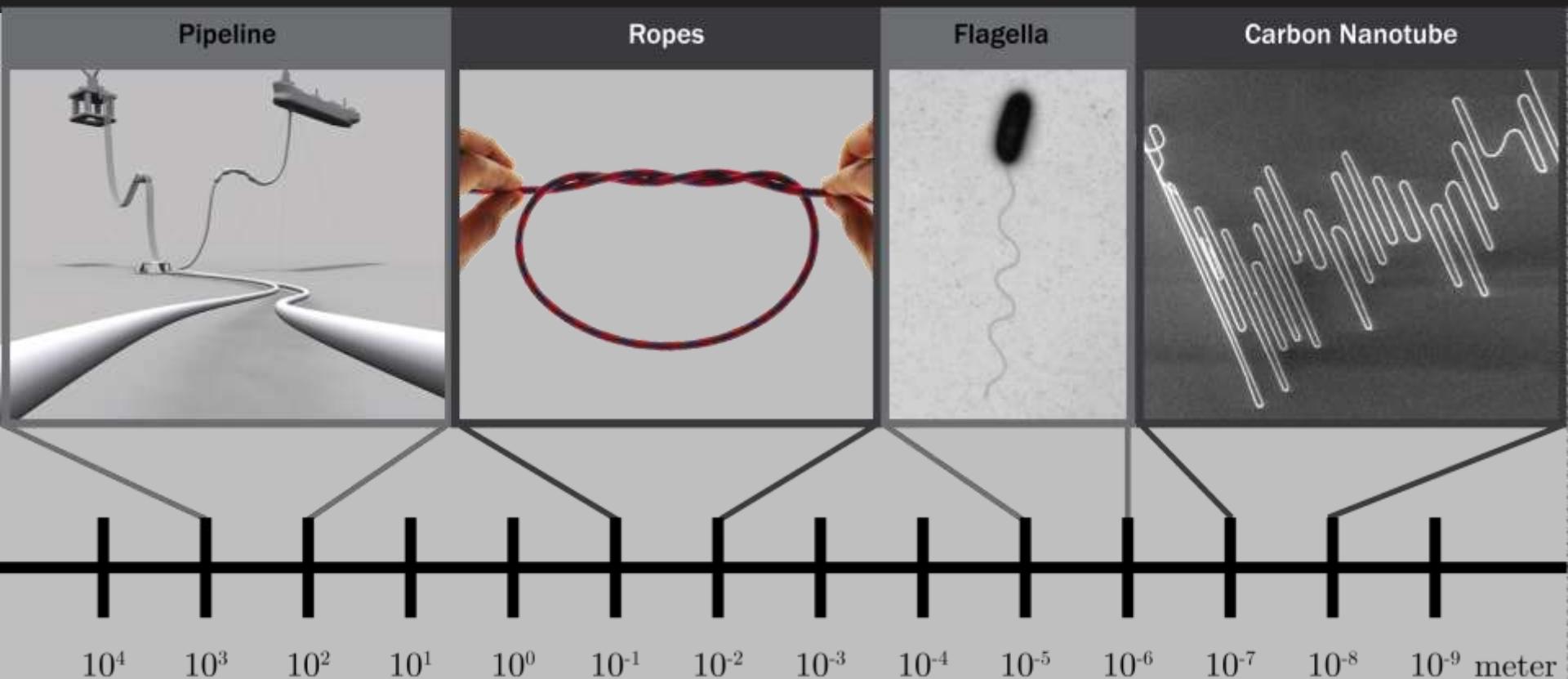


Mechanics of thin elastic rods

Engineering meets computer graphics



Khalid Jawed

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Department of Mechanical Engineering
Massachusetts Institute of Technology



Agenda

Part 1

Pipeline



10⁴ 10³ 10² 10¹

Ropes



10⁰ 10⁻¹ 10⁻² 10⁻³

Part 2

Flagella



10⁻⁴ 10⁻⁵ 10⁻⁶

Carbon Nanotube



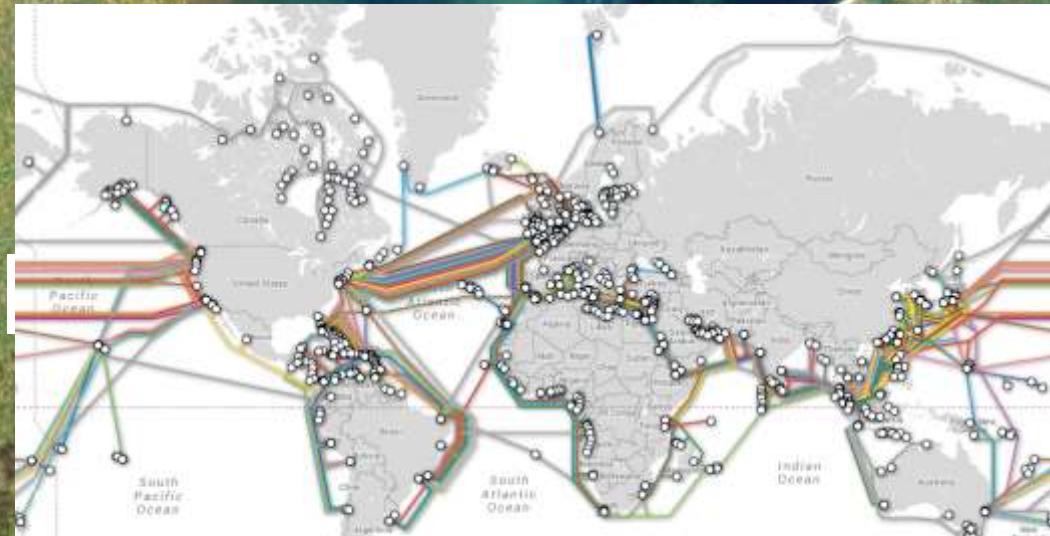
10⁻⁶ 10⁻⁷ 10⁻⁸ 10⁻⁹ meter

Coiling of rods

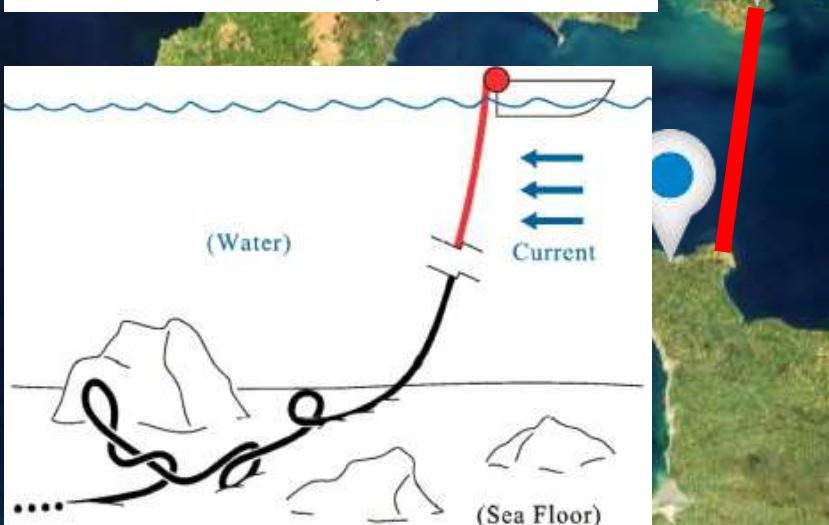
Propulsion of
flagella



Submarine Cable Deployment
Global Marine Systems



Submarine Cable Map



Tangles in cables
Goyal et al, *Int. J. Non Linear Mech.* 2008



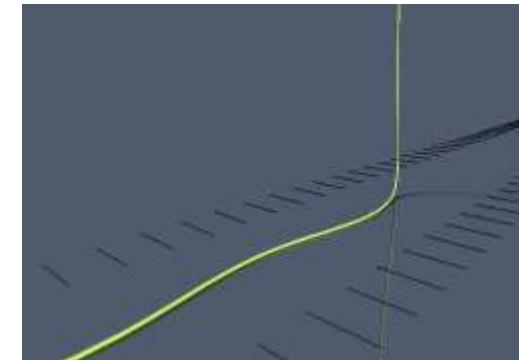
With Fang Da, J. Joo, Eitan Grinspun
Columbia Computer Graphics Group
Columbia University



PipeLine Under The Ocean (PLUTO), 1944

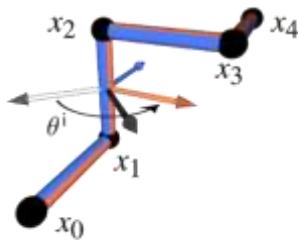
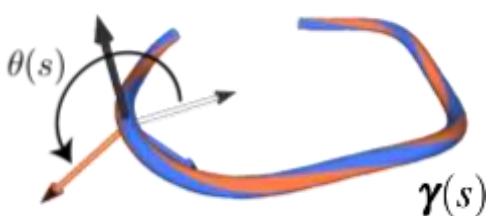
Numerics: Discrete Elastic Rods

Bergou et al, SIGGRAPH 2008



Twist-free frame

Material frame



$$\begin{aligned} E_{bending} &= \frac{1}{2} \int EI\kappa^2 ds & \kappa &= \frac{d^2\gamma}{ds^2} \\ E_{twisting} &= \frac{1}{2} \int GJ\psi^2 ds & \psi &= \frac{d\theta}{ds} \\ E_{stretching} &= \frac{1}{2} \int EA\alpha^2 ds & \alpha &= \frac{\Delta l}{l_o} \end{aligned}$$

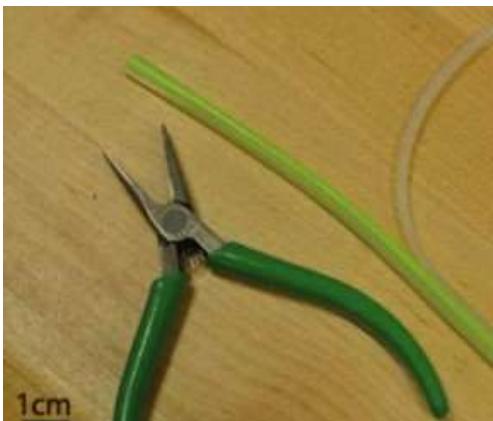
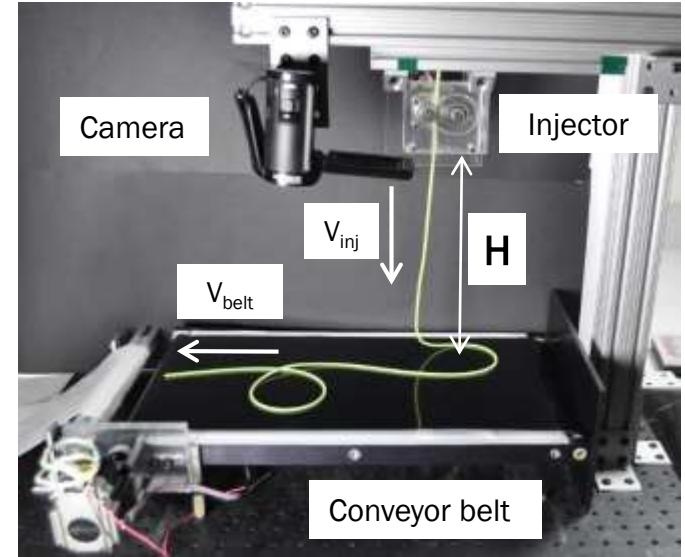
Discrete formulation
Discrete Differential Geometry

$$\mathbf{F} = -\nabla E$$

$$\mathbf{F}_{bending} + \mathbf{F}_{twisting} + \mathbf{F}_{stretching} = \mathbf{F}_{ext}$$

