

Studying Gravitational Waves with the PENCIL CODE

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PENCIL CODE as a GW solver

PENCIL CODE special modules for GWs:

- gravitational_waves.f90;
- gravitational_waves_hij6.f90;
- ✓ gravitational_waves_hTXk.f90.

Coupled to HD/MHD source module:

• hydro.f90

• ...

- \cdot hydro_kinematic.f90
- magnetic.f90
- magnetic/maxwell.f90

or another special modules (?)

Studying GWs using the PENCIL CODE

- The timestep constraint in solving the gravitational wave equations sourced by hydromagnetic turbulence [Roper Pol, Brandenburg, et al. 2020]
- Numerical simulations of gravitational waves from early-universe turbulence [Roper Pol, Mandal, et al. 2020]
- *Circular polarization of gravitational waves from early-Universe helical turbulence* [Kahniashvili et al. 2021]
- *Relic gravitational waves from the chiral magnetic effect* [Brandenburg, He, et al. 2021]
- Can we observe the QCD phase transition-generated gravitational waves through pulsar timing arrays? [Brandenburg, Clarke, et al. 2021]
- Spectrum of turbulence-sourced gravitational waves as a constraint on graviton mass [He, Brandenburg, and Sinha 2021]
- The scalar, vector, and tensor modes in gravitational wave turbulence simulations [Brandenburg, Gogoberidze, et al. 2021]

GW equation in terms of comoving x, conformal t and scaled source T_{ij} in real space is¹

$$(\partial_t^2 - \boldsymbol{\nabla}^2) h_{ij}(\mathbf{x}, t) = \mathcal{G}(t) T_{ij}(\mathbf{x}, t), \tag{1}$$

where $\mathcal{G}(t) = 6/t$, and in **k** space is

$$\ddot{\tilde{h}}_{ij}(\boldsymbol{k},t) + \boldsymbol{k}^{2}\tilde{\tilde{h}}_{ij}(\boldsymbol{k},t) = \mathcal{G}(t)\tilde{T}_{ij}(\boldsymbol{k},t).$$
(2)

Can evolve the wave equation as

$$\begin{pmatrix} \boldsymbol{k}\tilde{\boldsymbol{h}} - \boldsymbol{k}^{-1}\mathcal{G}\tilde{\boldsymbol{T}} \\ \dot{\tilde{\boldsymbol{h}}} \end{pmatrix}_{+,\times}^{t+\delta t} = \begin{pmatrix} \cos \boldsymbol{k}\delta t & \sin \boldsymbol{k}\delta t \\ -\sin \boldsymbol{k}\delta t & \cos \boldsymbol{k}\delta t \end{pmatrix} \begin{pmatrix} \boldsymbol{k}\tilde{\boldsymbol{h}} - \boldsymbol{k}^{-1}\mathcal{G}\tilde{\boldsymbol{T}} \\ \dot{\tilde{\boldsymbol{h}}} \end{pmatrix}_{+,\times}^{t}.$$
(3)

¹Roper Pol, Brandenburg, et al. 2020.

• Why?

 \rightarrow Phase transitions during the early radiation era \Rightarrow inevitable turbulence.

- How?
 - \rightarrow Construct stress-energy tensor \tilde{T}_{ij} for GWs

$$\tilde{T}_{ij} = \frac{4}{3}\gamma^2 \rho u_i u_j - B_i B_j + \cdots$$
(4)

from turbulent **u**, **B** fields.

The continuity, velocity and magnetic induction equations:

$$\frac{\partial \ln \rho}{\partial t} = -\frac{4}{3} (\nabla \cdot \boldsymbol{u} + \boldsymbol{u} \cdot \nabla \ln \rho) + \frac{1}{\rho} [\boldsymbol{u} \cdot (\boldsymbol{J} \cdot \boldsymbol{B}) + \eta \boldsymbol{J}^{2}]$$

$$\frac{\partial u}{\partial t} = \frac{\boldsymbol{u}}{3} (\nabla \cdot \boldsymbol{u} + \boldsymbol{u} \cdot \nabla \ln \rho) - \frac{\boldsymbol{u}}{\rho} [\boldsymbol{u} \cdot (\boldsymbol{J} \times \boldsymbol{B}) + \eta \boldsymbol{J}^{2}] - \frac{1}{4} \nabla \ln \rho + \frac{3}{4\rho} \boldsymbol{J} \times \boldsymbol{B} + \frac{2}{\rho} \nabla \cdot (\rho \nu \mathbf{S})$$

