$





(van Vreeswijk & Sompolinsky 1996, 1998)

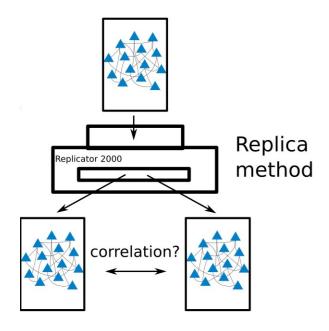
NETWORK-SIZE DEPENDENT TRANSITION



N → infinity: always chaotic Van Vreeswijk & Sompolinsky 1996

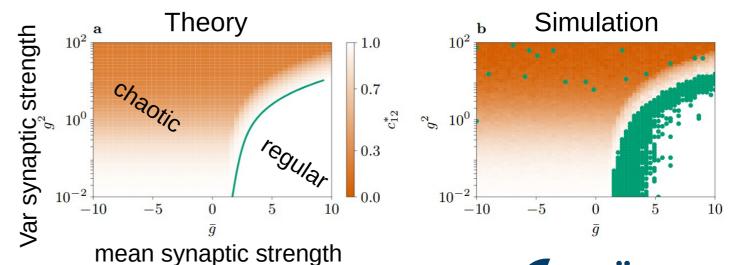
TRANSITION TO CHAOS IN BINARY NETWORKS

Replica decorrelation



condition for finite-size transition to chaos

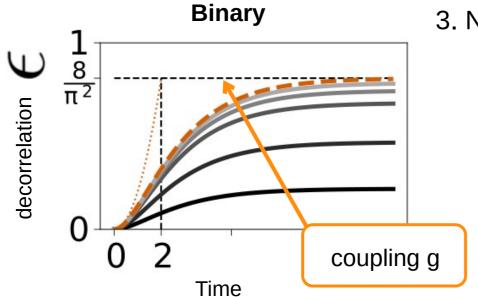




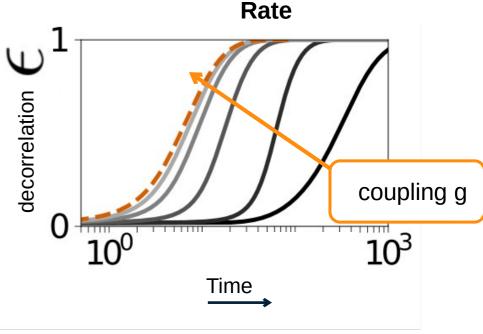
CHAOS IN BINARY NETWORKS

Differences to continuous rate networks

- 1. Mutually exclusive regimes.
- 2. Limited chaotic attractor.



3. No critical slowing down.

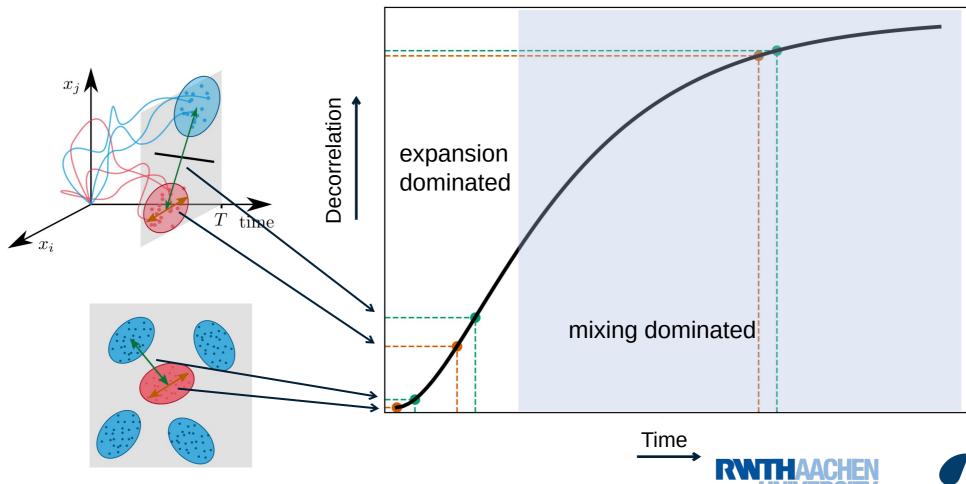






DECORRELATION CURVE

Inter-class distance increases compared to intra-class distance



TRANSIENT CHAOTIC DIMENSIONALITY EXPANSION

Classification in chaotic binary networks

• Input data: 50 Gaussian classes in 8 dim. (not linearly separable) optimal signal after $2 \ln(2) \sim 1.5$ • Linear readout accuracy peaks during expansion phase activations per neuron



 $2 \ln(2)$

coupling g



10

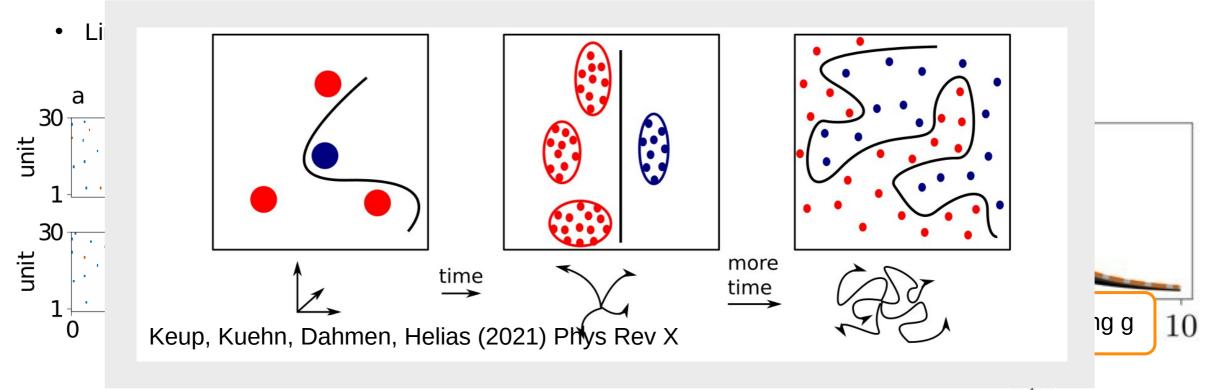
time $[\tau]$

TRANSIENT CHAOTIC DIMENSIONALITY EXPANSION

Classification in chaotic binary networks

• Input data: 50 Gaussian classes in 8 dim. (not linearly separable)

optimal signal after 2 ln(2) ~ 1.5 activations per neuron







Acknowledgments







Federal Ministry of Education and Research







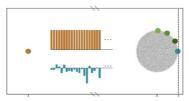


SUMMARY

novel type of critical state
 implied by wide distribution of correlations
 dynamics close to linear instability and chaos
 caused by disorder of connectivity

chaotic dynamics enhances separability
 discrete coupling: stereotypical and fast
 quick separation of signals by recurrent networks

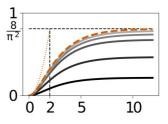




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