Experiments: rod fabrication

Miller et al, *Phys. Rev. Lett.* 2014
Jawed et al, *Proc. Natl. Acad. Sci.* 2014



Demolding



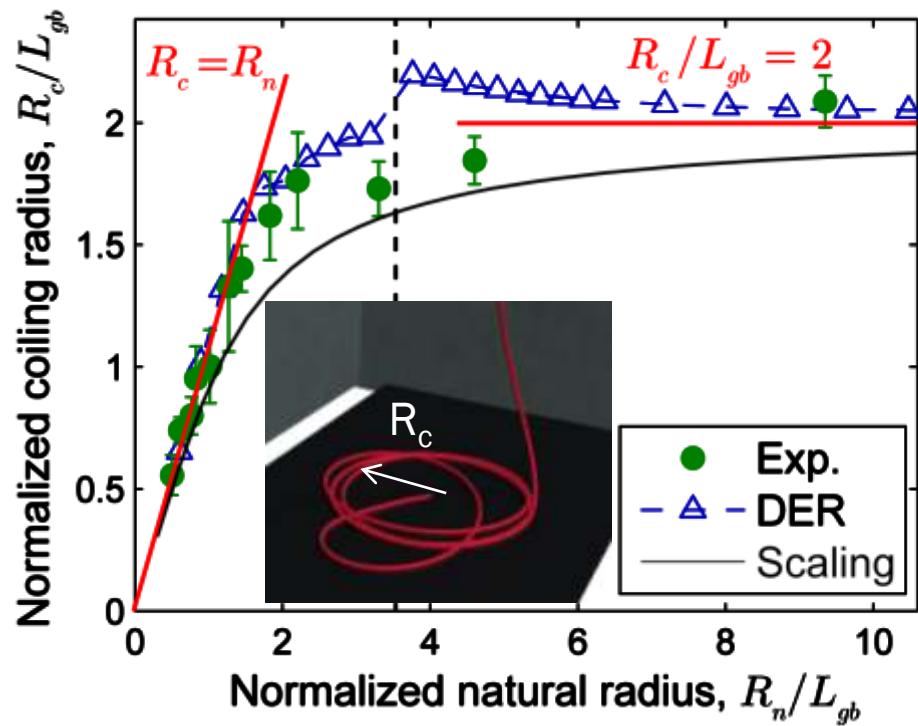
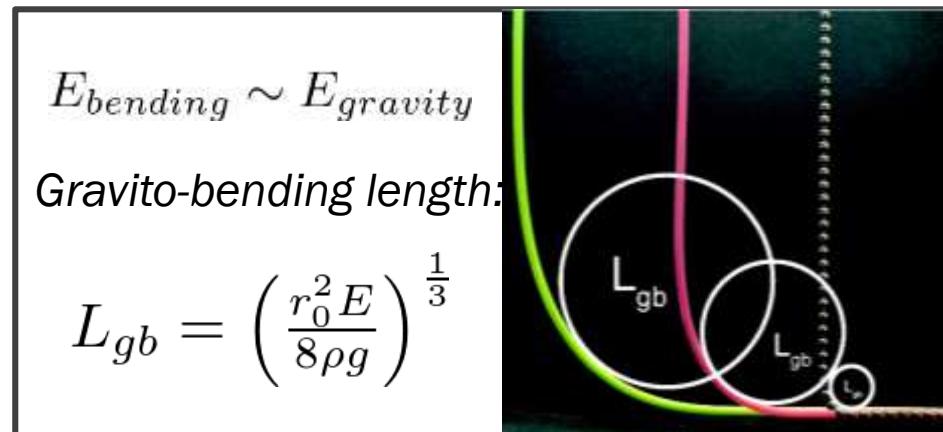
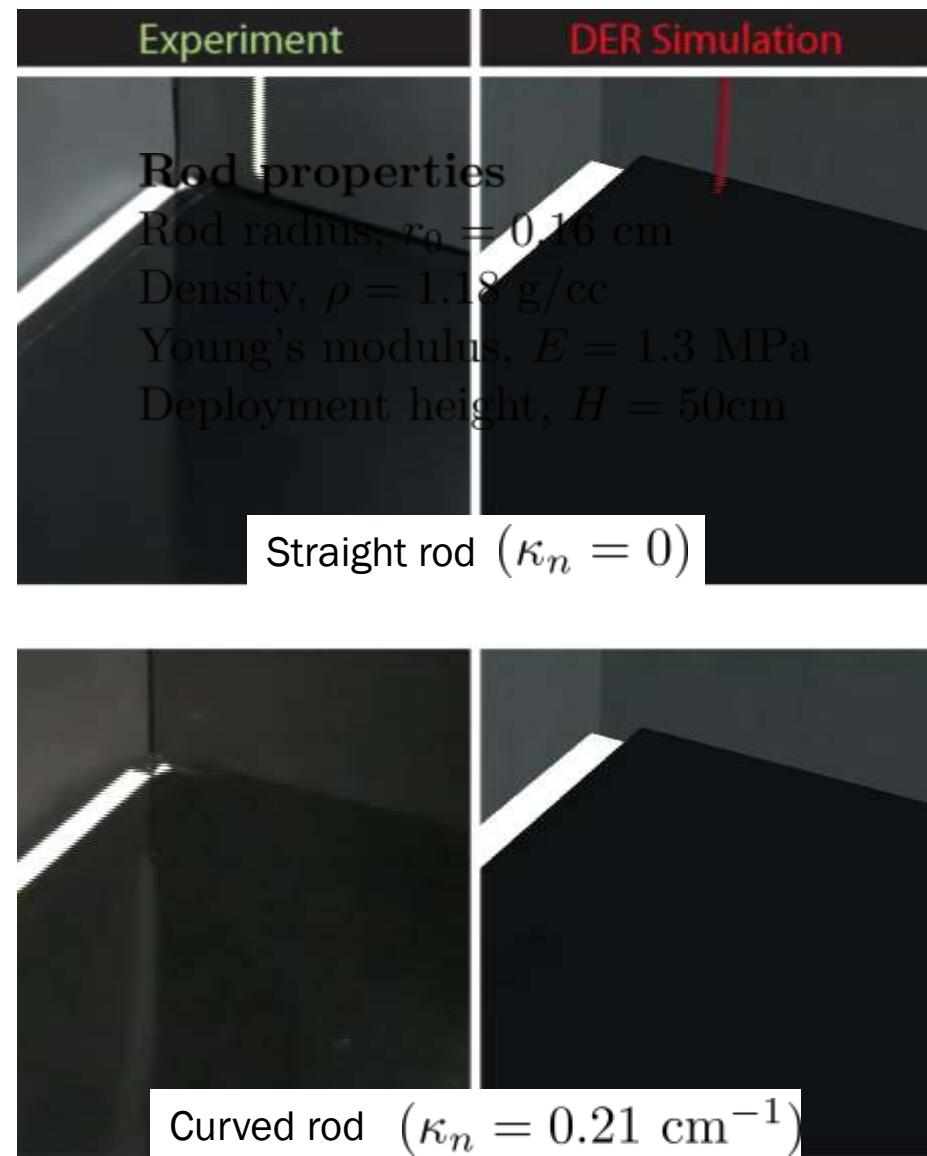
Curvature in engineering
Jindaltubes.com