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\boldsymbol{u} \times \boldsymbol{B} - \eta \boldsymbol{J} + \boldsymbol{\mathcal{F}}),$$
(5)

where $J = \nabla \times B$ and $D/Dt = \partial/\partial t + u \cdot \nabla$. The rate-of-strain tensor elements are $S_{ij} = \frac{1}{2}(u_{i,j} + u_{j,i}) - \frac{1}{3}\delta_{ij}u_{k,k}$. Also, relativistic EOS $p = \rho/3$ has been used.

The forcing to model generation of **B**:

$$\mathcal{F}(\mathbf{x},t) = \operatorname{Re}[\mathcal{N}\tilde{f}(\mathbf{k})\exp(i\mathbf{k}\cdot\mathbf{x}+i\varphi)].$$
(8)

Each time step *t*:

• Random \boldsymbol{k} and φ .

$$\hat{f}(k) = (k imes e)/\sqrt{k^2 - (k \cdot e)^2}$$
, where $e
min k$.

• Normalisation $\mathcal{N} = f_0 / \delta t^{1/2}$, with f_0 varied by hand to achieve approximately the same $\mathcal{E}_{\mathrm{B}}^{\max} = \langle \mathbf{B}^2 \rangle / 2$ for all runs (later).

An application: massive GWs

Modified action²

$$S = \int d^4 x (\mathcal{L}_{\rm EH} + \mathcal{L}_{m_{\rm g}} + \mathcal{L}_{\rm mat}), \text{ where } \mathcal{L}_{m_{\rm g}} = \frac{1}{4} m_{\rm g}^2 \sqrt{-g} \Big(h_{\mu\nu}^2 - \frac{1}{2} h^2 \Big).$$
(9)

• Modified GW equation in *k* space

$$\ddot{\tilde{h}}_{+,\times}(\boldsymbol{k},t) + (\boldsymbol{k}^2 + m_{\rm g}^2)\tilde{h}_{+,\times}(\boldsymbol{k},t) = \mathcal{G}(t)\tilde{T}_{+,\times}(\boldsymbol{k},t).$$
(10)

• Nonlinear dispersion:

$$\omega = \sqrt{\mathbf{k}^2 + \omega_{\rm cut}^2},\tag{11}$$

where $\omega_{\rm cut} = m_{\rm g} c^2 / \hbar$ is an effective graviton mass term³.

²Fierz and Pauli 1939. ³Lee 2013.

An application: massive GWs

Spatial spectrum:

- 0.3 $\leq \omega_{\mathrm{cut}} \leq$ 10;
- $\label{eq:mg} \begin{array}{l} \cdot \ 2.2 \times 10^{-23} \mathrm{eV} \lesssim m_{\mathrm{g}} \lesssim \\ 7.4 \times 10^{-22} \mathrm{eV}; \end{array}$
- $\Omega_{\rm GW} \sim k^1 \Rightarrow k^3;$
- $h_c \sim k^{-1/2} \Rightarrow k^{3/2};$
- Consistent change for different k_f.



Figure 1: Spatial energy & strain spectra. $k_f = 2$ (up) & $k_f = 6$ (down).

An application: massive GWs

Temporal spectrum:

- · $\omega_{\mathrm{cut}} \in \{0,1\}$,
- + $m_{
 m g} \approx 7.4 \times 10^{-23} {
 m eV};$
- Clear cut-off;
- $f \sim 10 \mathrm{nHz};$
- $\Delta t_{\rm phys} \sim {
 m yr}.$



Figure 2: Temporal strain spectra. $\omega_{\mathrm{cut}}=0$ (up) and $\omega_{\mathrm{cut}}=1$ (down). 9

PENCIL CODE for GW spectral analysis:

- Turbulent source: kinetic, (electro)magnetic...
- Early universe: EWPT, QCDPT...
- Modified gravity: massive gravity and more.

Questions?