$$\epsilon = \frac{v_{inj} - v_{belt}}{v_{inj}}$$

Control parameter:
Dimensionless speed mismatch

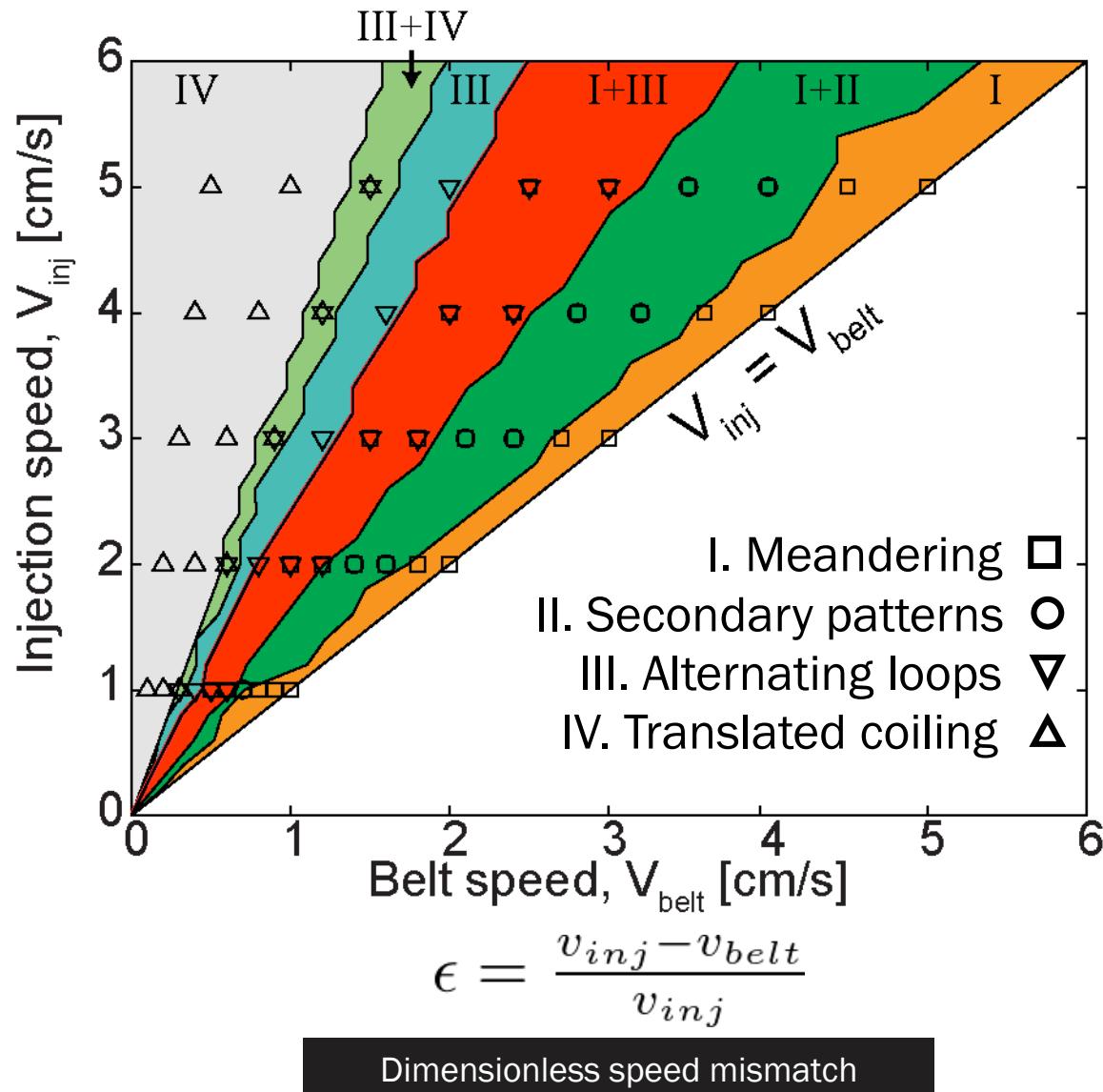
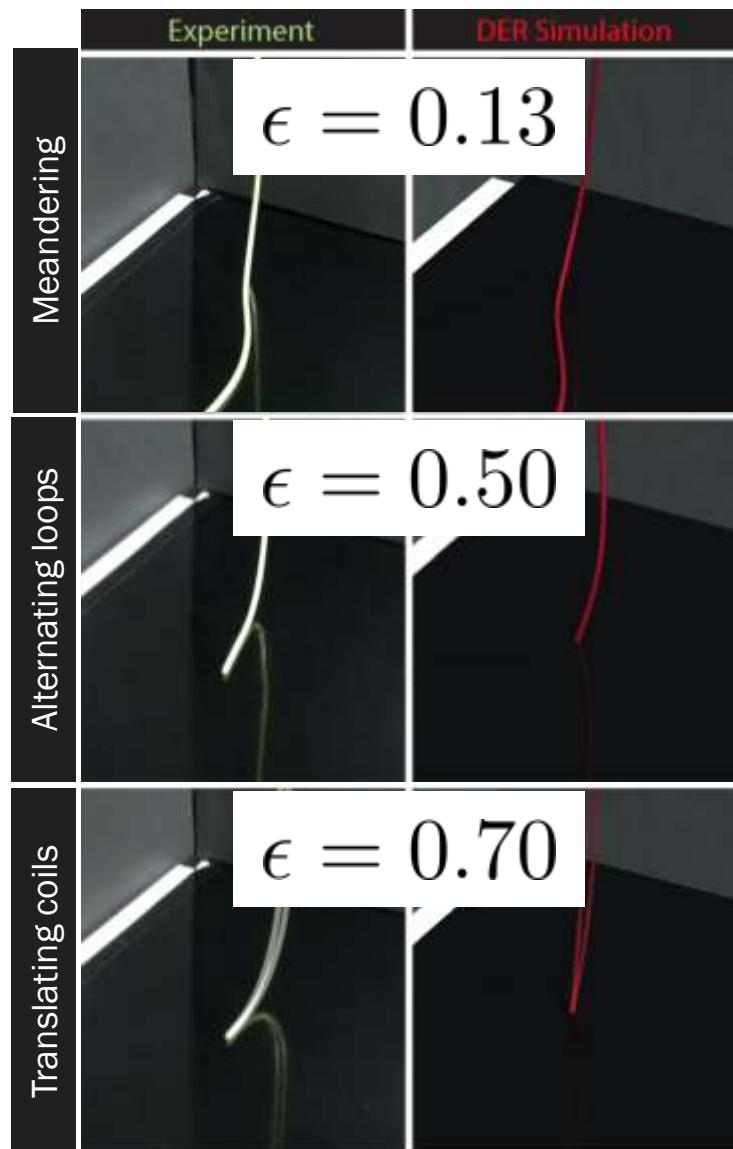
Static coiling

Jawed et al, Proc. Natl. Acad. Sci. 2014



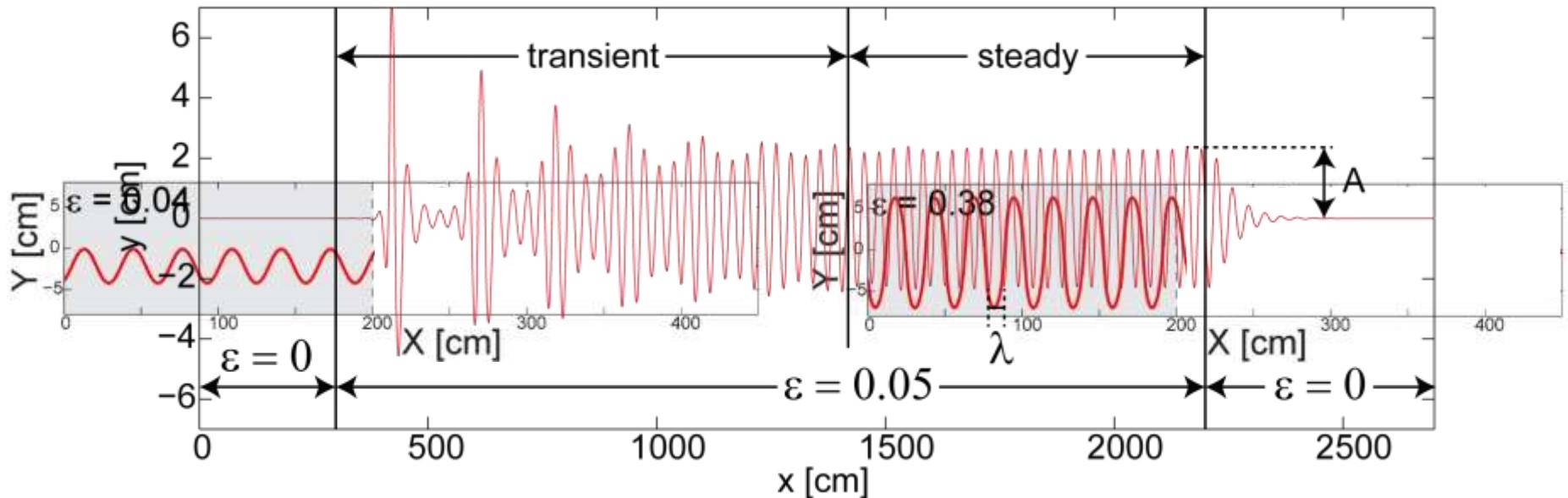
Dynamic coiling: straight rod

Jawed et al, Proc. Natl. Acad. Sci. 2014
Jawed et al, Extreme Mechanics Letters, 2014

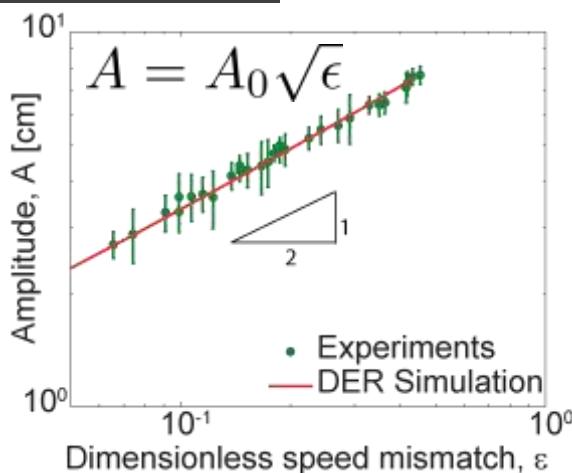


Meandering: Hopf bifurcation?

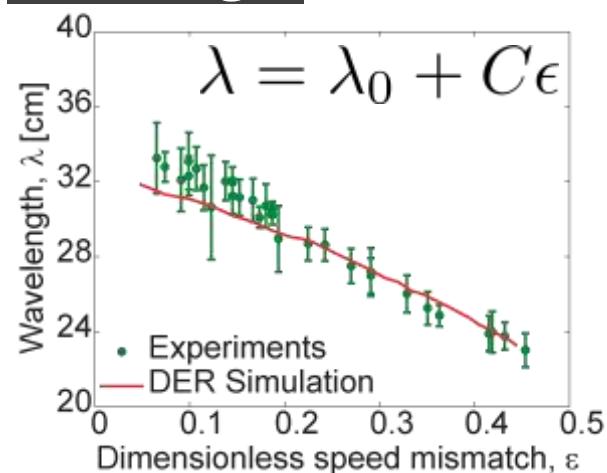
Jawed et al, Proc. Natl. Acad. Sci. 2014



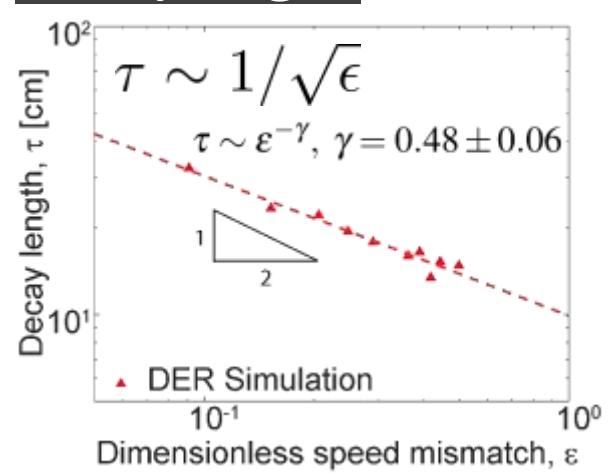
Amplitude



Wavelength



Decay length



Meandering length scale

Jawed et al, Proc. Natl. Acad. Sci. 2014

Geometry

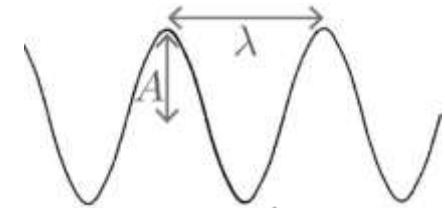
$$A = A_0 \sqrt{\epsilon}$$

$$\lambda \approx A_0(-2.48\epsilon + 3.2)$$

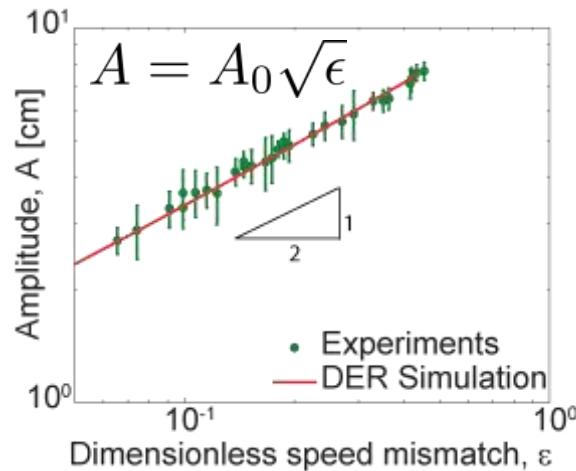
Onset wavelength

$$\lambda_0 = \lambda(\epsilon \rightarrow 0) \approx 3.2A_0$$

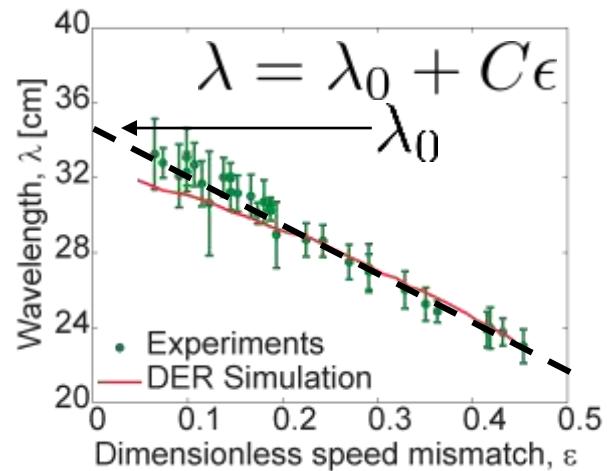
What sets λ_0 ?



Amplitude



Wavelength



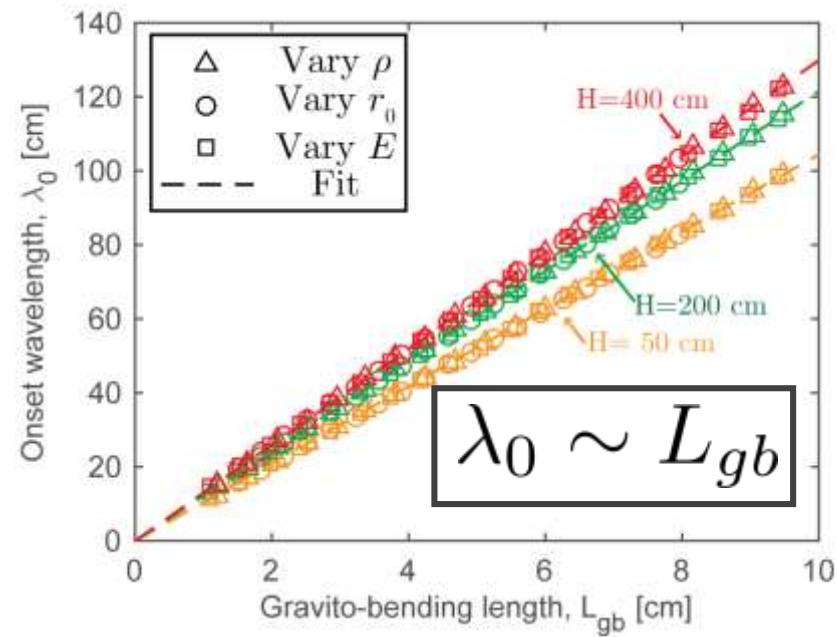
Meandering length scale

Jawed et al, Proc. Natl. Acad. Sci. 2014

Geometry

$$A = A_0 \sqrt{\epsilon}$$

$$\lambda \approx A_0(-2.48\epsilon + 3.2)$$



Onset wavelength

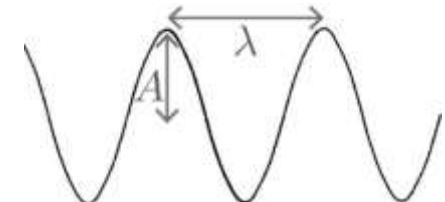
$$\lambda_0 = \lambda(\epsilon \rightarrow 0) \approx 3.2A_0$$

What sets λ_0 ?

Physical parameters

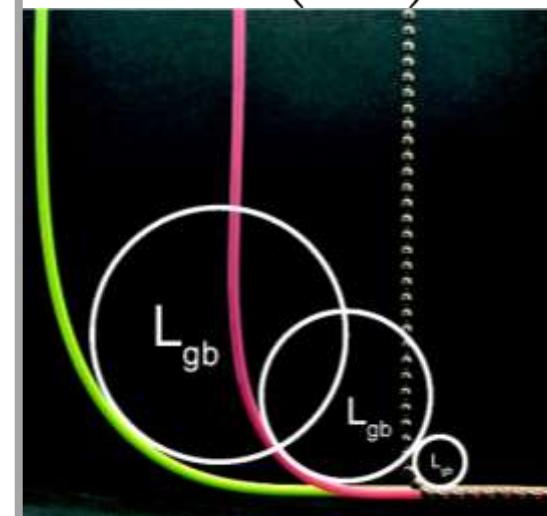
Density, ρ
Young's modulus, E
Rod radius, r_0

Numerical Experiment
Vary one parameter
keeping others fixed



Gravito-bending length scale

$$L_{gb} = \left(\frac{r_0^2 E}{8\rho g} \right)^{\frac{1}{3}}$$



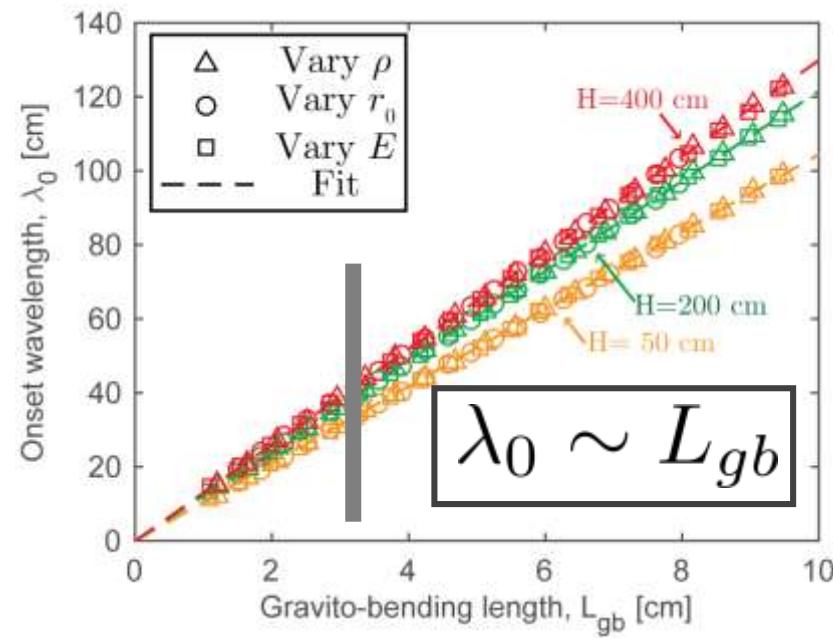
Meandering length scale

Jawed et al, Proc. Natl. Acad. Sci. 2014

Geometry

$$A = A_0 \sqrt{\epsilon}$$

$$\lambda \approx A_0(-2.48\epsilon + 3.2)$$

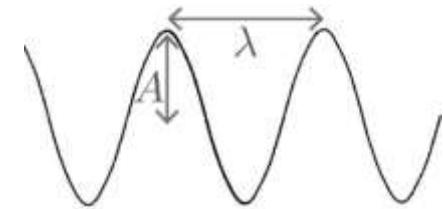
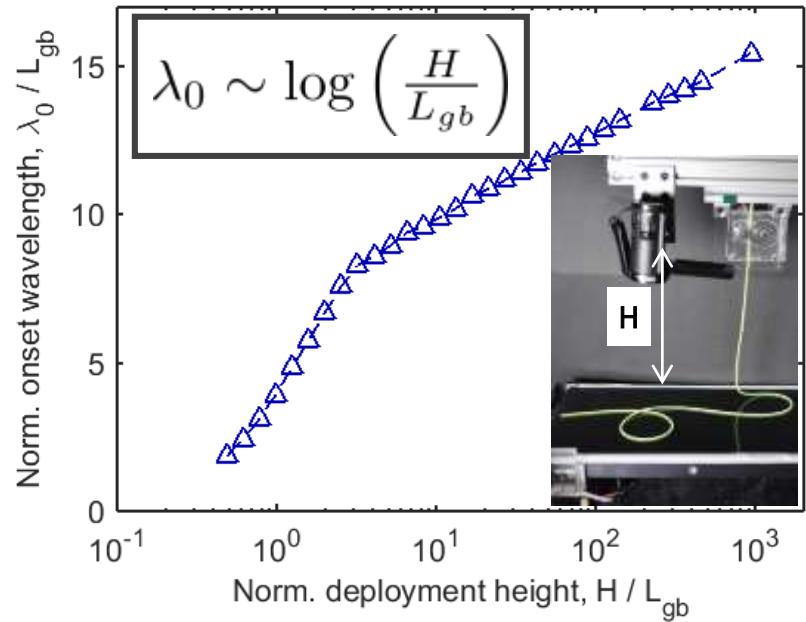


Onset wavelength

$$\lambda_0 = \lambda(\epsilon \rightarrow 0) \approx 3.2A_0$$

What sets λ_0 ?

Deployment height, H



Meandering length scale

Jawed et al, Proc. Natl. Acad. Sci. 2014

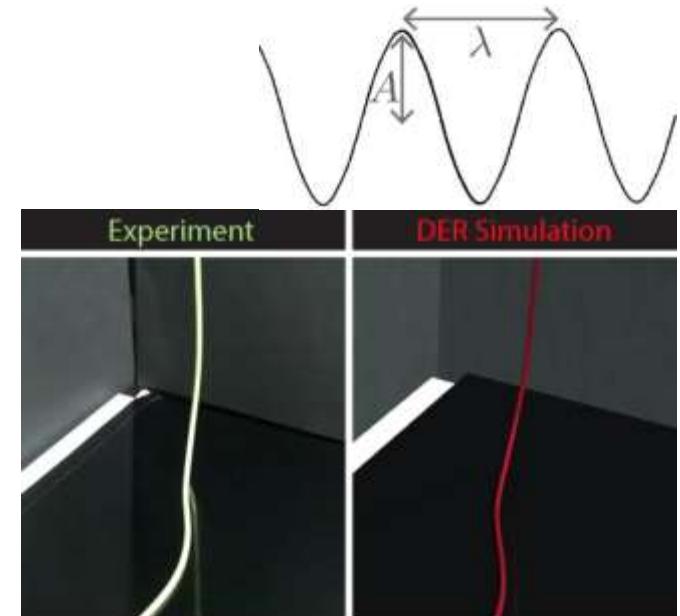
Putting it all together:

Onset wavelength

$$\lambda_0 = L_{gb} \left(D_1 \log \left(\frac{H}{L_{gb}} \right) + \beta \right)$$

$$D_1 = 1.22 \pm 0.01$$

$$\beta = 7.14 \pm 0.02$$



Meandering wavelength

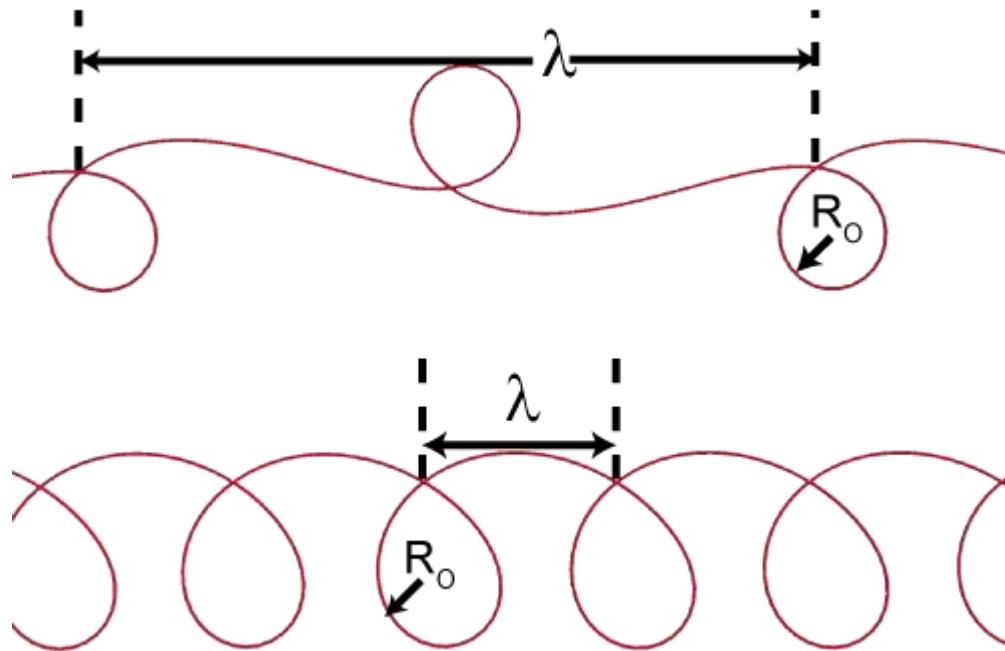
$$\lambda = \frac{-2.48\epsilon+3.2}{3.2} L_{gb} \left(D_1 \log \left(\frac{H}{L_{gb}} \right) + \beta \right)$$

Meandering amplitude

$$A = \frac{L_{gb}}{3.2} \left(D_1 \log \left(\frac{H}{L_{gb}} \right) + \beta \right) \sqrt{\epsilon}$$

Meandering Coiling length scale

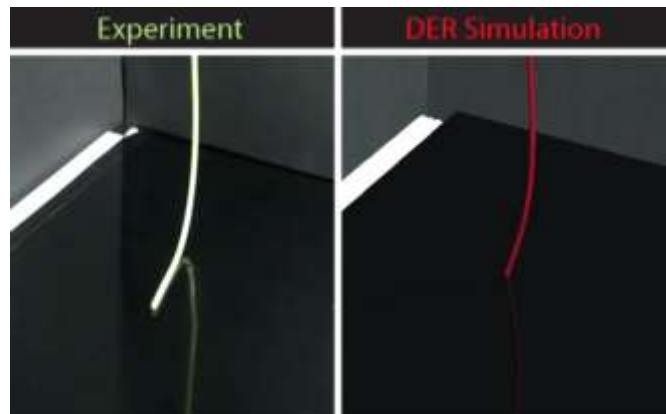
Jawed et al, *Extreme Mechanics Letters* 2014



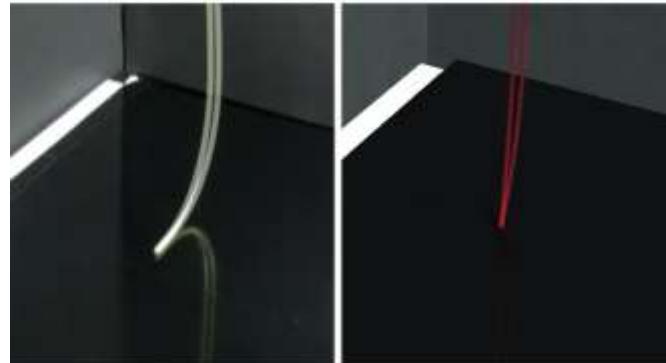
$$\lambda \sim L_{gb} \log \left(\frac{H}{L_{gb}} \right)$$

$$R_o \sim L_{gb} \log \left(\frac{H}{L_{gb}} \right)$$

Alternating loops



Translated coiling



Prominence of geometry

Fluid-mechanical sewing machine

Stephen Morris, U of Toronto

Meandering thread experiment

Viscous syrup falling onto a moving belt

Eitan Grinspun, Columbia U.
Meandering

Why similarity?

Jawed, M. K. et al
A Geometric Model
for the Coiling of
an Elastic Rod
Deployed Onto a
Moving Substrate
Journal of Applied
Mechanics, 82
(12):121007,
2015

Molten glass sewing machine

Brun et al, MIT

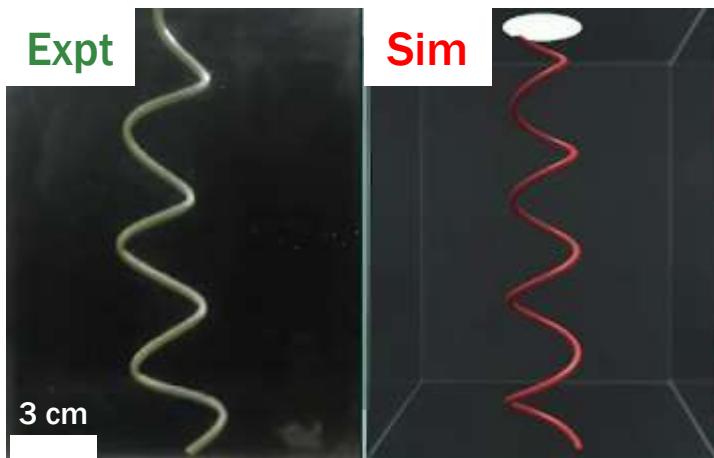
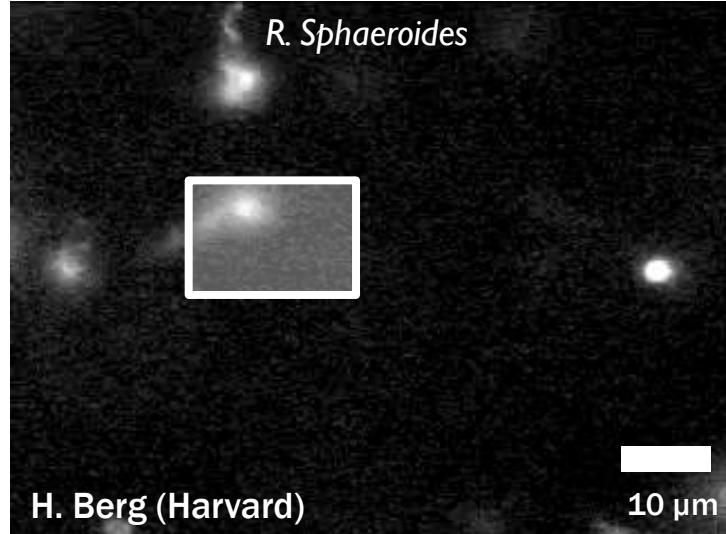
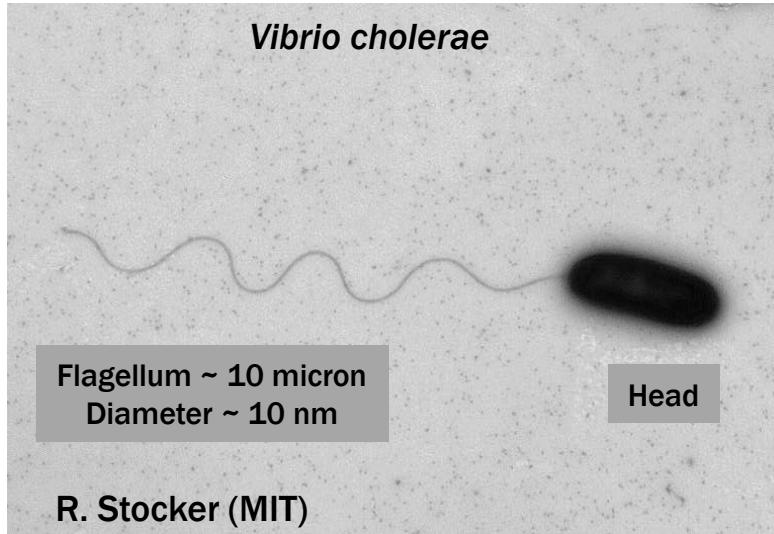
Pendulum
Coil

COILING SPEED U_c
 $U_c = \Omega_c R_c$

Cooking of Jalebi



Propulsion and instability of bacterial flagella



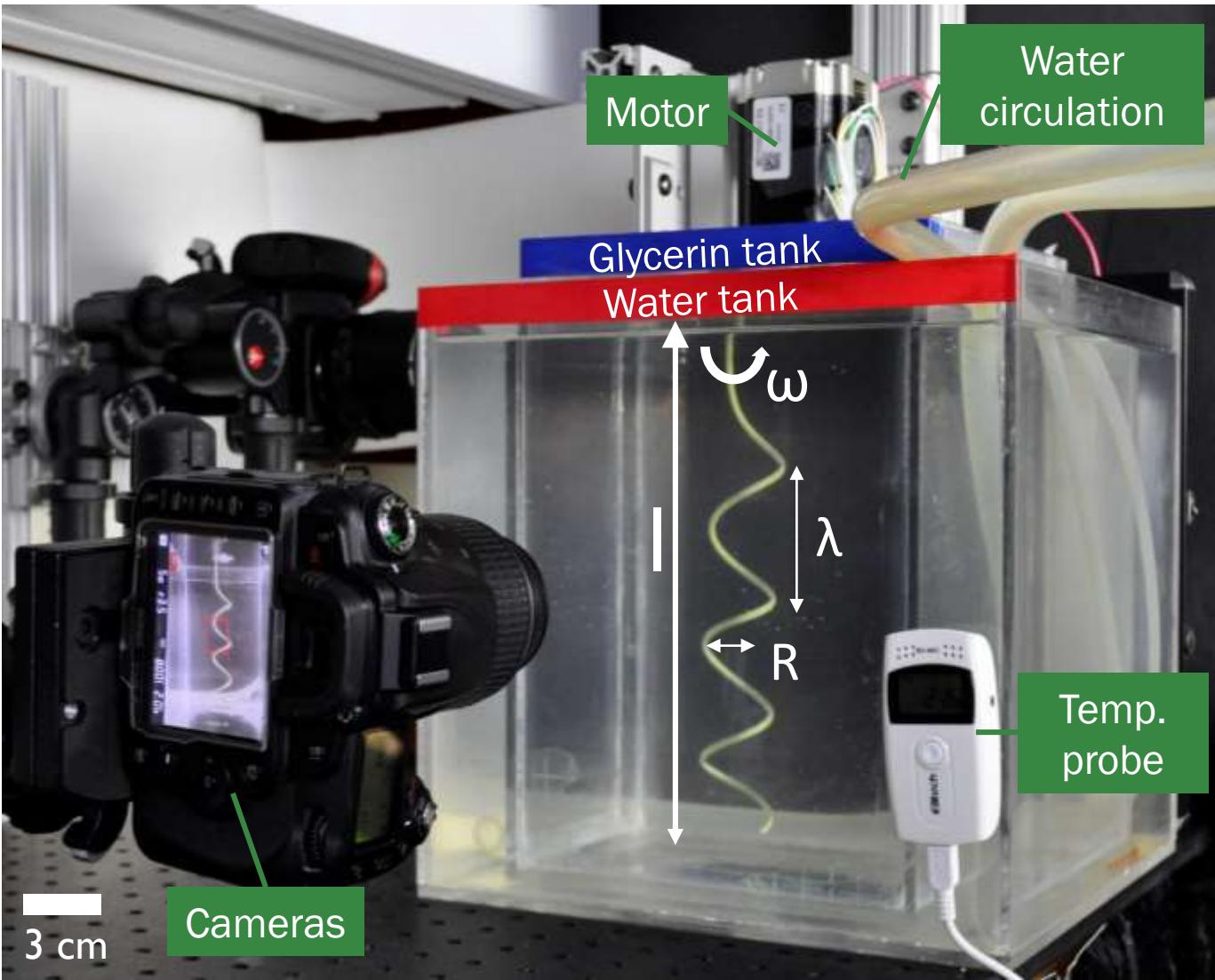
Khalid Jawed
Noor Khouri
Pedro M. Reis
Elasticity, Geometry and Statistics Lab
Massachusetts Institute of Technology



Fang Da
Eitan Grinspun
Columbia Computer Graphics Group
Columbia University



Experiments



Physical properties

Axial length
 $12 \leq l[\text{cm}] \leq 20$

Helix pitch
 $3 \leq \lambda[\text{cm}] \leq 12$

Helix radius
 $0.95 \leq R[\text{cm}] \leq 2.22$

Rod radius
 $0.94 \leq r_0[\text{mm}] \leq 2.38$

Rod density
 $1.13 \leq \rho_r[\text{g/cm}^3] \leq 1.37$

Young's Modulus
 $636 \leq E[\text{kPa}] \leq 1255$

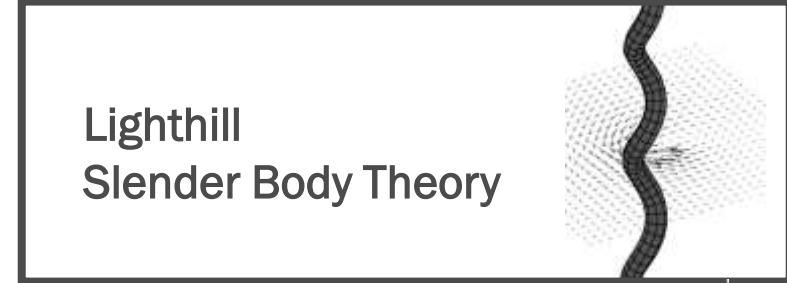
Temperature of glycerin
 $7.55 \leq \theta[\text{°C}] \leq 20$

Viscosity of glycerin
 $0.5 \leq \mu[\text{Pa} \cdot \text{s}] \leq 4.45$

Numerics



Discrete
Elastic
Rod



Lighthill Slender Body Theory

Relation between velocity $\mathbf{u}(s)$ at each point on flagellum and the force $\mathbf{f}(s)$ exerted by the fluid.

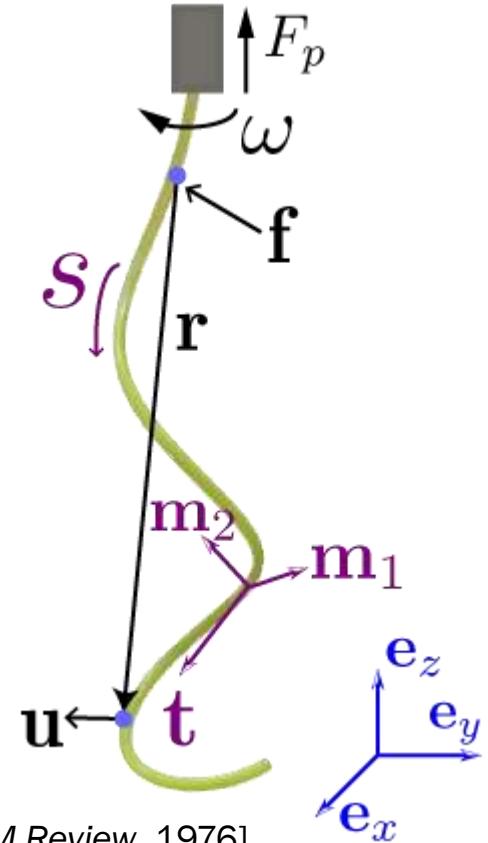
local non-local

$$\mathbf{u}(s) = \frac{\mathbf{f}_\perp}{4\pi\mu} + \int_{r(s',s)>\delta} \mathbf{f}(s') \cdot \mathbf{J}(\mathbf{r}) ds'$$

where $\mathbf{f}_\perp = \mathbf{f} \cdot (\mathbf{I} - \mathbf{t}\mathbf{t}^T)$

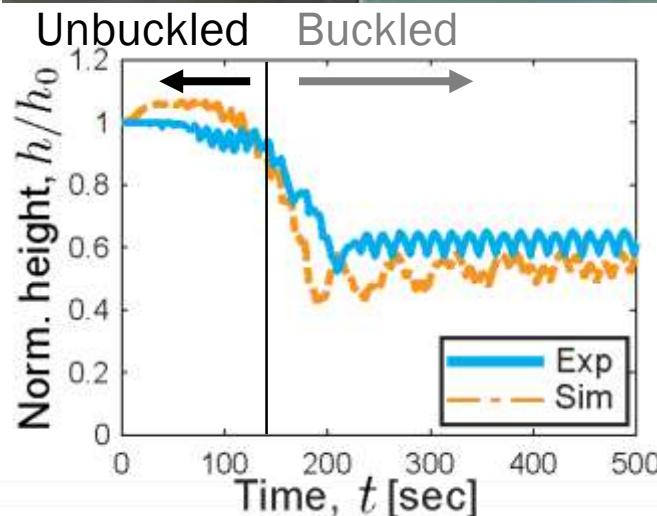
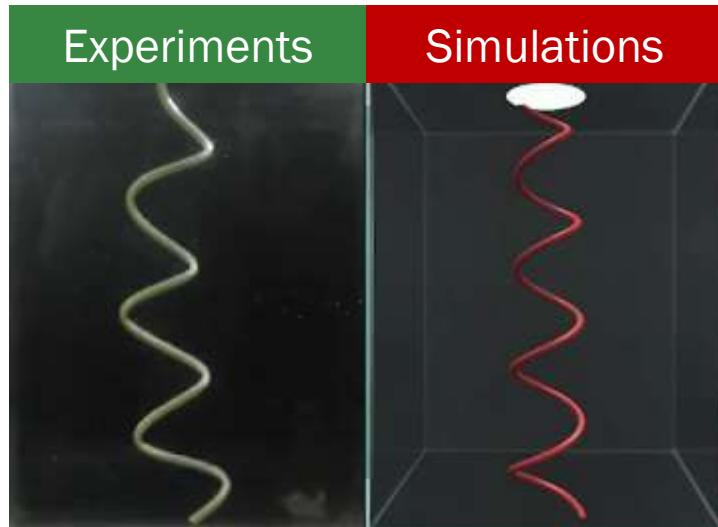
$$\delta = \frac{1}{2} r_0 \sqrt{e}$$

$$\mathbf{J}(\mathbf{r}) = \frac{1}{8\pi\mu} \left(\frac{\mathbf{I}}{|\mathbf{r}|} + \frac{\mathbf{r}\mathbf{r}^T}{|\mathbf{r}|^3} \right)$$



[Lighthill SIAM Review, 1976]

Flagellar dynamics



Physical properties

$$l = 20 \pm 0.5 \text{ cm}$$

$$Y = 1255 \pm 49 \text{ kPa}$$

$$\lambda = 5 \pm 0.5 \text{ cm}$$

$$\mu = 1.6 \pm 0.05 \text{ Pa} \cdot \text{s}$$

$$R = 1.59 \pm 0.1 \text{ cm}$$

$$\rho_m = 1.240 \text{ g/cm}^3$$

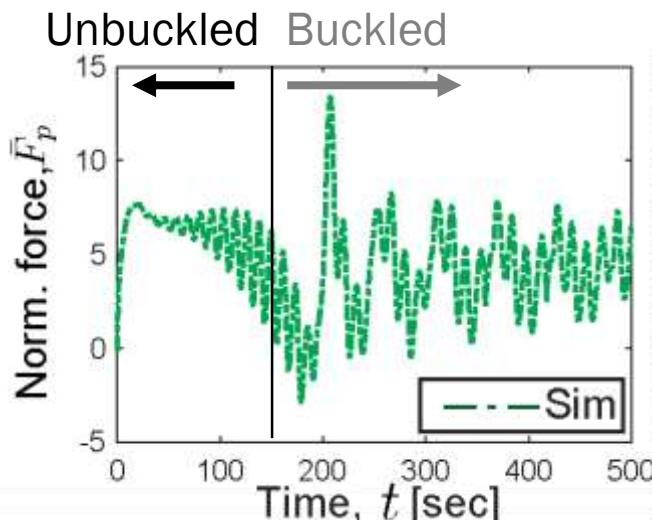
$$r_0 = 1.58 \pm 0.02 \text{ mm}$$

$$Re \sim 10^{-2}$$

$$\rho_r = 1.273 \pm 0.022 \text{ g/cm}^3$$

Bacteria $3 \leq \lambda/R \leq 11$
 $2.7 \leq l/\lambda \leq 11$

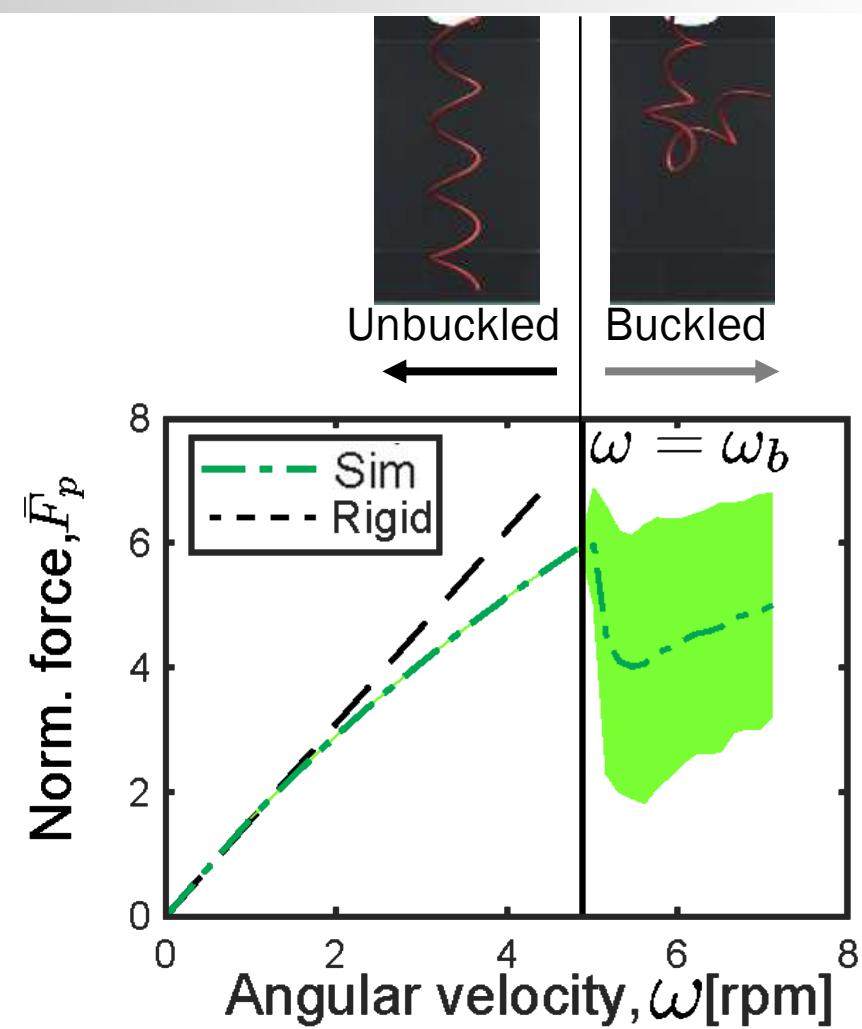
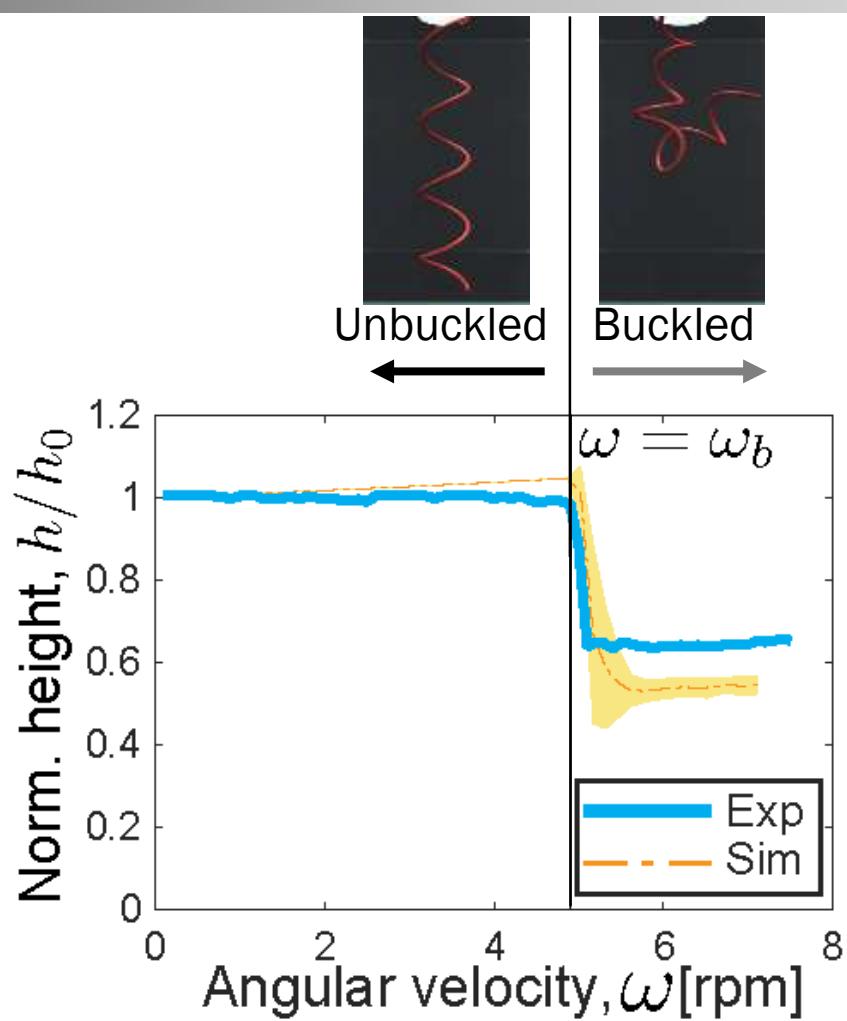
Model $\lambda/R = 3.14$
 $l/\lambda = 4.0$



Fluid loading can cause buckling

$$\bar{F}_p = \frac{F_p}{EI/l^2}$$

Critical buckling velocity



Critical angular velocity for buckling: ω_b

Buckling characterization

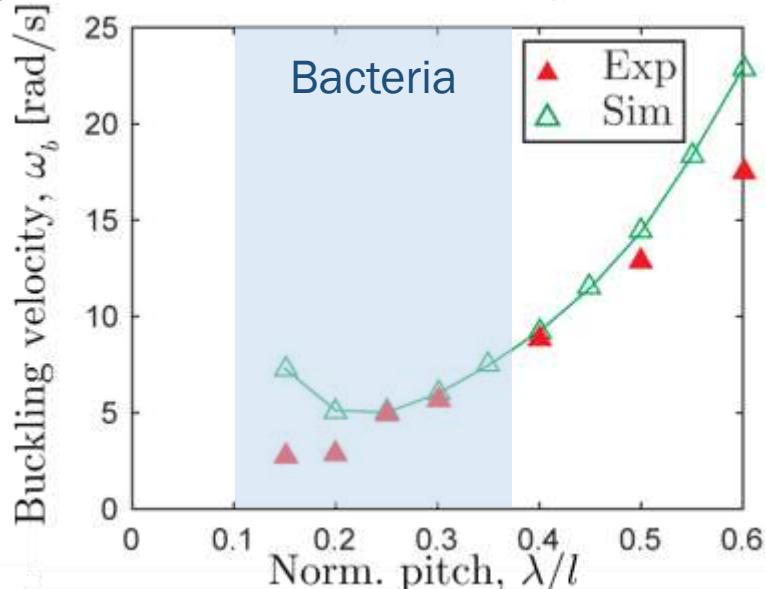
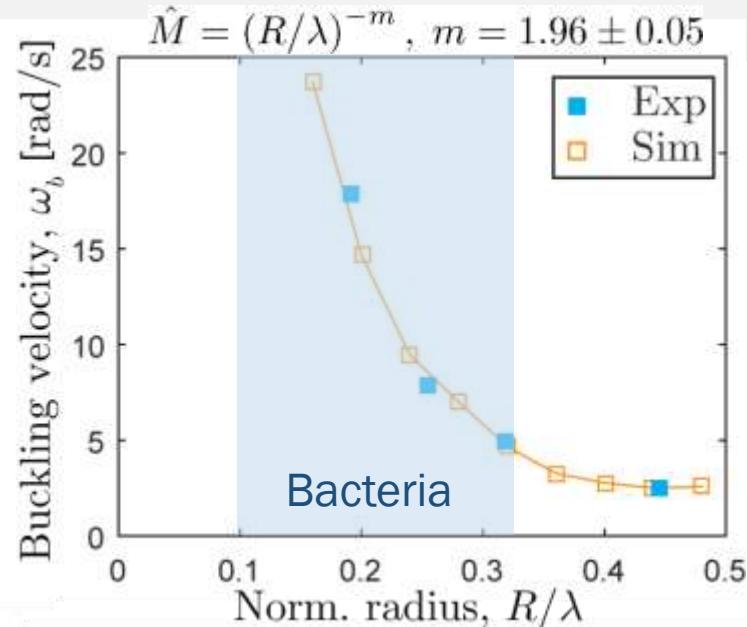
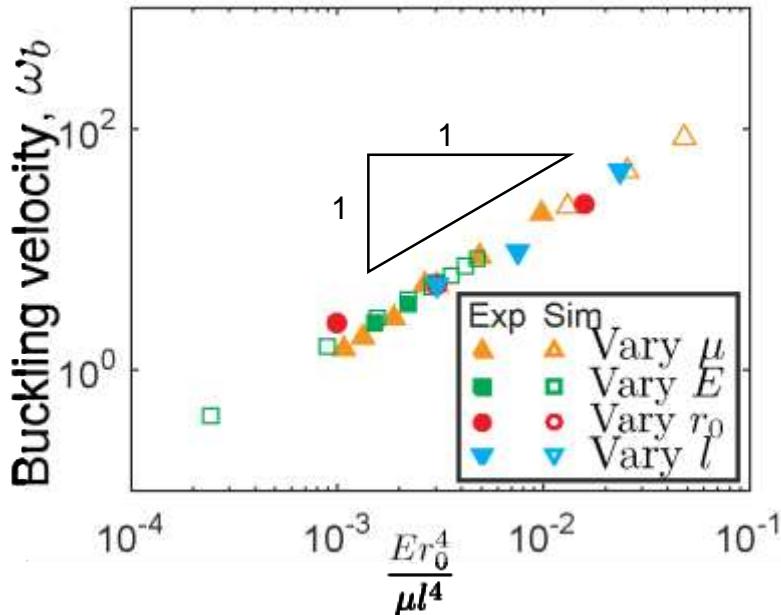
Critical force \sim Viscous force

$$F_c \sim F_v$$

$$\omega_b = \frac{E r_0^4}{\mu l^4} \bar{\omega}_b$$

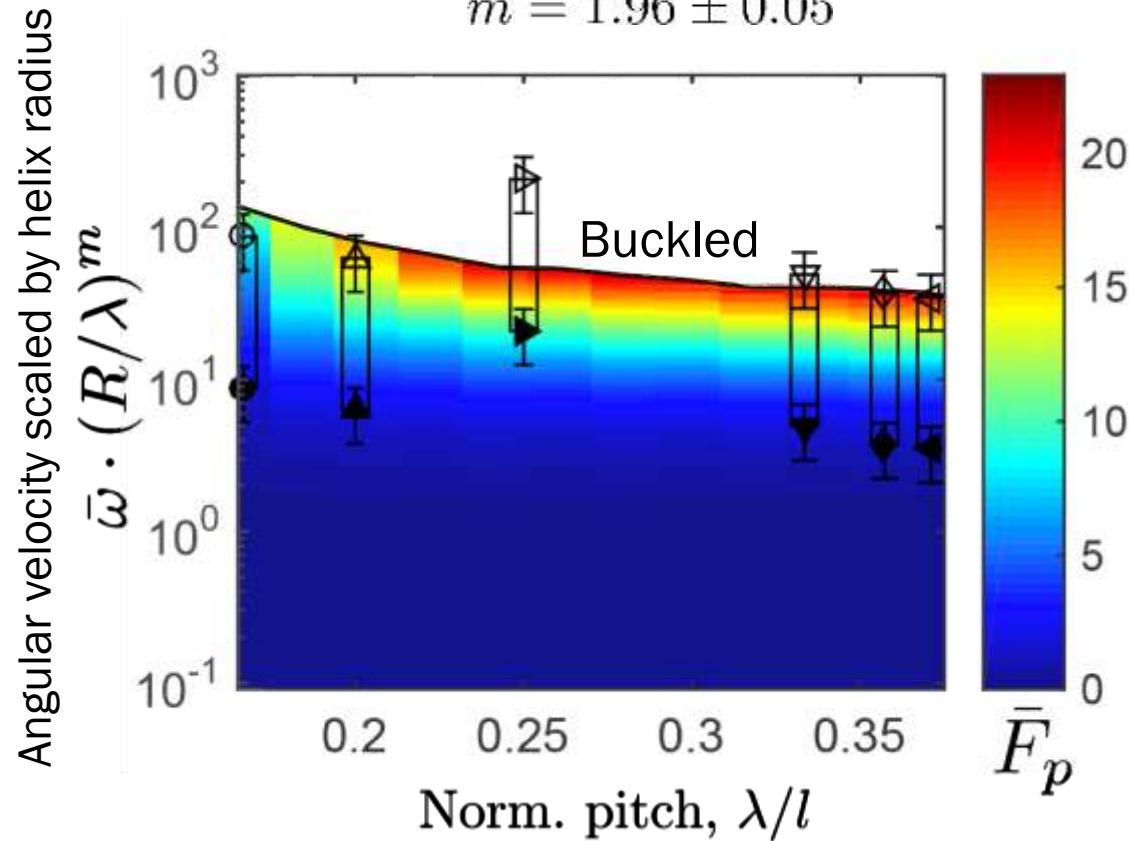
Geometry dependent coefficient

$$\bar{\omega}_b(R/\lambda, \lambda/l) = \hat{M}(R/\lambda) \hat{N}(\lambda/l)$$



Phase diagram

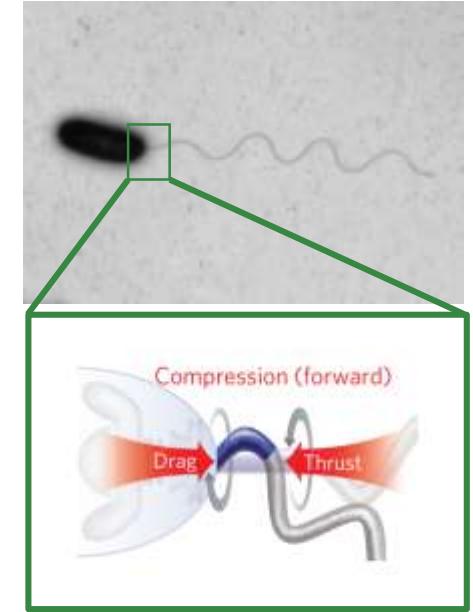
Norm. buckling velocity: $\bar{\omega}_b = (R/\lambda)^{-m} \hat{N}(\lambda/l)$
 $m = 1.96 \pm 0.05$



- Caulobacter crescentus (Wild)
- ◆ Escherichia coli
- ▼ Rhizobium lupini (Semicoiled)

- ▲ R. lupini (Curly)
- ▶ Salmonella (Wild)
- ◀ Vibrio alginolyticus

Failure to functionality

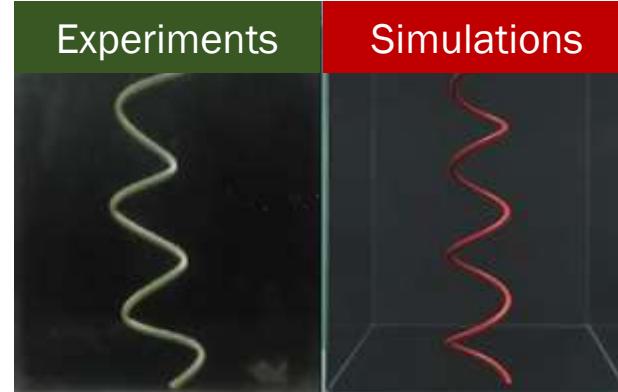
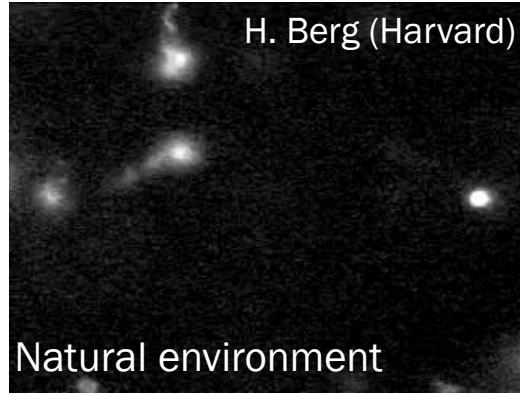


Son, K., Guasto, J., Stocker, R. (2013).
Bacteria can exploit a flagellar buckling
instability to change direction.
Nature physics, 9(8), 494-498.

Filled symbols: $EI \sim 10^{-23}$ N/m²
Y. Takano et al. JSME Int J., 2003

Open symbols: $EI \sim 10^{-24}$ N/m²
S. Fujime et al. J. Mol. Biol. (1972)

Conclusion



- Coupled effect of flexible structure + fluid
 - Structure: Discrete Elastic Rods
 - Fluid: Lighthill slender body theory
- Flagella can buckle above a threshold angular velocity
- Bacteria may exploit buckling for functionality

Jawed, M. K., Khouri, N. K., Da, F., Grinspun, E., & Reis, P. M.
Propulsion and instability of a flexible helical rod rotating in a viscous fluid
Physical Review Letters, 115(16):168101, 2015

Jawed, M. K. & Reis, P. M.
Deformation of a soft helical filament in an axial flow at low Reynolds number
Soft Matter, 2016

Email: khalidjm@mit.edu

Web: khalidjawed.com

Discrete Elastic Rods

Bergou et al, SIGGRAPH 2008
Bergou et al, SIGGRAPH 2010

Simulation loop

Letting $\mathbf{q} = (\mathbf{x}_0, \theta^0, \dots, \mathbf{x}_n, \theta^n, \mathbf{x}_{n+1})^T$ and $h_k = t_{k+1} - t_k$, we apply Euler integration (implicit on internal and explicit on external forces), using Newton's method to solve

$$\begin{aligned} M\Delta\dot{\mathbf{q}} - h_k \mathbf{F}_{\text{int}}(\mathbf{q}(t_k) + \Delta\mathbf{q}) &= h_k \mathbf{F}_{\text{ext}}(t_k, \mathbf{q}(t_k), \dot{\mathbf{q}}(t_k)) \\ \Delta\mathbf{q} - h_k \Delta\dot{\mathbf{q}} &= h_k \dot{\mathbf{q}}(t_k) \end{aligned}$$

Explicit energy gradient and Hessian

$$E_s = \frac{1}{2} \sum_{j=0}^n k_s^j \left(\varepsilon^j \right)^2 |\bar{\mathbf{e}}^j|,$$

$$E_t = \frac{1}{2} \sum_{i=1}^n \beta_i \frac{(m_i - \bar{m}_i)^2}{\bar{l}_i},$$

$$E_b = \frac{1}{2} \sum_{i=1}^n \frac{1}{\bar{l}_i} (\boldsymbol{\kappa}_i - \bar{\boldsymbol{\kappa}}_i)^T \mathbf{B}_i (\boldsymbol{\kappa}_i - \bar{\boldsymbol{\kappa}}_i)$$

Time-parallel frame

Given the reference frame at time t_{k-1} , the time-parallel reference frame $\underline{\mathbf{d}}_\alpha^j(t_k)$ only depends on $\mathbf{t}^j(t_k)$, and the only nonzero terms in (3) are $i \in \{j, j+1\}$. Thus, the force stencil is local and the energy Hessian is banded.